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# Evaluation of Human Exposure to Magnetic Fields Generated by Electric Power Systems in Complex Configurations 

Ph.D. Thesis
by

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#### Abstract

The international growing concern for the human exposure to magnetic fields generated by electric power lines has unavoidably led to imposing legal limits. Respecting these limits, implies being able to calculate easily and accurately the generated magnetic field also in complex configurations. Twisting of phase conductors is such a case. The consolidated exact and approximated theory regarding a single-circuit twisted three-phase power cable line has been reported along with the proposal of an innovative simplified formula obtained by means of an heuristic procedure. This formula, although being dramatically simpler, is proven to be a good approximation of the analytical formula and at the same time much more accurate than the approximated formula found in literature. The double-circuit twisted threephase power cable line case has been studied following different approaches of increasing complexity and accuracy. In this framework, the effectiveness of the above-mentioned innovative formula is also examined. The experimental verification of the correctness of the twisted double-circuit theoretical analysis has permitted its extension to multiple-circuit twisted three-phase power cable lines. In addition, appropriate 2D and, in particularly, 3D numerical codes for simulating real existing overhead power lines for the calculation of the magnetic field in their vicinity have been created. Finally, an innovative 'smart' measurement and evaluation system of the magnetic field is being proposed, described and validated, which deals with the experimentally-based evaluation of the total magnetic field $B$ generated by multiple sources in complex three-dimensional arrangements, carried out on the basis of the measurement of the three Cartesian field components and their correlation with the field currents via multilinear regression techniques. The ultimate goal is verifying that magnetic induction intensity is within the prescribed limits.


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## CHAPTER 1

## INTRODUCTION

Within the past years there has been very significant growth of man-made, extremely low frequency (ELF) magnetic fields at frequencies of 50 and 60 Hz predominantly from electric energy generation, transmission, distribution and use. Man-made ELF fields are many orders of magnitude greater than the naturally arising fields at 50 and 60 Hz . Within all organisms are endogenous electromagnetic fields and currents that play a role in the complex mechanisms of physiological control such as neuromuscular activity, glandular secretion, cellmembrane function and tissue development, growth and repair. It is not surprising that, because of the role of electromagnetic fields and currents in so many basic physiological processes, questions arise concerning possible effects of artificially produced fields on biological systems. With advances in technology and the ever increasing need for electric energy, human exposure to $50 / 60 \mathrm{~Hz}$ magnetic fields has increased to the point that valid questions are raised concerning safe limits of such exposure. Public concern is growing and in many countries regulatory and advisory agencies have been requested to evaluate possible adverse effects of ELF electromagnetic fields on human health.

Exposure standards have been developed internationally, that provide adequate protection against all known adverse effects of exposure to EMF. The guidelines developed by the International Commission on Non Ionizing Radiation Protection (ICNIRP) [1] are widely recognized and have formed the basis for national regulations in several countries. European Union has also established a common framework [2], [3] for giving the general public a high level of protection against the potential harmful effects of exposure to electromagnetic fields, particularly by limiting exposure to sources of non-ionizing radiation. However, social concerns for hypothesized long-term effects of chronic exposure to low-level EMF have created a demand for precautionary measures beyond the standards for recognized, acute effects.

As far as Italy is concerned, the limits currently in force for magnetic fields generated by 50 Hz power lines are set by [4]:
$\rightarrow$ The exposure limit: $100 \mu \mathrm{~T}$ (rms value). This limit must never be exceeded in case of general public exposure.
$\rightarrow$ The attention value: $10 \mu \mathrm{~T}$ (rms value). This limit is a cautionary measure, adopted in children's playgrounds, residential dwellings, school premises and in areas where people are staying for 4 hours or more per day, in order to protect against any possible long-term effects that might be related to power frequency ( 50 Hz ) magnetic fields. The attention value is the median of values recorded over 24 hours, under normal operational conditions.
$\rightarrow$ The quality objective: $3 \mu \mathrm{~T}$ (rms value). This limit is adopted for the purpose of minimizing progressively the exposures to magnetic fields generated by 50 Hz power lines and is applied in designing new power lines in the neighborhood of children's playgrounds, residential dwellings, school premises, and in areas where people are staying for 4 hours or more per day, as well as in planning developments in the proximity of existing electric power lines and installations, including the categories mentioned above. The quality objective is the median of values recorded over 24 hours, under normal operational conditions.

The above considerations have stimulated the search for methods of arranging the conductors of power lines in such a way that the surrounding magnetic fields will be greatly reduced. Focusing on power distribution cables, a solution adopted for practical reasons when laying the cables and that has proved to be very effective also in mitigating the magnetic field is twisting the phase conductors. Practical reasons and field mitigation effectiveness have made twisted three-phase cables the standard solution for medium voltage (MV) and low voltage (LV) power cables in the electric power distribution networks. In turn, at the distribution level both overhead and underground twisted three-phase cables are overwhelming traditional overhead lines with bare conductors, for many reasons, e.g.:

- Buildings are spreading over larger and larger areas, thereby making the use of traditional overhead lines possible over less and less broad areas;
- Overhead cables are often replacing traditional overhead lines also in rural areas, in order to reduce fault rate and maintenance of the lines;
- twisted cables are becoming more and more popular with the growing spread of renewable sources, since underground twisted three-phase cable lines are the preferred solution for connecting wind-generators and photovoltaic systems to the distribution grid, due to their very low environmental impact.

All this has unavoidably led to the ever increasing use of twisted three-phase cables.

The calculation of the low-frequency magnetic field in the vicinity of a three-wire twisted cable carrying three-phase current, even if it is drastically reduced, is essential for the evaluation of its impact near sensitive receptors and also for the calculation of the distance corresponding to a maximum limit value of the rms magnetic induction
[4], [5]. This calculation problem has been discussed in literature since 1937, when Buchholz gave an analytical solution for the twisted pair in form of an infinite series containing modified Bessel functions in [6] and [7]. Later on, several studies have been made and published, dealing mostly with approximations to the rather bulky series-type solution, some keeping only the dominant first term of the series [8] - [10] and others attempting direct ways of approximation [11]. Most recently, this issue has been treated analytically in [12] - [23] and particularly in [14] and [15] by Pettersson et al., who presented a complete and exact theory of the power-frequency magnetic field emitted by a twisted three-phase configuration and also provided experimental results for demonstrating the correctness of the analytical solution.

As already mentioned, the exact formula for the calculation of the magnetic field generated by a three-phase twisted configuration uses an infinite-term series of modified Bessel functions of the first and second kind and their derivatives in a reference system with cylindrical coordinates [14] - [15]. In the literature an approximate formula also exists that, for distances comparable to the pitch of the helix, gives results not far from those of the exact formula for distances exceeding, say, 1 m from the axis of the helix, and is in fact often used in these conditions. On the contrary, as the field-point approaches the twisted three-phase arrangement, the approximate formula results in steeply-increasing errors and the exact formula is recommended. Note that distances close to the conductors are of major interest in the case of MV and LV cables: indeed, such cables - in the form of overhead or underground power lines - often cross densely populated areas and are sometimes integrated within inhabited buildings; moreover, when the current rms value is relatively low then the generated level of the magnetic field is significant only in areas close to the conductors.

In this PhD Thesis, via an heuristic procedure, an innovative simplified formula for the rms magnetic induction is obtained which is much simpler than the rigorous analytical one and provides a much smaller relative error compared to the approximated one from the literature, especially for small distances from the helix axis [24] - [27]. The effectiveness of the innovative expression is evaluated by carrying out some numerical simulations relevant to a typical MV cable in order to compare the results provided by the exact and the approximate formulae from the literature with those obtained via the innovative simplified formula.

Subsequently, due to the increasing interest of utilities worldwide in double-circuit twisted overhead, as well as underground, cable lines, the theoretical treatment of a double-circuit twisted cable line was effectuated, following different approaches of increasing complexity and accuracy, all based on the superposition of the effects of each single-circuit twisted three-phase cable. Resorting to the correct vectorial approach, the total magnetic field is calculated as the vector sum of the two fields generated by each twisted cable using the exact expressions of the three field components generated by each cable, consisting in the above-described series of

Bessel functions. The objective difficulties derived from this calculus - such as knowing the precise geometrical arrangement of the three phases or being sure of the constancy of the pitch of the single helix and also of the constancy of the perfect parallelism of the two helixes - make the resulting formula far more complex than that relevant to one single three-phase twisted cable. As a second and rather necessary step for the purpose of simplicity, based on the "worst case assumption", the rms magnetic field $B$ generated by a double-circuit twisted cable line is calculated as the algebraic sum of the two total fields from the single-circuits, each computed via the exact, the approximated from the literature and the innovative simplified formula, thereby providing the most conservative value of $B$ that serves as an upper reference limit for exposure evaluation [28] - [32]. The analytic theory is supported by numerical evaluation as well as experimental results by actual measurements in situ.

Finally, the calculation of the magnetic field is further extended to a multiplecircuit twisted cable line, again following both the exact vectorial and the "worst case" approach where the field from each cable circuit is computed via both the exact and the innovative simplified formula, thereby providing the most conservative value of $B$. The effectiveness of the simplified formula for the case of a triple-circuit twisted cable lines is also proven by the simulation of some practical case-studies related to MV triple-circuit twisted cable lines [33].

The calculation of the magnetic field generated by overhead power lines, [34] [51], was also treated in this PhD Thesis for the sake of creating appropriate bi- and tri-dimensional numerical codes. The bi-dimensional calculation of the magnetic field in the vicinity of overhead power lines overestimates the field since it does not consider the catenary form of the line and effectuates the calculation at the point of maximum exposure, i.e. at the mid-span between two adjacent towers. Doing so, the two-dimensional numerical codes for the calculation of magnetic field provide isolines of magnetic induction at $100 \mu \mathrm{~T}, 10 \mu \mathrm{~T}$ and $3 \mu \mathrm{~T}$ that spread over wider areas compared to the real ones, leading to a reduction in beneficial construction space. On the other hand, the tri-dimensional calculation of the magnetic field examines the real geometrical configuration of the power line, providing thus, results closer to the real ones.

In the framework of determining the general public exposure to magnetic fields from complex 3D arrangements of the field sources from power systems - particularly overhead transmission and distribution lines - measurement and evaluation systems of the magnetic field generated are needed that can provide reliable and accurate indications. Such arrangements occur more and more frequently in the vicinity of residential and industrial areas, where overhead and underground lines of different voltage rating and geometry, as well as more or less wide substations, often lie close to each other. The final part of this Thesis proposes an innovative measurement and evaluation system capable of matching this need. The innovative "smart" measurement device measures and also records the rms values of the magnetic field components $B_{x}$,
$B_{y}$ and $B_{z}$ as a function of time (unlike the usual magnetic field measuring devices that record and store only the rms value of total magnetic induction field $B$ ) for subsequent analytical processing. The three-dimensional magnetic field measurement device has been entirely created, calibrated and characterized in the laboratory [52] - [56], and has been tested with success by experimental campaigns in the field in the presence of multiple current sources. The innovative evaluation system uses multilinear regression algorithms developed in Matlab ${ }^{\circledR}$ environment for correlating the measurement results of $B_{x}, B_{y}$ and $B_{z}$ at a given field point to the relevant values of time-varying currents enabling, in this way, the extrapolation of the rms value of the total magnetic field to any combination of source currents of interest [57] - [59]. The best performances of that innovative measurement and evaluation system have been published [60], [61] and discussed.

## CHAPTER 2

## Single - Circuit Twisted Three - Phase Cable Lines

### 2.1 THE EXACT FORMULATIONS

### 2.1.1 Single Wire Helix

According to the Biot-Savart law, the magnetic flux density vector of a helical line source carrying current $I$ is given by the line integral

$$
\begin{equation*}
B=\frac{\mu_{0} I}{4 \pi} \int \frac{d \vec{r} \times\left(r-\vec{r}^{\prime}\right)}{\left|r-\vec{r}^{\prime}\right|^{3}} \tag{2.1}
\end{equation*}
$$

along the helix, where $\mu_{0}$ is the permeability of the free space $\left(\mu_{0}=4 \pi 10^{-7} \mathrm{H} / \mathrm{m}\right), \vec{r}$ is the field point, $\vec{r}^{\prime}$ is the source point variable, $p$ is the helix pitch and $\alpha$ is the helix radius (see Fig. 2.1).


Figure 2.1: Helical line current (after [15]).
This integral cannot be calculated analytically in a direct manner. However, the integrand can be series expanded in terms whose integrals can be given in the form of Bessel functions.


Figure 2.2: Field components generated by a one-wire helix, expressed in cylindrical coordinates (after [14]).

The exact theory regarding this configuration consists in calculating the magnetic field in cylindrical coordinates $r, \varphi$ and $z$ (respectively radial, azimuthal and axial, see Fig. 2.2), using the following expressions for the components $B_{r}, B_{\varphi}$, e $B_{z}$ of the magnetic induction field [15]:

$$
\begin{gather*}
B_{r}=\frac{\mu_{0} I \alpha}{\pi r^{2}}(k r)^{2} \sum_{n=1}^{\infty} n I_{n}^{\prime}(n k \alpha) K_{n}^{\prime}(n k r) \sin \left[n\left(\phi-\phi_{0}-k z\right)\right]  \tag{2.2.a}\\
B_{\phi}=\frac{\mu_{0} I}{2 \pi r}+\frac{\mu_{0} I \alpha}{\pi r^{2}}(k r) \sum_{n=1}^{\infty} n I_{n}^{\prime}(n k \alpha) K_{n}(n k r) \cos \left[n\left(\phi-\phi_{0}-k z\right)\right]  \tag{2.2.b}\\
B_{z}=-\frac{\mu_{0} I \alpha}{\pi r^{2}}(k r)^{2} \sum_{n=1}^{\infty} n I_{n}^{\prime}(n k \alpha) K_{n}(n k r) \cos \left[n\left(\phi-\phi_{0}-k z\right)\right] \tag{2.2.c}
\end{gather*}
$$

where $k=2 \pi / p$ and $I_{n}(z), K_{n}(z)$ are the modified Bessel functions of first and second kind of order $n$, and $I_{n}{ }^{\prime}(z), K_{n}{ }^{\prime}(z)$ their derivatives.

For very large distances, the three Bessel function sums may be shown to go to zero much faster than the inverse $-r$ term of $B_{\varphi}$, so that the only field component ultimately left will be $\mathrm{B}_{\phi}=\mu_{0} I / 2 \pi r$ which is the field of an infinitely long straight conductor carrying current $I$. That means that the field from a single conductor does not vanish simply by twisting it.

For $p \rightarrow 0$, the current distribution will approach a purely azimuthal surface current on the cylinder surface. In this case, all three field components vanish, which is expected from the solenoids theory where the field is concentrated to the inner of the cylinder.

### 2.1.2 Two Wire Helix

Without loss of generality, by setting $\varphi_{0}=0$ for the conductor carrying current $I$ and $\varphi_{0}=\pi$ for the conductor carrying current $-I$, the field from each of the helices is given by eqs. (2.2.a), (2.2.b) and (2.2.c) and the total field is found by summation. Thus as the even order terms will cancel pairwise while the odd will double, the following expressions derive [15]:

$$
\begin{align*}
& B_{r}=2 \frac{\mu_{0} I \alpha}{\pi r^{2}} \gamma^{2} \sum_{n}^{\infty} n I_{n}^{\prime}(n \eta) K_{n}^{\prime}(n \gamma) \sin [n(\phi-k z)]  \tag{2.3.a}\\
& B_{\phi}=2 \frac{\mu_{0} I \alpha}{\pi r^{2}} \gamma \sum_{n}^{\infty} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \cos [n(\phi-k z)]  \tag{2.3.b}\\
& B_{z}=-2 \frac{\mu_{0} I \alpha}{\pi r^{2}} \gamma^{2} \sum_{n}^{\infty} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \cos [n(\phi-k z)] \tag{2.3.c}
\end{align*}
$$

where the summations range over $n=1,3,5$. For sinusoidal current of angular velocity $\omega, I$ is to be replaced by $\hat{I} \sin (\omega t)$, where $\widehat{I}$ is the peak value. Up to sign, the expressions will also hold as they stand for the effective values of the $B$-components with $I$ denoting the effective value of the current.

### 2.1.3 Three Wire Helix

For the three-phase case-study the conductor arrangement in a transverse plane will form an equilateral triangle since the three coaxial helices have equal radii. Therefore, let us number the wires as $i=1,2,3$ and set $\alpha_{i}=(i-1) 2 \pi / 3, \quad \phi_{i}=(i-1) 2 \pi / 3$ and $I_{i}=\hat{I} \sin \left(\omega t+\alpha_{i}\right)$, where $\alpha_{i}$ are the phase angles of the three helices and $\phi_{i}$ are the location parameters of the three helices.Term-wise addition of the three fields yields after some elementary calculations using the auxiliary geometric relations

$$
\begin{align*}
& \sum_{i=1}^{3} \sin \left(\omega t+a_{i}\right) \sin \left[n\left(\phi-\phi_{i}-k z\right)\right]=\mp \frac{3}{2} \cos (\omega t \pm n \Phi)  \tag{2.4.a}\\
& \sum_{i=1}^{3} \sin \left(\omega t+a_{i}\right) \cos \left[n\left(\phi-\phi_{i}-k z\right)\right]=\frac{3}{2} \sin (\omega t \pm n \Phi) \tag{2.4.b}
\end{align*}
$$

with $\Phi=\phi-k z$
for $n=1,2,4,5,7 \ldots$ and zero for $n=3,6,9 \ldots$, where the upper sign applies for $n=$ $2,5,8 \ldots$ and the lower for $n=1,4,7 \ldots$.

Thus, the time-dependent field components are [15]:

$$
\begin{align*}
& B_{r}=\frac{3}{2} \frac{\mu_{0} \hat{I} \alpha}{\pi r^{2}} \gamma^{2} \sum_{n}(\mp n) I_{n}^{\prime}(n \eta) K_{n}^{\prime}(n \gamma) \cos (\omega t \pm n \Phi)  \tag{2.5.a}\\
& B_{\phi}=\frac{3}{2} \frac{\mu_{0} \hat{I} \alpha}{\pi r^{2}} \gamma \sum_{n} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \sin (\omega t \pm n \Phi)  \tag{2.5.b}\\
& B_{z}=-\frac{3}{2} \frac{\mu_{0} \widehat{I} \alpha}{\pi r^{2}} \gamma^{2} \sum_{n} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \sin (\omega t \pm n \Phi) \tag{2.5.c}
\end{align*}
$$

The effective values of the field components can be seen to be given by [15]:

$$
\begin{align*}
& B_{r}=\frac{3}{2} B_{0} \gamma^{2}\left[\sum_{n} \sum_{m}(\mp n)(\mp m) I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta) K_{n}^{\prime}(n \gamma) K_{m}^{\prime}(m \gamma) \cos ( \pm n \mp m) \Phi\right]^{1 / 2}  \tag{2.6.a}\\
& B_{\phi}=\frac{3}{2} B_{0} \gamma\left[\sum_{n} \sum_{m} n m I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta) K_{n}(n \gamma) K_{m}(m \gamma) \cos ( \pm n \mp m) \Phi\right]^{1 / 2}  \tag{2.6.b}\\
& B_{z}=\frac{3}{2} B_{0} \gamma^{2}\left[\sum_{n} \sum_{m} n m I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta) K_{n}(n \gamma) K_{m}(m \gamma) \cos ( \pm n \mp m) \Phi\right]^{1 / 2} \tag{2.6.c}
\end{align*}
$$

with

$$
\begin{equation*}
B_{0}=\frac{\mu_{0} I \alpha}{\pi r^{2}} \tag{2.7}
\end{equation*}
$$

where the index and sign rules of above apply and $I$ denotes the effective value of the current.

### 2.1.4 Simplification of the Three Wire Helix Case

Considering an imagined helix of pitch $p$ through the field-point (see Fig. 2.3 after [14]) and using the helical coordinates $r$ (radial), $b$ (binormal) and $s$ (tangential) the field can be described by only two components, i.e. the radial component $B_{r}$ and the binormal component $B_{b}$, since the tangential component $B_{s}$ is zero.


Figure 2.3: Field components in helical coordinates (after [14]).
The relation between the two reference systems is given by

$$
\begin{align*}
& B_{s}=B_{z} \sin \psi+B_{\phi} \cos \psi  \tag{2.8.a}\\
& B_{b}=B_{z} \cos \psi-B_{\phi} \sin \psi \tag{2.8.b}
\end{align*}
$$

with $\psi=\tan ^{-1}(k r)$ the pitch angle of the field-point helix.
Now, cancelling the single term $\mu_{0} I / 2 \pi r$ of $B_{\varphi}$ by applying an imagined return current on the cylinder axis, the time dependent field-vector components can be written as follows [14]:

$$
\begin{gather*}
B_{r}=B_{0} \gamma^{2} \sum_{n=1}^{\infty} n I_{n}^{\prime}(n \eta) K_{n}^{\prime}(n \gamma) \sin (n \Phi) \sin (\omega t)  \tag{2.9.a}\\
B_{b}=-B_{0} \gamma \sqrt{1+\gamma^{2}} \sum_{n=1}^{\infty} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \cos (n \Phi) \sin (\omega t) \tag{2.9.b}
\end{gather*}
$$

where $B_{0}=\mu_{0} I_{p} \alpha / \pi r^{2}, \eta=k \alpha, \gamma=k r, \Phi=\phi-\phi_{0}-k z$ and $I_{p}$ is the peak value of the sinusoidal current $i(t)=I_{p} \sin \omega t$ of angular frequency $\omega$.

For the three-phase case-study we number the wires by $i=1,2,3$ and set $I_{i}=\hat{I} \sin \left(\omega t+\alpha_{i}\right), \alpha_{i}=(i-1) 2 \pi / 3, \phi_{i}=(i-1) 2 \pi / 3$, where $\alpha_{i}$ are the phase angles of the three helices and $\phi_{i}$ are the location parameters of the three helices.

Term-wise addition of the three fields yields after calculations using the auxiliary geometric expressions (2.4) and the following results represent the time-dependent field components [14]:

$$
\begin{align*}
& B_{r}=\frac{3}{2} \frac{\mu_{0} I_{p} \alpha}{\pi r^{2}} \gamma^{2} \sum_{n}(\mp n) I_{n}^{\prime}(n \eta) K_{n}^{\prime}(n \gamma) \cos (\omega t \pm n \Phi)  \tag{2.10.a}\\
& B_{b}=-\frac{3}{2} \frac{\mu_{0} I_{p} \alpha}{\pi r^{2}} \gamma \sqrt{1+\gamma^{2}} \sum_{n} n I_{n}^{\prime}(n \eta) K_{n}(n \gamma) \sin (\omega t \pm n \Phi) \tag{2.10.b}
\end{align*}
$$

where $\Phi=\phi-k z$ and the summations range over $n=1,2,4,5,7 \ldots$ i.e. all positive integers except $n=3,6,9 \ldots$; the upper sign applies for $n=2,5,8 \ldots$ and the lower for $n=1,4,7 \ldots$

The rms values of the components and of the total field $B$ are given by [14]:

$$
\begin{gather*}
B_{r}=\frac{3}{2} B_{0} \gamma^{2}\left[\sum_{n} \sum_{m}(\mp n)(\mp m) I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta) K_{n}^{\prime}(n \gamma) K_{m}^{\prime}(m \gamma) \cos ( \pm n \mp m) \Phi\right]^{1 / 2}  \tag{2.11.a}\\
B_{b}=\frac{3}{2} B_{0} \gamma \sqrt{1+\gamma^{2}}\left[\sum_{n} \sum_{m} n m I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta) K_{n}(n \gamma) K_{m}(m \gamma) \cos ( \pm n \mp m) \Phi\right]^{1 / 2}  \tag{2.11.b}\\
B=\frac{3}{2} B_{0} \gamma^{2}\left\{\sum_{n} \sum_{m} n m I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta)\left[(\mp 1)(\mp 1) K_{n}^{\prime}(n \gamma) K_{m}^{\prime}(m \gamma)+\frac{1+\gamma^{2}}{\gamma^{2}} K_{n}(n \gamma) K_{m}(m \gamma)\right] \cos ( \pm n \mp m) \Phi\right\}^{1 / 2} \tag{2.12}
\end{gather*}
$$

Formulae (2.11) - (2.12) can be applied to the special case of an untwisted configuration (with $\eta \rightarrow 0, \gamma \rightarrow 0$ and $p \rightarrow \infty$ ) using the small argument approximations for the Bessel functions:

$$
\begin{equation*}
I_{n}(\eta)=\frac{(n / 2)^{n}}{n!} K_{n}(\gamma)=\frac{1}{2} \frac{(n-1)!}{(\gamma / 2)^{n}} \text { with } \eta, \gamma \ll 1 \tag{2.13}
\end{equation*}
$$

The result is [14]:

$$
\begin{gather*}
B_{r}=\frac{3}{2} B_{0}\left\{\sum_{n} \sum_{m}\left(\frac{\alpha}{r}\right)^{n+m-2} \frac{1}{4}(\mp 1)(\mp 1) \cos ( \pm n \mp m) \Phi\right\}^{1 / 2}  \tag{2.14.a}\\
B_{b}=\frac{3}{2} B_{0}\left\{\sum_{n} \sum_{m}\left(\frac{\alpha}{r}\right)^{n+m-2} \frac{1}{4} \cos ( \pm n \mp m) \Phi\right\}^{1 / 2} \tag{2.14.b}
\end{gather*}
$$

$$
\begin{equation*}
B=\frac{3}{2} B_{0}\left\{\sum_{n} \sum_{m}\left(\frac{\alpha}{r}\right)^{n+m-2} \frac{1}{4}[(\mp 1)(\mp 1)+1] \cos ( \pm n \mp m) \Phi\right\}^{1 / 2} \tag{2.15}
\end{equation*}
$$

### 2.2 THE APPROXIMATE FORMULATION FROM THE LITERATURE

For certain values of parameters $\alpha$ and $p$, the first term of the series expansions will be so dominant that it can serve as an approximation of the whole sum for proper values of the variable $r$. One important example of such a situation is when the configuration has a loose twist $(\alpha \ll p)$ and the field-point is far from the axis of the helix ( $r \gg p$, hence $\gamma \gg 1$ ). In this case the total field reduces to [14], [15]:

$$
\begin{equation*}
B=\frac{3}{2} B_{0} \gamma^{2} I_{n}^{\prime}(\eta)\left[K_{1}^{2}(\gamma)+\frac{1+\gamma^{2}}{\gamma^{2}} K_{1}^{2}(\gamma)\right]^{1 / 2} \tag{2.16}
\end{equation*}
$$

Using the small argument approximation $I_{n}{ }^{\prime}(\eta)=1 / 2$, holding for $\eta \rightarrow 0$, and the large argument approximation $K_{1}(\gamma) \approx-K_{1}^{\prime}(\gamma) \approx \sqrt{2 /(\pi \gamma)} \mathrm{e}^{-\gamma}$ holding for $\gamma \gg 1$, one obtains from eq. (2.16) [14], [15]:

$$
\begin{equation*}
B \approx 3 / 4 F \sqrt{2} B_{0} \tag{2.17}
\end{equation*}
$$

where $F$ is the so-called "twist factor" [14], [15]:

$$
\begin{equation*}
F=\sqrt{\pi / 2} \gamma^{3 / 2} \mathrm{e}^{-\gamma} \tag{2.18}
\end{equation*}
$$

Formulae (2.17) - (2.18) illustrate that the field reduction involved by twisting the conductors is expressed by the twist factor $F$. For example, for distance to pitch ratios, $r / p$, equal to 1 and $2, F$ is equal to 0.037 and 0.0002 respectively. Thus, a dramatic decay of field with distance is observed.

### 2.3 THE SIMPLIFIED INNOVATIVE FORMULATION

An unavoidable preliminary step was to implement in Matlab ${ }^{\text {TM }}$ environment both the approximate (eq. (2.17)) and the exact (eq. (2.12)) expressions, so as to reproduce the literature values of magnetic field reported in [14], [15]. Since the exact formula of eq. (2.12) was in turn verified satisfactorily via a sound experimental campaign
illustrated in detail in [14], [15], it can be argued that the validation of the script by reproducing the literature values of magnetic field according to both the approximate (eq. (2.17)) and the exact (eq. (2.12)) expressions - is fundamental. An extensive comparison of the results provided by the implemented script with those reported in [14], [15] has been carried out in [24] - [27], not only from the viewpoint of total $B$ field, but also of radial and binormal components, and has provided excellent results.

### 2.3.1 Angle $\Phi$

In search for an innovative simplified expression of the magnetic induction, it is convenient to analyze the behaviour of $B$ as a function of the distance $r$ only. To do so, the dependence of the total field and the relevant components on angle $\Phi$ (which indicates the angular position of the field-point in the twisted configuration) has been studied in [24], [25]. The results are depicted in Fig. 2.4 and can be summarised as follows:

1. the radial component $B_{r}$ prevails in the composition of the total magnetic field $B$;
2. the radial component $B_{r}$, and therefore the total magnetic field $B$, reaches its maximum value at angles $\Phi=60^{\circ}, \Phi=180^{\circ}$ and $\Phi=300^{\circ}$. At these angle values the azimuthal, $B_{\varphi}$, and the axial, $B_{z}$, components reach their minimum value;
3. the radial component $B_{r}$, and therefore the total magnetic field $B$, presents its minimum value at angles $\Phi=0^{\circ}, \Phi=120^{\circ}$ and $\Phi=240^{\circ}$. At these angle values the azimuthal, $B_{\varphi}$, and the axial, $B_{z}$, components reach their maximum value.


Figure 2.4: $B, B_{r}, B_{\varphi}, B_{z}$ vs. angle $\Phi(r=0.3 m, I=280 A, p=1.37 m, \alpha=0.02 m)$

So, by setting and keeping constant $\Phi=60^{\circ}$ during the search of the innovative simplified formula, the maximum value of the total magnetic field $B$ at every field point was guaranteed, providing, thus, a conservative estimation of the magnetic field.

### 2.3.2 Simplification of the Bessel Functions

In order to find an alternative approximate expression that simplifies the exact calculation, it is useful to set a finite maximum value for the indexes $n$ and $m$ of the Bessel functions in eq. (2.12).Such maximum value has to be high enough to let the approximate solution converge to the exact one, but low enough to make the approximate solution easily computable. From an analysis of the convergence of $B$ it can be deduced that:

1. with increasing $n$ and $m$, the approximate solution tends to the exact one;
2. the first-order approximation ( $n=m=1$ ) corresponds to the approximate solution given by eq. (2.16). This approximation of the magnetic field does not depend on the angle $\Phi$ and seems to give a good average value of the field $B$;
3. for greater distances the speed of convergence to the exact solution increases. In particular, for distances $>1 \mathrm{~m}$ the second-order approximation ( $n=m=2$ ) is an excellent approximation of the exact solution. The second-order approximation results always better than the first-order one and moreover always overestimates the field compared to the exact one (which is desirable from an engineering point of view, since it is conservative).

By setting:

$$
\begin{equation*}
A=\left\{\sum_{n} \sum_{m} n m I_{n}^{\prime}(n \eta) I_{m}^{\prime}(m \eta)\left[(\mp 1)(\mp 1) K_{n}^{\prime}(n \gamma) K_{m}^{\prime}(m \gamma)+\frac{1+\gamma^{2}}{\gamma^{2}} K_{n}(n \gamma) K_{m}(m \gamma)\right] \cos ( \pm n \mp m) \Phi\right\}^{1 / 2} \tag{2.19}
\end{equation*}
$$

equation (2.12) simplifies as follows:

$$
\begin{equation*}
B=\frac{3}{2} B_{0} \gamma^{2} A=\frac{3}{2} \frac{\mu_{0} I \alpha}{\pi}\left(\frac{2 \pi}{p}\right)^{2} A=\frac{2.4 \pi^{2} I \alpha}{p^{2}} A \tag{2.20}
\end{equation*}
$$

with $\mu_{0}=4 \pi 10^{-1} \mu \mathrm{H} / \mathrm{m}$ in order for $B$ to result directly in $\mu \mathrm{T}$.
The term $A$ is the square root that contains the series of the Bessel functions and the rms magnetic induction $B$ is directly proportional to it. Since the quantity $2.4 \pi^{2} I \alpha / p^{2}$ does not depend on the distance $r$, the root $A$ is the single term that determines the dependence of the magnetic field on distance $r$.

For all distances, the second-order $(n=m=2)$ approximation $B_{2}$ of the exact value of magnetic induction $B$ - based on the second-order approximation $A_{2}$ of the exact value of $A$ - implies always a percent relative error $e_{2}$ smaller than that implied by the firstorder approximation; $e_{2}$ is defined as:

$$
\begin{equation*}
e_{2}(r)=100 \frac{A_{2}(r)-A(r)}{A(r)}=100 \frac{B_{2}(r)-B(r)}{B(r)} \tag{2.21}
\end{equation*}
$$

Moreover, what is important from a practical point of view, $e_{2}$ is always low and positive. The value of $A_{2}$ derived from the second-order approximation is always greater than the exact value of $A$, thereby leading to a precautionary overestimation of the magnetic field. Therefore, the exact expression of the total magnetic field $B$ can be reasonably approximated via the second-order approximation, as proved by the analysis made in [24], [25].

### 2.3.3 The Innovative Formula

A fundamental clue for finding the innovative formula has been gained by plotting the curve of $B$ vs. the radial distance from the helix axis, $r$, in semi-logarithmic coordinates. Such dependence is nearly linear, with a slight deviation from linearity (of hyperbolic type) only for small values of $r$. This means that the field can be approximated with the equation of a straight line plus a hyperbolic term which vanishes rapidly by increasing $r$. Thus the following innovative simplified expression has been conceived:

$$
\begin{equation*}
\ln B=a_{0}^{\prime}+a_{1}^{\prime} r+\frac{a_{2}^{\prime}}{r^{a_{3}^{\prime}}} \tag{2.22}
\end{equation*}
$$

Subsequently, on the basis of eq. (2.20) relationship (2.22) has been reprocessed as follows:

$$
\begin{equation*}
\ln B \approx \ln B_{2}=\ln \left(\frac{2.4 \pi^{2} \alpha}{p^{2}}\right)+\ln I+\ln A_{2} \approx \ln G+\ln I+\left(a_{0}+a_{1} r+\frac{a_{2}}{r^{a_{3}}}\right) \tag{2.23}
\end{equation*}
$$

being:

- $\ln I=f(I)$, known and trivial function of $I$;
- $\ln G=\ln \left(2.4 \pi^{2} \alpha / p^{2}\right)=f(\alpha, p)$, known function of $\alpha, p$;
- $\quad \ln A_{2} \approx a_{0}+a_{1} r+a_{2} / r^{a_{3}}=f(r ; \alpha, p)$, known function of $r$ and unknown function of $\alpha, p$.
$B$ is proportional to $I$, so the dependence of $\ln B$ on the current $I$ is trivially expressed by using the natural logarithm of $I$. The dependence on helix radius and pitch is expressed partly in an exact form via the logarithm of $G$, where $G=2.4 \pi^{2} \alpha / p^{2}$, and partly in an approximated form via the logarithm of $A_{2}$. In this way, only the square root of $A_{2}$ is approximated and not the entire expression of the field $B_{2}$.

The unknown parameters of the innovative approximate expression (2.23) are $a_{0}$, $a_{1}, a_{2}$ and $a_{3}$. These parameters have been derived as a function of $\alpha$ and $p$ with best-
fitting techniques developed in Matlab ${ }^{\mathrm{TM}}$ and aimed at minimizing the overall error on $B$ as a function of $r$, under the constraint that a positive error, possibly $>10 \%$, is to be privileged. Accurate values for parameters $a_{1}$ and $a_{0}$ are essential, since they govern the linear part of the $\ln B_{2}$ vs. $r$, and for this purpose an appropriate solution has been to limit the fitting of the linear part to the interval [ $1.2 \mathrm{~m} ; 2.0 \mathrm{~m}$ ] instead of the entire range $[0.3 \mathrm{~m} ; 2.0 \mathrm{~m}]$. After several tests, the value of $a_{2}=0.1$ has been set, while for parameter $a_{3}$ the value that minimizes the overall error has been chosen case by case with varying $\alpha, p$ according to the above requirement of a positive error not exceeding $10 \%$.

The three-dimensional Figures 2.5-2.7 show the values of parameters $a_{0}, a_{1}, a_{3}$ obtained by means of the above-described best-fitting procedure as a function of $p \in[0.8 \mathrm{~m} ; 2.0 \mathrm{~m}]$ and $\alpha \in[0.01 \mathrm{~m} ; 0.10 \mathrm{~m}]$; note that such wide intervals include most - if not all - the values of $\alpha, p$ that have practical interest.


Figure 2.5: Parameter $a_{0}$ versus radius $\alpha$ and pitch $p$ of the helix.


Figure 2.6: Parameter $a_{l}$ versus radius $\alpha$ and pitch $p$ of the helix.


Figure 2.7: Parameter $a_{3}$ versus radius $\alpha$ and pitch $p$ of the helix.
A second phase of analysis is based on the fact that the parameters $a_{0}, a_{1}$ and $a_{3}$ are functions of both pitch and radius of the helix [24], [25]. Therefore, using the techniques outlined above for the minimization of the error with respect to the exact formulation (eq. (2.12)), analytical approximated expressions of $a_{0}, a_{1}$ and $a_{3}$ have been obtained that represent the dependence of these parameters on both radius and pitch of the helix. To do so, a great number of simulations has been performed, taking again into consideration the various types of cables commercially available, each one characterized by different values of $\alpha$ and $p$, and so by different values of the parameters $a_{0}, a_{1}$ and $a_{3}$. All these simulations show that parameters $a_{0}, a_{1}$ and $a_{3}$ exhibit the following common behaviour for all the cables treated (see Appendix 1).

Parameter $a_{0}$ is a second-degree function of the pitch $p$, so an adequate representation of its dependence on the pitch is as follows: $a_{0}=a_{00}+a_{01} \cdot p+a_{02} \cdot p^{2}$; parameters $a_{00}, a_{01}$ and $a_{02}$ are in their turn second-degree functions of the radius $\alpha$. Considering all the above (Fig. 2.5) the following analytical approximated expression of $a_{0}$ as a function of p and $\alpha$ has been obtained:

$$
\begin{align*}
& a_{0}=\left[\left(12.8712 \cdot \alpha^{2}\right)+(0.2107 \cdot \alpha)-0.1383\right] \cdot p^{2}+ \\
& +\left[\left(-48.4318 \cdot \alpha^{2}\right)+(-0.6919 \cdot \alpha)+0.9094\right] \cdot p+  \tag{2.24}\\
& +\left[\left(49.6932 \cdot \alpha^{2}\right)+(0.5724 \cdot \alpha)-1.3991\right]
\end{align*}
$$

As far as parameter $a_{1}$ is concerned, the results of all the cases of the cables treated show that the dependence of this parameter on the radius $\alpha$, with $p$ fixed, can be represented as an horizontal straight line. Thus parameter $a_{1}$ depends only on $p$ and more precisely it is a third-degree function of the pitch $p$. Taking into consideration these arguments (Fig. 2.6) and minimizing the error, the following analytical approximated expression of $a_{1}$ as a function of $p$ has been obtained:

$$
\begin{equation*}
a_{1}=\left(2.1018 \cdot p^{3}\right)+\left(-11.6186 \cdot p^{2}\right)+(23.2879 \cdot p)-20.4446 \tag{2.25}
\end{equation*}
$$

Finally, all the cables treated show that parameter $a_{3}$ is a function of both pitch and radius of the helix (Fig. 2.7): in fact, $a_{3}$ is a first-degree function of the pitch $p$ so it can be approximated with a straight line $a_{3}=a_{30}+a_{31} p$, where the parameters $a_{30}$ and $a_{31}$ are also first-degree functions of the radius $\alpha$. Minimizing the error, the following analytical expression of $a_{3}$ as a function of $p$ and $\alpha$ has been obtained:

$$
\begin{equation*}
a_{3}=(0.2739 \cdot \alpha+0.2430) \cdot p+(0.1036 \cdot \alpha)+1.4444 \tag{2.26}
\end{equation*}
$$

At this point, the heuristic parametric analysis is concluded having established two levels of approximation of different complexity:

1. the first level of approximation, which consists in calculating the numerical values of parameters $a_{0}, a_{1}$ and $a_{3}$ for various types of twisted cables commercially available, therefore for a defined set of values for pitch and radius;
2. the second level of approximation, which consists in explaining the dependence of these parameters on radius and pitch of the helix by defining the approximate functional relationships based on the numerical values of $a_{0}, a_{1}$ and $a_{3}$ previously obtained. The advantage of the $2^{\text {nd }}$ level of approximation is the analytical representation of the dependence of $B$ on $p$ and $\alpha$, whereas its disadvantage is the greater error on the value of $B$ because of the further level of heuristic approximation.

### 2.4 APPLICATIONS - SIMULATION RESULTS

The simplified innovative formula developed here is applied to a type of threephase twisted cable frequently used in MV in Italy, named ARG7H1RX after [63], with the following main characteristics: conductor cross-section $3 \times 120 \mathrm{~mm}^{2}$, rated voltage 12 (phase to ground) $/ 20$ (phase-to-phase) kV , ampacity $I=280 \mathrm{~A}$, pitch $p=1.37$ m and radius $\alpha=0.02 \mathrm{~m}$.

Regarding the considered cable used in overhead power lines, the values of $B_{\text {exact }}$, $B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ as a function of the distance $r$ from the axis of the helix are shown in Fig. 2.8 in linear coordinates and in Table 2.1 along with the percent errors of $B_{l i t}$, $B_{\text {simpl }}, B_{\text {simp } 2}$ with respect to $B_{\text {exact }}$. The error of the simplified innovative formulain in the first level of approximation is always much smaller than the error of the approximate formula from the literature, never exceeding a few $\%$ even for small values of the distance $r$ apart a value $\sim 10 \%$ for $r=0.3 \mathrm{~m}$. As to the second level of approximation, the same considerations hold, apart distances higher than, say, 1.3 m , where the absolute error of the simplified innovative formula, $B_{\text {simp } 2}$, is practically the same as the error of the approximate formula: this is the price paid to the further level of approximation. However, in this range (and mostly elsewhere) the second level of approximation overestimates $B$, whereas eq. (2.17) underestimates $B$; hence - as the first level of approximation - also the second level provides a conservative estimate
with respect to the literature approximation.


Figure 2.8: $B_{\text {exact, }}, B_{\text {lit, }} B_{\text {simpl }}, B_{\text {simp } 2} v$ v. distance r from the axis of the helix for the cable ARG7HIRX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$ in overhead line configuration in linear coordinates.

Table 2.1: $B_{\text {exact, }} B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2} v s$. distance r from the axis of the helix and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARG7HIRX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$ in overhead line configuration.

| $\mathbf{r}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {it }}$ <br> $[\boldsymbol{\mu \mathbf { T } ]}$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu \mathbf { T }}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp }}$ <br> $[\boldsymbol{\mu \mathbf { T } ]}$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.3 | 22.92 | 13.49 | -41.15 | 25.34 | 10.56 | 25.17 | 9.81 |
| 0.4 | 10.94 | 7.385 | -32.50 | 10.84 | -0.94 | 10.87 | -0.66 |
| 0.5 | 5.695 | 4.175 | -26.68 | 5.553 | -2.49 | 5.598 | -1.70 |
| 0.6 | 3.112 | 2.410 | -22.57 | 3.066 | -1.47 | 3.103 | -0.30 |
| 0.7 | 1.752 | 1.410 | -19.53 | 1.756 | 0.21 | 1.782 | 1.71 |
| 0.8 | 1.007 | 0.8339 | -17.21 | 1.026 | 1.91 | 1.045 | 3.71 |
| 0.9 | 0.5873 | 0.4970 | -15.37 | 0.6071 | 3.38 | 0.6194 | 5.47 |
| 1.0 | 0.3461 | 0.2980 | -13.90 | 0.3619 | 4.55 | 0.3701 | 6.92 |
| 1.1 | 0.2057 | 0.1796 | -12.68 | 0.2168 | 5.39 | 0.2223 | 8.04 |
| 1.2 | 0.1231 | 0.1087 | -11.66 | 0.1304 | 5.92 | 0.1339 | 8.83 |
| 1.3 | 0.0740 | 0.0660 | -10.79 | 0.0786 | 6.16 | 0.0809 | 9.32 |
| 1.4 | 0.0447 | 0.0402 | -10.05 | 0.0475 | 6.13 | 0.0490 | 9.54 |
| 1.5 | 0.0271 | 0.0246 | -9.40 | 0.0287 | 5.85 | 0.0297 | 9.50 |
| 1.6 | 0.0165 | 0.0150 | -8.84 | 0.0174 | 5.36 | 0.0180 | 9.24 |
| 1.7 | 0.0101 | 0.0092 | -8.34 | 0.0105 | 4.67 | 0.0109 | 8.77 |
| 1.8 | 0.0062 | 0.0057 | -7.89 | 0.0064 | 3.82 | 0.0067 | 8.12 |
| 1.9 | 0.0038 | 0.0035 | -7.50 | 0.0039 | 2.82 | 0.0040 | 7.31 |
| 2.0 | 0.0023 | 0.0021 | -7.14 | 0.0024 | 1.67 | 0.0025 | 6.36 |

Dealing now with the considered cable used in underground power lines at a burial depth of 0.80 m , the values of $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ calculated at the soil level as a
function of the horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane are shown in Fig. 2.9 in linear coordinates. From the viewpoint of the performances of $B_{\text {simp } 1}$ and $B_{\text {simp } 2}$ vs. $B_{\text {lit }}$, considerations very similar to those made for the previous case hold in this case, too. Indeed, Fig. 2.9 shows clearly that the error of the simplified innovative formula, in both levels of approximation, is always much smaller than the error of the approximate formula from the literature. For very small values of $B$ the curves in Fig. 2.9 tend to overlap; thus, it is better to refer to the numerical values of $B_{\text {exact }}, B_{l i t}, B_{\text {simpl }}, B_{\text {simp } 2}$ and to the percent errors of $B_{\text {lit }}, B_{\text {simpl }}$, $B_{\text {simp } 2}$ with respect to $B_{\text {exact }}$ reported in Table 2.2. The Table confirms that the percent errors provided by innovative eq. (2.23) in first level of approximation remain much lower - never exceeding a few $\%$ even for small values of the distance $x$ and tending to $0-$ than those provided by eq. (2.17) after [14], [15] and always positive.As to the second level of approximation, the errors appear slightly larger, but never exceeding $10 \%$, and always positive, also providing a conservative estimate with respect to the literature approximation.


Figure 2.9: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpp }}, B_{\text {simp } 2}$ at ground level vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $\mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.2: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at ground level vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $k V 3 x 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m.

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 1.007 | 0.8339 | -17.21 | 1.026 | 1.91 | 1.045 | 3.71 |
| 0.1 | 0.9735 | 0.8073 | -17.08 | 0.9930 | 2.01 | 1.011 | 3.83 |
| 0.2 | 0.8809 | 0.7336 | -16.71 | 0.9011 | 2.29 | 0.9176 | 4.17 |
| 0.3 | 0.7499 | 0.6287 | -16.16 | 0.7704 | 2.74 | 0.7852 | 4.70 |
| 0.4 | 0.6050 | 0.5114 | -15.47 | 0.6250 | 3.31 | 0.6376 | 5.38 |
| 0.5 | 0.4663 | 0.3978 | -14.69 | 0.4846 | 3.93 | 0.4950 | 6.14 |
| 0.6 | 0.3461 | 0.2980 | -13.90 | 0.3619 | 4.55 | 0.3701 | 6.92 |
| 0.7 | 0.2492 | 0.2165 | -13.10 | 0.2619 | 5.12 | 0.2683 | 7.66 |
| 0.8 | 0.1750 | 0.1534 | -12.34 | 0.1848 | 5.59 | 0.1896 | 8.32 |
| 0.9 | 0.1205 | 0.1065 | -11.62 | 0.1276 | 5.94 | 0.1312 | 8.86 |
| 1.0 | 0.0817 | 0.0727 | -10.95 | 0.0867 | 6.14 | 0.0892 | 9.25 |
| 1.1 | 0.0546 | 0.0490 | -10.33 | 0.0580 | 6.17 | 0.0598 | 9.48 |
| 1.2 | 0.0362 | 0.0327 | -9.77 | 0.0384 | 6.04 | 0.0396 | 9.55 |
| 1.3 | 0.0238 | 0.0216 | -9.25 | 0.0251 | 5.74 | 0.0260 | 9.45 |
| 1.4 | 0.0155 | 0.0141 | -8.77 | 0.0163 | 5.29 | 0.0169 | 9.19 |
| 1.5 | 0.0101 | 0.0092 | -8.34 | 0.0105 | 4.67 | 0.0109 | 8.77 |

For the same underground configuration of Fig. 2.9, the values of $B_{\text {exact }}, B_{l i t}, B_{\text {simp }}$, $B_{\text {simp2 }}$ are calculated 1 m above the soil level, i.e. where the human vital organs are situated, as a function of the horizontal distance $x$ from the orthogonal projection of the cable axis on the plane parallel to the ground at the same height are shown in Fig. 2.10 in linear coordinates.


Figure 2.10: $B_{\text {exact }}, B_{l i t}, B_{\text {simpl }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARE4H1RX 12/20 $\mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}$ ( $I=280 \mathrm{~A}$ ) in underground line configuration at a depth of 0.80 m in linear coordinates.

The values of $B$ are very small and the curves in Fig. 2.10 tend to overlap, although the better performances of $B_{\text {simpl }}$ and $B_{\text {simp2 }}$, with respect to $B_{l i t}$ remain evident; however, for a less subjective evaluation, the numerical values of $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}$, $B_{\text {simp } 2}$ and the percent errors of $B_{l i t}, B_{\text {simp } 1}, B_{\text {simp } 2}$, with respect to $B_{\text {exact }}$ are reported in Table 2.3.

Table 2.3: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {it }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp }}$ <br> $[\boldsymbol{\mu \mathbf { T } ]}$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.0062 | 0.0057 | -7.89 | 0.0064 | 3.82 | 0.0067 | 8.12 |
| 0.1 | 0.0061 | 0.0056 | -7.88 | 0.0063 | 3.79 | 0.0066 | 8.10 |
| 0.2 | 0.0058 | 0.0054 | -7.85 | 0.0060 | 3.71 | 0.0063 | 8.04 |
| 0.3 | 0.0054 | 0.0050 | -7.79 | 0.0056 | 3.58 | 0.0059 | 7.93 |
| 0.4 | 0.0050 | 0.0046 | -7.71 | 0.0051 | 3.39 | 0.0053 | 7.78 |
| 0.5 | 0.0044 | 0.0041 | -7.62 | 0.0045 | 3.15 | 0.0047 | 7.58 |
| 0.6 | 0.0038 | 0.0035 | -7.51 | 0.0039 | 2.84 | 0.0041 | 7.33 |
| 0.7 | 0.0032 | 0.0030 | -7.38 | 0.0033 | 2.47 | 0.0035 | 7.03 |
| 0.8 | 0.0027 | 0.0025 | -7.24 | 0.0027 | 2.03 | 0.0029 | 6.66 |
| 0.9 | 0.0022 | 0.0020 | -7.09 | 0.0022 | 1.52 | 0.0023 | 6.23 |
| 1.0 | 0.0017 | 0.0016 | -6.94 | 0.0018 | 0.95 | 0.0018 | 5.74 |
| 1.1 | 0.0014 | 0.0013 | -6.78 | 0.0014 | 0.30 | 0.0014 | 5.17 |
| 1.2 | 0.0010 | 0.0010 | -6.62 | 0.0010 | -0.43 | 0.0011 | 4.54 |
| 1.3 | 0.0008 | 0.0007 | -6.46 | 0.0008 | -1.22 | 0.0008 | 3.84 |
| 1.4 | 0.0006 | 0.0006 | -6.30 | 0.0006 | -2.08 | 0.0006 | 3.07 |
| 1.5 | 0.0004 | 0.0004 | -6.14 | 0.0004 | -3.01 | 0.0004 | 2.23 |

All the comments relevant to previous Fig. 2.9 hold in the case of Fig. 2.10 too - in particular a further reduction of the field values is observed, as expected due to the further increase of the distance between field points and field source. Also in this case, the error of the simplified innovative formula, in both levels of approximation, is always smaller than the error of the approximate formula from the literature, being limited to a few $\%$ even for small values of the distance $x$ and tending to 0 , but now the difference between the simplified innovative formula and the approximate formula from the literature is smaller, due to the greater distance of the field-points from the helix axis. In particular, for these distances, $B_{\text {simp } 1}$ starts to underestimate the field providing a negative error. This is due to the fact that the approximation of the exact formula has been effectuated considering interval up to 2.0 m for the linear fitting of $\ln B_{2}$ vs. $r$. This means that after 2.0 m it is possible that the simplified formula presents different behaviour than that expected. This inconvenience is compensated by the fact that at distances higher than 2.0 m the magnetic field is so drastically reduced that results practically zero. For this reason it was chosen to approximate with as much higher precision and accuracy the magnetic field up to 2.0 m near the field source and not to focus on the effects of the field where results negligible.

The better performances of the proposed simplified innovative formula compared to the approximated formula after [14], [15] is supported by two more practical cases, based on the same type of cable ARG7H1RX, but with different conductor crosssections with respect to the previous one, hence with different values of $p$ and $\alpha$. These two cases are as follows:
$I^{\text {st }}$ case: conductor cross-section $3 \times 70 \mathrm{~mm}^{2}$ [63], rated voltage 12 (phase to ground)/20(phase-to-phase) kV , ampacity $I=200 \mathrm{~A}, p=1.24 \mathrm{~m}, \alpha=0.018 \mathrm{~m}$;
$2^{\text {nd }}$ case: conductor cross-section $3 \times 185 \mathrm{~mm}^{2}$ [63], rated voltage 12 (phase to ground) $/ 20$ (phase-to-phase) kV , ampacity $I=360 \mathrm{~A}, p=1.5 \mathrm{~m}, \alpha=0.022 \mathrm{~m}$.

As shown in Figs. 2.11-2.16 and Tables 2.4-2.9, such cases confirm that the error involved by the innovative formula with respect to the exact formulation is not only much lower than that brought about by the approximated formula from the literature, but also fairly small (say, below $10 \%$ for distances $>0.3 \mathrm{~m}$ and mostly around a few percent) and positive, thereby implying a conservative overestimation of the field. Since - as previously highlighted - the exact formula was extensively validated experimentally in [14], [15], this proves the satisfactory behaviour of the proposed innovative simplified formula not only in comparison with the approximated formula of the literature, but also in absolute terms.


Figure 2.11: $B_{\text {exact, }} B_{\text {lit, }} B_{\text {simpl }}, B_{\text {simp } 2}$ vs. distance r from the axis of the helix for the cable ARG7HIRX $12 / 20 \mathrm{kV} \mathrm{3x70} \mathrm{~mm}{ }^{2}(I=200 \mathrm{~A})$ in overhead line configuration in linear coordinates.

Table 2.4: $B_{\text {exact, }} B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ vs. distance $r$ from the axis of the helix and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARG7H1RX $12 / 20 \mathrm{kV} 3 x 70 \mathrm{~mm}^{2}(I=200 \mathrm{~A})$ in overhead line configuration.

| $\mathbf{r}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.3 | 14.06 | 8.718 | -38.02 | 15.45 | 9.86 | 15.39 | 9.43 |
| 0.4 | 6.478 | 4.548 | -29.79 | 6.443 | -0.54 | 6.450 | -0.44 |
| 0.5 | 3.240 | 2.451 | -24.36 | 3.178 | -1.93 | 3.190 | -1.55 |
| 0.6 | 1.697 | 1.348 | -20.55 | 1.681 | -0.95 | 1.691 | -0.35 |
| 0.7 | 0.9143 | 0.7519 | -17.77 | 0.9202 | 0.64 | 0.9275 | 1.44 |
| 0.8 | 0.5023 | 0.4237 | -15.64 | 0.5136 | 2.25 | 0.5185 | 3.22 |
| 0.9 | 0.2798 | 0.2407 | -13.97 | 0.2900 | 3.64 | 0.2932 | 4.80 |
| 1.0 | 0.1574 | 0.1376 | -12.62 | 0.1649 | 4.74 | 0.1670 | 6.07 |
| 1.1 | 0.0893 | 0.0790 | -11.51 | 0.0943 | 5.54 | 0.0956 | 7.04 |
| 1.2 | 0.0510 | 0.0456 | -10.59 | 0.0541 | 6.03 | 0.0549 | 7.70 |
| 1.3 | 0.0293 | 0.0264 | -9.80 | 0.0311 | 6.24 | 0.0316 | 8.07 |
| 1.4 | 0.0169 | 0.0153 | -9.13 | 0.0179 | 6.19 | 0.0182 | 8.18 |
| 1.5 | 0.0097 | 0.0089 | -8.54 | 0.0103 | 5.91 | 0.0105 | 8.05 |
| 1.6 | 0.0057 | 0.0052 | -8.03 | 0.0060 | 5.42 | 0.0061 | 7.71 |
| 1.7 | 0.0033 | 0.0030 | -7.58 | 0.0034 | 4.74 | 0.0035 | 7.18 |
| 1.8 | 0.0019 | 0.0018 | -7.17 | 0.0020 | 3.90 | 0.0020 | 6.47 |
| 1.9 | 0.0011 | 0.0010 | -6.81 | 0.0012 | 2.91 | 0.0012 | 5.61 |
| 2.0 | 0.0007 | 0.0006 | -6.49 | 0.0007 | 1.79 | 0.0007 | 4.61 |



Figure 2.12: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simp } 1}, B_{\text {simp } 2}$ at ground level vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $k V 3 x 70 \mathrm{~mm}^{2}(I=200 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.5: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at ground level vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $k V 3 x 70 \mathrm{~mm}^{2}(I=200 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\boldsymbol{\%}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.5023 | 0.4237 | -15.64 | 0.5136 | 2.25 | 0.5185 | 3.22 |
| 0.1 | 0.4842 | 0.4090 | -15.52 | 0.4955 | 2.34 | 0.5003 | 3.33 |
| 0.2 | 0.4344 | 0.3684 | -15.19 | 0.4458 | 2.61 | 0.4502 | 3.64 |
| 0.3 | 0.3648 | 0.3112 | -14.68 | 0.3759 | 3.04 | 0.3798 | 4.11 |
| 0.4 | 0.2890 | 0.2484 | -14.05 | 0.2993 | 3.57 | 0.3026 | 4.72 |
| 0.5 | 0.2178 | 0.1887 | -13.35 | 0.2268 | 4.16 | 0.2295 | 5.39 |
| 0.6 | 0.1574 | 0.1376 | -12.62 | 0.1649 | 4.74 | 0.1670 | 6.07 |
| 0.7 | 0.1101 | 0.0970 | -11.90 | 0.1159 | 5.28 | 0.1174 | 6.72 |
| 0.8 | 0.0749 | 0.0665 | -11.21 | 0.0791 | 5.72 | 0.0803 | 7.28 |
| 0.9 | 0.0498 | 0.0446 | -10.55 | 0.0528 | 6.04 | 0.0537 | 7.72 |
| 1.0 | 0.0326 | 0.0293 | -9.95 | 0.0346 | 6.22 | 0.0352 | 8.02 |
| 1.1 | 0.0210 | 0.0190 | -9.38 | 0.0223 | 6.24 | 0.0227 | 8.17 |
| 1.2 | 0.0134 | 0.0122 | -8.87 | 0.0142 | 6.10 | 0.0145 | 8.16 |
| 1.3 | 0.0084 | 0.0077 | -8.40 | 0.0089 | 5.80 | 0.0091 | 7.98 |
| 1.4 | 0.0053 | 0.0049 | -7.97 | 0.0056 | 5.34 | 0.0057 | 7.65 |
| 1.5 | 0.0033 | 0.0030 | -7.58 | 0.0034 | 4.74 | 0.0035 | 7.18 |



Figure 2.13: $B_{\text {exact, }}, B_{\text {lit }}, B_{\text {simp } 1}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 70 \mathrm{~mm}^{2}$ ( $I=200 \mathrm{~A}$ ) in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.6: $B_{\text {exact, }}, B_{\text {lit }}, B_{\text {simp }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 70 \mathrm{~mm}^{2}(I=200 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\boldsymbol{\%}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.0019 | 0.0018 | -7.17 | 0.0020 | 3.90 | 0.0020 | 6.47 |
| 0.1 | 0.0019 | 0.0018 | -7.16 | 0.0020 | 3.87 | 0.0020 | 6.45 |
| 0.2 | 0.0018 | 0.0017 | -7.13 | 0.0019 | 3.80 | 0.0019 | 6.38 |
| 0.3 | 0.0017 | 0.0016 | -7.08 | 0.0017 | 3.67 | 0.0018 | 6.27 |
| 0.4 | 0.0015 | 0.0014 | -7.01 | 0.0016 | 3.48 | 0.0016 | 6.11 |
| 0.5 | 0.0013 | 0.0012 | -6.92 | 0.0014 | 3.24 | 0.0014 | 5.90 |
| 0.6 | 0.0011 | 0.0011 | -6.82 | 0.0012 | 2.94 | 0.0012 | 5.63 |
| 0.7 | 0.0009 | 0.0009 | -6.71 | 0.0010 | 2.57 | 0.0010 | 5.31 |
| 0.8 | 0.0008 | 0.0007 | -6.58 | 0.0008 | 2.14 | 0.0008 | 4.93 |
| 0.9 | 0.0006 | 0.0006 | -6.45 | 0.0006 | 1.65 | 0.0006 | 4.48 |
| 1.0 | 0.0005 | 0.0004 | -6.31 | 0.0005 | 1.08 | 0.0005 | 3.97 |
| 1.1 | 0.0004 | 0.0003 | -6.17 | 0.0004 | 0.44 | 0.0004 | 3.39 |
| 1.2 | 0.0003 | 0.0003 | -6.02 | 0.0003 | -0.27 | 0.0003 | 2.74 |
| 1.3 | 0.0002 | 0.0002 | -5.87 | 0.0002 | -1.04 | 0.0002 | 2.03 |
| 1.4 | 0.0001 | 0.0001 | -5.73 | 0.0001 | -1.89 | 0.0001 | 1.24 |
| 1.5 | 0.0001 | 0.0001 | -5.58 | 0.0001 | -2.80 | 0.0001 | 0.39 |



Figure 2.14: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp2 }}$ vs. distance $r$ from the axis of the helix for the cable ARG7H1RX 12/20 kV 3x185 $\mathrm{mm}^{2}(I=360 \mathrm{~A})$ in overhead line configuration in linear coordinates.

Table 2.7: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ vs. distance $r$ from the axis of the helix and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARG7H1RX $12 / 20 \mathrm{kV} 3 x 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$ in overhead line configuration.

| $\mathbf{r}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu \mathbf { T } ]}$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.3 | 33.56 | 18.76 | -44.09 | 37.03 | 10.36 | 36.85 | 9.82 |
| 0.4 | 16.47 | 10.69 | -35.10 | 16.18 | -1.73 | 16.22 | -1.52 |
| 0.5 | 8.849 | 6.295 | -28.94 | 8.560 | -3.26 | 8.610 | -2.70 |
| 0.6 | 5.004 | 3.776 | -24.54 | 4.898 | -2.10 | 4.939 | -1.30 |
| 0.7 | 2.921 | 2.299 | -21.27 | 2.912 | -0.28 | 2.942 | 0.73 |
| 0.8 | 1.741 | 1.415 | -18.76 | 1.768 | 1.53 | 1.789 | 2.73 |
| 0.9 | 1.054 | 0.8774 | -16.77 | 1.087 | 3.11 | 1.101 | 4.48 |
| 1.0 | 0.6453 | 0.5475 | -15.16 | 0.6735 | 4.36 | 0.6834 | 5.89 |
| 1.1 | 0.3985 | 0.3434 | -13.84 | 0.4195 | 5.27 | 0.4262 | 6.96 |
| 1.2 | 0.2478 | 0.2163 | -12.72 | 0.2623 | 5.84 | 0.2668 | 7.69 |
| 1.3 | 0.1549 | 0.1367 | -11.78 | 0.1644 | 6.11 | 0.1675 | 8.10 |
| 1.4 | 0.0973 | 0.0866 | -10.97 | 0.1032 | 6.09 | 0.1053 | 8.22 |
| 1.5 | 0.0613 | 0.0550 | -10.26 | 0.0649 | 5.83 | 0.0663 | 8.09 |
| 1.6 | 0.0388 | 0.0351 | -9.65 | 0.0409 | 5.33 | 0.0418 | 7.72 |
| 1.7 | 0.0246 | 0.0224 | -9.10 | 0.0258 | 4.64 | 0.0264 | 7.15 |
| 1.8 | 0.0157 | 0.0143 | -8.61 | 0.0162 | 3.77 | 0.0167 | 6.40 |
| 1.9 | 0.0100 | 0.0092 | -8.18 | 0.0102 | 2.74 | 0.0105 | 5.48 |
| 2.0 | 0.0064 | 0.0059 | -7.79 | 0.0065 | 1.58 | 0.0066 | 4.42 |



Figure 2.15: $B_{\text {exact }}, B_{\text {lit, }} B_{\text {simpl }}, B_{\text {simp2 }}$ at ground level vs. horizontal distance x from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $k V 3 \times 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.8: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at ground level vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7H1RX 12/20 $k V 3 \times 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m.

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\boldsymbol{\%}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 1.741 | 1.415 | -18.76 | 1.768 | 1.53 | 1.789 | 2.73 |
| 0.1 | 1.687 | 1.373 | -18.62 | 1.715 | 1.64 | 1.735 | 2.85 |
| 0.2 | 1.537 | 1.257 | -18.22 | 1.567 | 1.95 | 1.586 | 3.19 |
| 0.3 | 1.323 | 1.090 | -17.62 | 1.355 | 2.43 | 1.372 | 3.72 |
| 0.4 | 1.084 | 0.9009 | -16.87 | 1.117 | 3.03 | 1.131 | 4.39 |
| 0.5 | 0.8509 | 0.7145 | -16.03 | 0.8823 | 3.69 | 0.8946 | 5.14 |
| 0.6 | 0.6453 | 0.5475 | -15.16 | 0.6735 | 4.36 | 0.6834 | 5.89 |
| 0.7 | 0.4759 | 0.4078 | -14.30 | 0.4995 | 4.97 | 0.5073 | 6.60 |
| 0.8 | 0.3431 | 0.2969 | -13.47 | 0.3619 | 5.48 | 0.3679 | 7.22 |
| 0.9 | 0.2430 | 0.2122 | -12.68 | 0.2572 | 5.86 | 0.2617 | 7.71 |
| 1.0 | 0.1696 | 0.1493 | -11.95 | 0.1799 | 6.08 | 0.1833 | 8.04 |
| 1.1 | 0.1171 | 0.1039 | -11.28 | 0.1242 | 6.13 | 0.1267 | 8.21 |
| 1.2 | 0.0800 | 0.0715 | -10.66 | 0.0849 | 6.01 | 0.0866 | 8.20 |
| 1.3 | 0.0543 | 0.0489 | -10.09 | 0.0574 | 5.72 | 0.0587 | 8.01 |
| 1.4 | 0.0367 | 0.0332 | -9.57 | 0.0386 | 5.26 | 0.0395 | 7.66 |
| 1.5 | 0.0246 | 0.0224 | -9.10 | 0.0258 | 4.64 | 0.0264 | 7.15 |



Figure 2.16: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simp1 }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$ from the orthogonal projection of the cable axis on the ground plane for the cable ARG7HIRX $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.9: $B_{\text {exact, }} B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance x from the orthogonal projection of the cable axis on the ground plane for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.0157 | 0.0143 | -8.61 | 0.0162 | 3.77 | 0.0167 | 6.40 |
| 0.1 | 0.0155 | 0.0141 | -8.60 | 0.0160 | 3.74 | 0.0164 | 6.37 |
| 0.2 | 0.0149 | 0.0136 | -8.56 | 0.0154 | 3.66 | 0.0158 | 6.30 |
| 0.3 | 0.0140 | 0.0128 | -8.50 | 0.0145 | 3.53 | 0.0149 | 6.18 |
| 0.4 | 0.0128 | 0.0118 | -8.42 | 0.0133 | 3.34 | 0.0136 | 6.01 |
| 0.5 | 0.0115 | 0.0106 | -8.31 | 0.0119 | 3.09 | 0.0122 | 5.79 |
| 0.6 | 0.0101 | 0.0093 | -8.19 | 0.0104 | 2.77 | 0.0106 | 5.51 |
| 0.7 | 0.0087 | 0.0080 | -8.05 | 0.0089 | 2.39 | 0.0091 | 5.16 |
| 0.8 | 0.0073 | 0.0067 | -7.90 | 0.0074 | 1.95 | 0.0076 | 4.76 |
| 0.9 | 0.0060 | 0.0056 | -7.74 | 0.0061 | 1.43 | 0.0063 | 4.28 |
| 1.0 | 0.0049 | 0.0045 | -7.57 | 0.0049 | 0.84 | 0.0051 | 3.74 |
| 1.1 | 0.0039 | 0.0036 | -7.40 | 0.0039 | 0.18 | 0.0040 | 3.12 |
| 1.2 | 0.0031 | 0.0028 | -7.22 | 0.0031 | -0.56 | 0.0031 | 2.43 |
| 1.3 | 0.0024 | 0.0022 | -7.05 | 0.0023 | -1.37 | 0.0024 | 1.67 |
| 1.4 | 0.0018 | 0.0017 | -6.87 | 0.0018 | -2.25 | 0.0018 | 0.84 |
| 1.5 | 0.0014 | 0.0013 | -6.69 | 0.0013 | -3.20 | 0.0014 | -0.06 |

The effectiveness of the innovative expression is also evaluated with some additional numerical simulations relevant to another typical MV cable, named ARE4H1RX after [64] with the following main characteristics: conductor crosssection $3 \times 120 \mathrm{~mm}^{2}$, rated voltage 12 (phase to ground) $/ 20$ (phase-to-phase) kV , ampacity $I=288 \mathrm{~A}, p=1.30 \mathrm{~m}$ and $\alpha=0.019 \mathrm{~m}$.

As to the considered cable used in overhead power lines, the values of $B_{\text {exact }}, B_{\text {lit }}$, $B_{\text {simpl }}, B_{\text {simp } 2}$ as a function of the distance $r$ from the axis of the helix are shown in Fig. 2.17 in linear coordinates and the numerical values of $B_{\text {exact, }} B_{l i t}, B_{\text {simp }}, B_{\text {simp } 2}$ and the percent errors of $B_{l i t}, B_{\text {simpl }}, B_{\text {simp } 2}$ with respect to $B_{\text {exact }}$ are reported in Table 2.10.

Figure 2.18 deals with the considered cable used in underground power lines at a burial depth of 0.80 m presenting the values of $B_{\text {exact, }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ calculated at the soil level as a function of the distance $x$ from the axis of the cable in linear coordinates and Table 2.11 reports the numerical values of $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simp }}, B_{\text {simp } 2}$ and the percent errors of $B_{l i t}, B_{\text {simpl }}, B_{\text {simp } 2}$ with respect to $B_{\text {exact }}$.

Finally, Fig. 2.19 shows the values of $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ calculated 1 m above the soil level as a function of the distance $x$ from the cable axis in linear for the same underground configuration of Fig. 2.18, and in Table 2.12 are reported the percent errors of $B_{\text {simpl }}, B_{\text {simp } 2}, B_{l i t}$ with respect to $B_{\text {exact }}$.


Figure 2.17: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ vs. distance $r$ from the axis of the helix for the cable ARE4HIRX $12 / 20 \mathrm{kV} \mathrm{3x120} \mathrm{~mm}{ }^{2}(I=288 \mathrm{~A})$ in overhead line configuration in linear coordinates.

Table 2.10: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simp } 1}, B_{\text {simp } 2}$ vs. distance $r$ from the axis of the helix and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=288 \mathrm{~A})$ in overhead line configuration.

| $\mathbf{r}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.3 | 21.88 | 13.24 | -39.49 | 24.14 | 10.32 | 23.99 | 9.65 |
| 0.4 | 10.26 | 7.072 | -31.05 | 10.19 | -0.67 | 10.21 | -0.46 |
| 0.5 | 5.232 | 3.901 | -25.44 | 5.119 | -2.16 | 5.153 | -1.51 |
| 0.6 | 2.798 | 2.196 | -21.49 | 2.765 | -1.17 | 2.792 | -0.19 |
| 0.7 | 1.540 | 1.254 | -18.58 | 1.547 | 0.45 | 1.567 | 1.73 |
| 0.8 | 0.8650 | 0.7235 | -16.37 | 0.8831 | 2.09 | 0.8965 | 3.64 |
| 0.9 | 0.4927 | 0.4207 | -14.62 | 0.5100 | 3.52 | 0.5189 | 5.33 |
| 1.0 | 0.2836 | 0.2461 | -13.21 | 0.2968 | 4.65 | 0.3026 | 6.71 |
| 1.1 | 0.1646 | 0.1447 | -12.05 | 0.1736 | 5.47 | 0.1774 | 7.77 |
| 1.2 | 0.0961 | 0.0855 | -11.09 | 0.1019 | 5.98 | 0.1043 | 8.52 |
| 1.3 | 0.0564 | 0.0506 | -10.26 | 0.0599 | 6.20 | 0.0615 | 8.98 |
| 1.4 | 0.0333 | 0.0301 | -9.56 | 0.0353 | 6.16 | 0.0363 | 9.16 |
| 1.5 | 0.0197 | 0.0179 | -8.94 | 0.0208 | 5.88 | 0.0215 | 9.10 |
| 1.6 | 0.0117 | 0.0107 | -8.41 | 0.0123 | 5.39 | 0.0127 | 8.81 |
| 1.7 | 0.0070 | 0.0064 | -7.93 | 0.0073 | 4.70 | 0.0075 | 8.33 |
| 1.8 | 0.0042 | 0.0038 | -7.51 | 0.0043 | 3.85 | 0.0045 | 7.67 |
| 1.9 | 0.0025 | 0.0023 | -7.13 | 0.0026 | 2.86 | 0.0027 | 6.86 |
| 2.0 | 0.0015 | 0.0014 | -6.79 | 0.0015 | 1.73 | 0.0016 | 5.91 |



Figure 2.18: $B_{\text {exact, }} B_{\text {lit, }} B_{\text {simpl }}, B_{\text {simp2 }}$ at ground level vs. horizontal distance x from the orthogonal projection of the cable axis on the ground plane for the cable ARE4H1RX 12/20 $k V 3 \times 120 \mathrm{~mm}^{2}(I=288 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.11: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at ground level vs. distance $x$ from the cable axis and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARE4HIRX 12/20 kV 3x120 $m^{2}(I=288 ~ A)$ in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {simp } 2}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.8650 | 0.7235 | -16.37 | 0.8831 | 2.09 | 0.8965 | 3.64 |
| 0.1 | 0.8349 | 0.6993 | -16.25 | 0.8532 | 2.19 | 0.8663 | 3.75 |
| 0.2 | 0.7522 | 0.6326 | -15.90 | 0.7708 | 2.47 | 0.7829 | 4.08 |
| 0.3 | 0.6359 | 0.5382 | -15.37 | 0.6544 | 2.90 | 0.6651 | 4.59 |
| 0.4 | 0.5082 | 0.4335 | -14.71 | 0.5258 | 3.45 | 0.5349 | 5.24 |
| 0.5 | 0.3872 | 0.3331 | -13.97 | 0.4029 | 4.05 | 0.4104 | 5.97 |
| 0.6 | 0.2836 | 0.2461 | -13.21 | 0.2968 | 4.65 | 0.3026 | 6.71 |
| 0.7 | 0.2011 | 0.1760 | -12.46 | 0.2116 | 5.20 | 0.2160 | 7.42 |
| 0.8 | 0.1389 | 0.1226 | -11.73 | 0.1468 | 5.66 | 0.1501 | 8.04 |
| 0.9 | 0.0940 | 0.0836 | -11.05 | 0.0996 | 5.99 | 0.1020 | 8.55 |
| 1.0 | 0.0625 | 0.0560 | -10.41 | 0.0664 | 6.18 | 0.0681 | 8.91 |
| 1.1 | 0.0411 | 0.0370 | -9.83 | 0.0436 | 6.20 | 0.0448 | 9.12 |
| 1.2 | 0.0267 | 0.0242 | -9.29 | 0.0283 | 6.07 | 0.0291 | 9.16 |
| 1.3 | 0.0172 | 0.0156 | -8.79 | 0.0181 | 5.77 | 0.0187 | 9.04 |
| 1.4 | 0.0110 | 0.0100 | -8.34 | 0.0115 | 5.31 | 0.0119 | 8.77 |
| 1.5 | 0.0070 | 0.0064 | -7.93 | 0.0073 | 4.70 | 0.0075 | 8.33 |



Figure 2.19: $B_{\text {exact }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. horizontal distance $x$
from the orthogonal projection of the cable axis on the ground plane for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=288 \mathrm{~A})$ in underground line configuration at a depth of 0.80 m in linear coordinates.

Table 2.12: $B_{\text {exact, }}, B_{\text {lit }}, B_{\text {simpl }}, B_{\text {simp } 2}$ at 1 m height from the ground vs. distance $x$ from the cable axis and relevant percent errors with respect to $B_{\text {exact }}$ for the cable ARE4H1RX 12/20 kV $3 \times 120 \mathrm{~mm}^{2}(I=288$ A) in underground line configuration at a depth of 0.80 m .

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {lit }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\mathbf{\%}]$ | $\mathbf{B}_{\text {simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\mathbf{\%}]$ | $\mathbf{B}_{\text {simp2 }}$ <br> $[\boldsymbol{\mu \mathbf { T } ]}$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 0.0042 | 0.0038 | -7.51 | 0.0043 | 3.85 | 0.0045 | 7.67 |
| 0.1 | 0.0041 | 0.0038 | -7.50 | 0.0042 | 3.83 | 0.0044 | 7.65 |
| 0.2 | 0.0039 | 0.0036 | -7.46 | 0.0041 | 3.75 | 0.0042 | 7.59 |
| 0.3 | 0.0037 | 0.0034 | -7.41 | 0.0038 | 3.62 | 0.0039 | 7.49 |
| 0.4 | 0.0033 | 0.0031 | -7.34 | 0.0034 | 3.43 | 0.0036 | 7.33 |
| 0.5 | 0.0029 | 0.0027 | -7.25 | 0.0030 | 3.19 | 0.0031 | 7.13 |
| 0.6 | 0.0025 | 0.0023 | -7.14 | 0.0026 | 2.89 | 0.0027 | 6.88 |
| 0.7 | 0.0021 | 0.0020 | -7.02 | 0.0022 | 2.52 | 0.0023 | 6.57 |
| 0.8 | 0.0017 | 0.0016 | -6.89 | 0.0018 | 2.08 | 0.0018 | 6.21 |
| 0.9 | 0.0014 | 0.0013 | -6.75 | 0.0014 | 1.58 | 0.0015 | 5.78 |
| 1.0 | 0.0011 | 0.0010 | -6.60 | 0.0011 | 1.01 | 0.0012 | 5.28 |
| 1.1 | 0.0008 | 0.0008 | -6.45 | 0.0009 | 0.37 | 0.0009 | 4.72 |
| 1.2 | 0.0006 | 0.0006 | -6.30 | 0.0006 | -0.35 | 0.0007 | 4.09 |
| 1.3 | 0.0005 | 0.0005 | -6.15 | 0.0005 | -1.13 | 0.0005 | 3.39 |
| 1.4 | 0.0004 | 0.0003 | -5.99 | 0.0003 | -1.98 | 0.0004 | 2.62 |
| 1.5 | 0.0003 | 0.0002 | -5.84 | 0.0003 | -2.91 | 0.0003 | 1.79 |

Considerations very similar to those made for the cable ARG7H1RX $12 / 20 \mathrm{kV}$ with conductor cross-section $3 \times 120 \mathrm{~mm}^{2}$ also hold in these cases. Figs. 2.17-2.19 and Tables $2.10-2.12$ prove once more that the error introduced by the simplified innovative formula, in the first and second level of approximation, is modest ( $<10 \%$ ) and mostly positive, providing a conservative estimate of the field. On the other hand, the literature approximation (eq. (2.17)) brings about an error always higher, in absolute terms, than that brought about by the proposed simplified formula, and always negative, underestimating the field.

The simulations results also include the calculation of the magnetic field isolines of $1.0 \mu \mathrm{~T}, 3.0 \mu \mathrm{~T}$ ('quality objective'), $5 \mu \mathrm{~T}, 10 \mu \mathrm{~T}$ ('attention value') and $100 \mu \mathrm{~T}$ ('exposure limit') for frequently used MV cable types, and are presented in Appendix 2 , facilitating the direct calculation of the distances corresponding to a maximum limit value of the rms magnetic induction generated by these cables carrying current equal at their ampacity.

### 2.5 CONCLUSIONS

As shown in this Chapter, the literature approximation for the calculation of the magnetic field generated by a twisted three-phase power cable carrying balanced three-phase currents provides acceptable results for large distances from the helix axis only, whereas for distances close to the conductor it exhibits large errors compared to the exact expression of the magnetic field. For this reason, a parametric heuristic analysis was performed that resulted in an innovative simplified expression of the rms magnetic induction as a function of the distance from the helix axis. This innovative expression approximates the logarithm of $B$ with the equation of a straight line plus a hyperbolic term for simulating the deviation from linearity close to the helix axis.

The effectiveness of the innovative expression is evaluated through some numerical simulations relevant to two typical MV cables i.e. ARG7H1RX $12 / 20 \mathrm{kV}$ with cross-sections $3 \times 70 \mathrm{~mm}^{2}, 3 \times 120 \mathrm{~mm}^{2}$ and $3 \times 185 \mathrm{~mm}^{2}$, and ARE4H1RX with cross-section $3 \times 120 \mathrm{~mm}^{2}$, in order to compare the results provided by the exact and the approximate formulae from the literature with those obtained via the innovative simplified formula. The very similar results obtained from the simulations concerning the two different MV cable types have proven the satisfactory behaviour of the proposed simplified formula in comparison with the approximated formula from the literature not only for the cases treated here, but also in broader terms. Thus, it can be argued that the considerations made about the proposed simplified formula in both levels of approximation, can be regarded as valid for all cable types commonly employed in distribution networks.

In conclusion, it can be said that the innovative formula results definitely much simpler than the exact one and provides a much smaller relative error compared to the approximated one from the literature, especially for small distances from the helix axis.

# Double-Circuit Twisted Three-Phase Cable Lines 

### 3.1 THEORY

Here, both the exact and the simplified innovative formulation are extended and applied to double-circuit twisted cable lines, making the hypothesis that the currents carried by both the cables are in-phase. Because of the linearity of the medium, the study of the magnetic field from adouble-circuit twisted three-phase cable line is based on the superposition principle, which implies that the magnetic field generated by each circuit can be calculated separately at every field-point and then the two individual magnetic fields can be composed linearly. As far as this composition is concerned, different approaches of increasing complexity and accuracy will be followed.

### 3.1.1 Exact Vector Analysis

The exact approach to the study of the magnetic field from a double-circuit twisted three-phase cable line is based, firstly, on the conversion between the cylindrical coordinate system (see Fig. 2.2) and the Cartesian coordinate system and, subsequently, on carrying out the vector sum of the Cartesian components of the magnetic field in a "reference plane" where the field points of interest lie. This plane commonly coincides with a line section taken as orthogonal to the ground, the latter assumed as flat. For overhead cable lines - where cables hang from one pole to another at variable distances from the ground in a catenary-like shape - this section is usually the mid-span section, where the conductors are closest to the ground and hence the field is highest at the ground level. For underground cable lines, where cables in general run parallel to each other and to the ground, this section can be a particular section close to sensitive receptors.

Hence, it can be concluded that both for overhead and underground double-circuit twisted three-phase cable lines the "reference plane" for calculation is a Cartesian $x-y$ plane assumed as orthogonal to the helix axis of both twisted three-phase cables and as matching the equation $z=0$. The helix axis - or cylindrical axis z of Fig. 2.2 - is the Cartesian $z$-axis; hence, the $z$ coordinate is the same in both systems.

In order to address better the "exact" vector sum of the Cartesian components of the magnetic field for double-circuit twisted three-phase cable lines, a "reference arrangement" of the double-circuit line is firstly considered, wherethe two twisted
cable circuits are attached on opposite sides of poles and placed at the same height from the ground, being $d$ the horizontal distance between the two cables: for the sake ofclarity, "cable number 1 " is the twisted cable located in the semi-space where $x<0$, while "cable number 2 " is that located in the semi-space where $x>0$.

Fig. 3.1 shows a sketch of the chosen "reference arrangement", as well as the radial field component, $B_{r}$, and the azimuthal field component, $B_{\varphi}$, generated by circuit 1 and 2 at five different equally spaced field-points lying along the $x$-axis and contained in the plane $z=0$. As a consequence of the chosen arrangement, the radial component generated by each cable at a certain field point has the direction of the straight line connecting the helix axis of each cable with that field point, whereas the azimuthal component generated by each cable is orthogonal to the radial component at every field-point. Note that the axial component, $B_{z}$, at each field-point is omitted in Fig. 3.1 since the $z$-axis is orthogonal to the plane $z=0$, as outlined in the Figure.


Figure 3.1: Sketch of the "reference arrangement" of the double-circuittwisted three-phase cable line chosen as first case-study for the vector sum of the Cartesian components of the magnetic field. The relevant radial ( $B_{r}$ ) and azimuthal $\left(B_{\varphi}\right)$ field components generated by circuit 1 (green arrows) and 2 (red arrows) at five equally spaced field-points along the $x$ axis are shown.

The magnetic field from the double-circuit, $B_{\text {DOUBLE }}$, is given by the following formula:

$$
\begin{equation*}
B_{D O U B L E}=\sqrt{B_{x}^{2}+B_{y}^{2}+B_{z}^{2}} \tag{3.1}
\end{equation*}
$$

being $\left\{\begin{array}{l}B_{x}=B_{r 1} \cos \theta_{1}+B_{r 2} \cos \theta_{2}+B_{\phi 1} \cos \varphi_{1}+B_{\phi 2} \cos \varphi_{2} \\ B_{y}=B_{r 1} \sin \theta_{1}+B_{r 2} \sin \theta_{2}+B_{\phi 1} \sin \varphi_{1}+B_{\phi 2} \sin \varphi_{2} \\ B_{z}=B_{z 1}+B_{z 2}\end{array}\right.$
where $B_{r 1}, B_{r 2}, B_{\varphi 1}, B_{\varphi 2}$ and $B_{z 1}, B_{z 2}$ are the exact values of the radial, azimuthal and axial components of the magnetic field from circuit 1 and from circuit 2, respectively, while $\theta_{1}, \theta_{2}$, and $\varphi_{1}, \varphi_{2}$ are the angles formed by the $x$-axis and the radial, $B_{r}$, and azimuthal components, $B_{\varphi}$, from circuit 1 and 2 , respectively (measured anticlockwise from the $x$-axis).

It holds:

$$
\left\{\begin{array}{l}
\theta_{1}=\arctan \left(y_{1} / x_{1}\right)  \tag{3.3}\\
\theta_{2}=\arctan \left(y_{2} / x_{2}\right)
\end{array}\right.
$$

and

$$
\left\{\begin{array}{l}
\varphi_{1}=\theta_{1}+\pi / 2  \tag{3.4}\\
\varphi_{2}=\theta_{2}+\pi / 2
\end{array}\right.
$$

where $\left(x_{1}, y_{l}\right)$ are the coordinates of cable 1 and $\left(x_{2}, y_{2}\right)$ are the coordinates of cable 2 in the plane $z=0$, where $\varphi=\Phi$.

The above theoretical approach has been converted into a calculation tool for the magnetic field from double-circuit twisted three-phase cable lines by implementing relationships (3.1) - (3.4), plus exact formulae (2.6) for $B_{r 1}, B_{r 2}, B_{\varphi 1}, B_{\varphi 2}$ and $B_{z 1}, B_{z 2}$, respectively, in a script in Matlab ${ }^{\mathrm{TM}}$ environment.

A proper approach to some practical problems involved by the vector sum can be accomplished by focusing more into the "reference arrangement" sketched in Fig. 3.1. So, for this first case study, the two twisted cable circuits displayed in Fig. 3.1 are ARG7H1RX [63] cables with cross-section $3 \times 120 \mathrm{~mm}^{2}$, rated voltage 12 (phase-toground) $/ 20$ (phase-to-phase) kV , ampacity $I=280 \mathrm{~A}$, pitch $p=1.37 \mathrm{~m}$, radius $\alpha=0.020 \mathrm{~m}$ and the distance between them has been chosen to be $d=1.0 \mathrm{~m}$. Further, five equallyspaced field points on a straight horizontal line that spreads from -1.0 m to 1.0 m at both ends of the line axis and whose vertical distance from the conductors is $H=0.5 \mathrm{~m}$ have been set.

### 3.1.1.1 Angle $\Phi$

As clearly shown by relationships (2.6) and (2.12), the total field and the relevant components are also functions of the angle $\Phi$ which indicates the angular position of the field-point in the twisted configuration (see Fig. 2.2). Therefore, performing the exact vector sum at each $x-y$ plane of interest for the calculation of magnetic field requires knowing the precise geometrical arrangement of the three phases of each twisted cable at that plane, but it can be argued that this geometrical arrangement - in
particular the angle $\Phi$ of the three phases of each circuit - can be hardly known precisely in practice.

Because of the above-mentioned practical problems, 16 possible phase arrangements have been chosen as "base arrangements", from which 8 are symmetric and 8 are anti-symmetric phase configurations, as shown in Tables 3.1 and 3.2. the angles $\theta_{l}$ and $\theta_{2}$, reported in Tables 3.1 and 3.2, are the angles of eq. 3.3.

Table 3.1: Geometrical arrangements of the 8 symmetric "base arrangements" with relevant angles $\Phi_{1}$ and $\Phi_{2}$.
SMMMETRIC PHASE ARRANGEMENTS

Table 3.2: Geometrical arrangements of the 8 anti-symmetric "base arrangements" with relevant angles $\Phi_{1}$ and $\Phi_{2}$.
(1.b)

For each base arrangement, the calculation of the total magnetic field was effectuated in function of the angles $\Phi_{1}$ and $\Phi_{2}$, which are not constant along the $x$ axis. A first remark on the results is that the 8 symmetric and the 8 anti-symmetric base arrangements provide the same value of the magnetic field at the same fieldpoints.

From the comparison between the values of $B_{\text {Double }}$ from Table 3.3, it can be deduced that the 16 base arrangements even if they are geometrically different as to the values of $\Phi_{1}, \Phi_{2}$ at every field-point, yield practically equivalent magnetic field profiles at the chosen line section. This is confirmed by Fig. 3.2, where the field profiles from -1.0 m to 1.0 m relevant to the first 8 out of 16 base arrangements are plotted: the eight profiles, labelled as $B_{\text {DOUBLEI }}, B_{\text {DoUBLE } 2}, \ldots, B_{\text {DOUBLE }}$, can hardly - if ever - be distinguished from each other. Moreover, none of the 16 base arrangements resulted as the most unfavourable of all in terms of exposure to magnetic field, since some give rise to highest values at some points, other at other points. Thus, the problem of choosing the most unfavourable geometrical arrangement among all possible arrangements arises.


Figure 3.2: Comparison between the values of the $B_{\text {Double }}$ of the first 8 hypothetical symmetric cases for $I=280 \mathrm{~A}$.

A hint at solving this problem in a conservative way comes from the parametrical analysis discussed in Chapter 2. This analysis concluded in setting a constant value for the angle $\Phi=60^{\circ}$ which yields the maximum value of $B$ and corresponds to the most dangerous case in terms of exposure to magnetic field. This choice corresponds to a geometrical arrangement that cannot be found in practice, but it provides at every field point a maximum value of the magnetic field generated by each single circuit. Nevertheless, this is not always true for their vector sum, $B_{\text {DOUBLE,max }}$. This is made clear in Tables 3.3 and 3.4 , where the values of $B_{\text {DOUBLE }}$ of the first 8 symmetric base arrangements are listed (the 8 anti-symmetric cases are omitted since they provide the same value for $B_{\text {DOUBLE }}$ ), and compared with the values of $B_{\text {DOUBLE, max }}$ for $\Phi=60^{\circ}$ through the relevant percent deviations $\Delta B \%$.

Table 3.3: $B_{\text {DOUBLE }}$ values of the first 8 symmetric base arrangementsfor $I=280$ A.

| $\begin{gathered} \mathbf{x} \\ {[\mathrm{m}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 1} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 2} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 3} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 4} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 5} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 6} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 7} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 8} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DovbLE,max }} \\ {[\mu \mathrm{T}]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 1.6986 | 1.6999 | 1.6999 | 1.6986 | 1.6986 | 1.6999 | 1.6986 | 1.6999 | 1.7001 |
| -0.9 | 2.4829 | 2.4864 | 2.4864 | 2.4829 | 2.4838 | 2.4855 | 2.4838 | 2.4855 | 2.4866 |
| -0.8 | 3.4681 | 3.4757 | 3.4757 | 3.4681 | 3.4717 | 3.4721 | 3.4717 | 3.4721 | 3.4757 |
| -0.7 | 4.5350 | 4.5469 | 4.5469 | 4.5350 | 4.5437 | 4.5382 | 4.5436 | 4.5382 | 4.5474 |
| -0.6 | 5.4192 | 5.4296 | 5.4295 | 5.4193 | 5.4321 | 5.4167 | 5.4320 | 5.4168 | 5.4336 |
| -0.5 | 5.7946 | 5.7946 | 5.7946 | 5.7946 | 5.8050 | 5.7842 | 5.8049 | 5.7843 | 5.8049 |
| -0.4 | 5.5038 | 5.4938 | 5.4937 | 5.5038 | 5.5065 | 5.4911 | 5.5061 | 5.4915 | 5.5076 |
| -0.3 | 4.7131 | 4.7020 | 4.7021 | 4.7130 | 4.7106 | 4.7046 | 4.7096 | 4.7056 | 4.7131 |
| -0.2 | 3.7868 | 3.7803 | 3.7810 | 3.7861 | 3.7846 | 3.7826 | 3.7823 | 3.7849 | 3.7853 |
| -0.1 | 3.0730 | 3.0692 | 3.0716 | 3.0706 | 3.0729 | 3.0694 | 3.0687 | 3.0735 | 3.0695 |
| -0.0 | 2.8197 | 2.8149 | 2.8200 | 2.8147 | 2.8200 | 2.8147 | 2.8149 | 2.8197 | 2.8136 |
| 0.1 | 3.1168 | 3.1109 | 3.1170 | 3.1106 | 3.1155 | 3.1121 | 3.1126 | 3.1151 | 3.1102 |
| 0.2 | 3.8460 | 3.8437 | 3.8461 | 3.8437 | 3.8451 | 3.8447 | 3.8449 | 3.8449 | 3.8436 |
| 0.3 | 4.7639 | 4.7685 | 4.7639 | 4.7685 | 4.7673 | 4.7652 | 4.7652 | 4.7673 | 4.7687 |
| 0.4 | 5.5409 | 5.5482 | 5.5409 | 5.5482 | 5.5499 | 5.5392 | 5.5391 | 5.5499 | 5.5510 |
| 0.5 | 5.8254 | 5.8254 | 5.8254 | 5.8254 | 5.8341 | 5.8166 | 5.8166 | 5.8342 | 5.8342 |
| 0.6 | 5.4474 | 5.4380 | 5.4474 | 5.4380 | 5.4497 | 5.4358 | 5.4357 | 5.4497 | 5.4511 |
| 0.7 | 4.5564 | 4.5453 | 4.5564 | 4.5453 | 4.5534 | 4.5483 | 4.5483 | 4.5534 | 4.5570 |
| 0.8 | 3.4806 | 3.4732 | 3.4806 | 3.4732 | 3.4767 | 3.4771 | 3.4771 | 3.4767 | 3.4806 |
| 0.9 | 2.4888 | 2.4853 | 2.4888 | 2.4853 | 2.4862 | 2.4879 | 2.4879 | 2.4862 | 2.4889 |
| 1.0 | 1.7010 | 1.6998 | 1.7010 | 1.6998 | 1.6998 | 1.7010 | 1.7010 | 1.6998 | 1.7013 |

Table 3.4: Relevant percent errors of the values $B_{\text {DOUBLE }}$ of first 8 symmetric base arrangements with respect to $B_{\text {Double, } \max }(I=280$ A).

| $\begin{gathered} \mathbf{x} \\ {[\mathbf{m}]} \end{gathered}$ | $\begin{gathered} \Delta B_{1} \\ {[\%]} \end{gathered}$ | $\begin{aligned} & \Delta B_{2} \\ & {[\%]} \end{aligned}$ | $\begin{gathered} \Delta B_{3} \\ {[\%]} \end{gathered}$ | $\begin{aligned} & \Delta B_{4} \\ & {[\%]} \end{aligned}$ | $\begin{aligned} & \Delta B_{5} \\ & {[\%]} \end{aligned}$ | $\begin{aligned} & \Delta B_{6} \\ & {[\%]} \end{aligned}$ | $\begin{aligned} & \Delta B_{7} \\ & {[\%]} \end{aligned}$ | $\begin{aligned} & \Delta B_{\mathbf{8}} \\ & {[\%]} \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | -0.09 | -0.02 | -0.02 | -0.09 | -0.09 | -0.02 | -0.09 | -0.02 |
| -0.9 | -0.15 | -0.01 | -0.01 | -0.15 | -0.11 | -0.04 | -0.11 | -0.04 |
| -0.8 | -0.22 | -0.00 | -0.00 | -0.22 | -0.12 | -0.10 | -0.12 | -0.10 |
| -0.7 | -0.27 | -0.01 | -0.01 | -0.27 | -0.08 | -0.20 | -0.08 | -0.20 |
| -0.6 | -0.26 | -0.07 | -0.07 | -0.26 | -0.03 | -0.31 | -0.03 | -0.31 |
| -0.5 | -0.18 | -0.18 | -0.18 | -0.18 | -0.01 | -0.36 | 0.00 | -0.36 |
| -0.4 | -0.07 | -0.25 | -0.25 | -0.07 | -0.02 | -0.30 | -0.03 | -0.29 |
| -0.3 | 0.00 | -0.23 | -0.23 | -0.00 | -0.05 | -0.18 | -0.07 | -0.16 |
| -0.2 | 0.04 | -0.13 | -0.11 | 0.02 | -0.02 | -0.07 | -0.08 | -0.01 |
| -0.1 | 0.11 | -0.01 | 0.07 | 0.04 | 0.11 | -0.00 | -0.03 | 0.13 |
| -0.0 | 0.22 | 0.05 | 0.23 | 0.04 | 0.23 | 0.04 | 0.05 | 0.22 |
| 0.1 | 0.21 | 0.02 | 0.22 | 0.02 | 0.17 | 0.06 | 0.08 | 0.16 |
| 0.2 | 0.06 | 0.00 | 0.07 | 0.00 | 0.04 | 0.03 | 0.04 | 0.03 |
| 0.3 | -0.10 | -0.00 | -0.10 | -0.00 | -0.03 | -0.08 | -0.07 | -0.03 |
| 0.4 | -0.18 | -0.05 | -0.18 | -0.05 | -0.02 | -0.21 | -0.21 | -0.02 |
| 0.5 | -0.15 | -0.15 | -0.15 | -0.15 | -0.00 | -0.30 | -0.30 | -0.00 |
| 0.6 | -0.07 | -0.24 | -0.07 | -0.24 | -0.03 | -0.28 | -0.28 | -0.03 |
| 0.7 | -0.01 | -0.26 | -0.01 | -0.26 | -0.08 | -0.19 | -0.19 | -0.08 |
| 0.8 | -0.00 | -0.21 | -0.00 | -0.21 | -0.11 | -0.10 | -0.10 | -0.11 |
| 0.9 | -0.01 | -0.14 | -0.01 | -0.14 | -0.11 | -0.04 | -0.04 | -0.11 |
| 1.0 | -0.02 | -0.09 | -0.02 | -0.09 | -0.09 | -0.02 | -0.02 | -0.09 |

The relevant percent deviations $\Delta B_{1}, \Delta B_{2}, \ldots, \Delta B_{8}$, from Table 3.4, never exceed $1 \%$ meaning that the choice of setting the angle $\Phi=60^{\circ}$ constant for both cables was reasonable.

### 3.1.1.2 System Symmetry

Considering the twisted double-circuit power line of the first case-study, the magnetic field generated by the 8 symmetric phase configurations along with $B_{\text {DOUBLE, max }}$ was also calculated on a straight horizontal line that spreads from -1.0 m to 1.0 m at both ends of the line axis and whose vertical distance from the conductors is $H=-0.5$ mthis time (see Fig. 3.3). This analysis has the purpose of examining the system symmetry in terms of magnetic induction generated and the results are depicted at Table 3.5.


Figure 3.3: Sketch of the "reference arrangement" of the double-circuit twisted three-phase cable line chosen as first case-study for the vector sum of the Cartesian magnetic field components calculated $H=0.5 \mathrm{~m}$ above the cables. The relevant radial $\left(B_{r}\right)$ and azimuthal $\left(B_{\varphi}\right)$ field components generated by circuit 1 (green arrows) and 2 (red arrows) at five equally spaced field-points along the $x$-axis are shown.

Table 3.5: $B_{\text {DoUBLE }}$ values of the first 8 symmetric base arrangements for the configuration of Fig. 3.3 for $I_{l}=I_{2}=280 \mathrm{~A}$.

| $\begin{gathered} \mathbf{x} \\ {[\mathrm{m}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 1} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 2} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 3} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 4} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 5} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, } 6} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE }, 7} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \boldsymbol{B}_{\text {DOUBLE, }, 8} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\boldsymbol{B}_{\text {DoUbLE, max }}$ $[\mu \mathrm{T}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 1.6998 | 1.7010 | 1.7010 | 1.6998 | 1.7010 | 1.6998 | 1.7010 | 1.6998 | 1.7013 |
| -0.9 | 2.4853 | 2.4888 | 2.4888 | 2.4853 | 2.4879 | 2.4862 | 2.4879 | 2.4862 | 2.4889 |
| -0.8 | 3.4732 | 3.4806 | 3.4806 | 3.4732 | 3.4771 | 3.4767 | 3.4771 | 3.4767 | 3.4806 |
| -0.7 | 4.5453 | 4.5564 | 4.5564 | 4.5453 | 4.5483 | 4.5534 | 4.5483 | 4.5534 | 4.5570 |
| -0.6 | 5.4380 | 5.4474 | 5.4474 | 5.4380 | 5.4358 | 5.4497 | 5.4357 | 5.4497 | 5.4511 |
| -0.5 | 5.8254 | 5.8254 | 5.8254 | 5.8254 | 5.8166 | 5.8341 | 5.8166 | 5.8342 | 5.8342 |
| -0.4 | 5.5482 | 5.5409 | 5.5409 | 5.5482 | 5.5392 | 5.5499 | 5.5391 | 5.5499 | 5.5510 |
| -0.3 | 4.7685 | 4.7639 | 4.7639 | 4.7685 | 4.7652 | 4.7673 | 4.7652 | 4.7673 | 4.7687 |
| -0.2 | 3.8437 | 3.8460 | 3.8461 | 3.8437 | 3.8447 | 3.8451 | 3.8449 | 3.8449 | 3.8436 |
| -0.1 | 3.1109 | 3.1168 | 3.1170 | 3.1106 | 3.1121 | 3.1155 | 3.1126 | 3.1151 | 3.1102 |
| -0.0 | 2.8149 | 2.8197 | 2.8200 | 2.8147 | 2.8147 | 2.8200 | 2.8149 | 2.8197 | 2.8136 |
| 0.1 | 3.0692 | 3.0730 | 3.0716 | 3.0706 | 3.0694 | 3.0729 | 3.0687 | 3.0735 | 3.0695 |
| 0.2 | 3.7803 | 3.7868 | 3.7810 | 3.7861 | 3.7826 | 3.7846 | 3.7823 | 3.7849 | 3.7853 |
| 0.3 | 4.7020 | 4.7131 | 4.7021 | 4.7130 | 4.7046 | 4.7106 | 4.7096 | 4.7056 | 4.7131 |
| 0.4 | 5.4938 | 5.5038 | 5.4937 | 5.5038 | 5.4911 | 5.5065 | 5.5061 | 5.4915 | 5.5076 |
| 0.5 | 5.7946 | 5.7946 | 5.7946 | 5.7946 | 5.7842 | 5.8050 | 5.8049 | 5.7843 | 5.8049 |
| 0.6 | 5.4296 | 5.4192 | 5.4295 | 5.4193 | 5.4167 | 5.4321 | 5.4320 | 5.4168 | 5.4336 |
| 0.7 | 4.5469 | 4.5350 | 4.5469 | 4.5350 | 4.5382 | 4.5437 | 4.5436 | 4.5382 | 4.5474 |
| 0.8 | 3.4757 | 3.4681 | 3.4757 | 3.4681 | 3.4721 | 3.4717 | 3.4717 | 3.4721 | 3.4757 |
| 0.9 | 2.4864 | 2.4829 | 2.4864 | 2.4829 | 2.4855 | 2.4838 | 2.4838 | 2.4855 | 2.4866 |
| 1.0 | 1.6999 | 1.6986 | 1.6999 | 1.6986 | 1.6999 | 1.6986 | 1.6986 | 1.6999 | 1.7001 |

Finally, the $B_{\text {DOUBLE, } \max }$ isolines of $3 \mu \mathrm{~T}, 10 \mu \mathrm{~T}$ and $100 \mu \mathrm{~T}$ were plotted in Fig. 3.4 for the twisted double-circuit power lines under exam.


Figure 3.4: $3 \mu T, 10 \mu T$ and $100 \mu T B_{\text {Double,max }}\left(\Phi_{1}=\Phi_{2}=60^{\circ}\right)$ isolines for the 'reference arrangement' of the twisted double-circuit for $I_{1}=I_{2}=280 \mathrm{~A}$.

### 3.1.2 "Worst Case" Approach

The simplest approach - but also the most approximated one - is the so-called "worst case" estimate. In this approach the magnetic field generated by the doublecircuit is calculated as the algebraic sum of the two rms values of the total field $B$ generated by each twisted cable. By calculating the rms values of the total field $B$ from circuit 1 and 2 according to the exact formula (2.12), $B_{1, \text { exact }}$ and $B_{2, \text { exact, }}$ respectively, one obtains the "worst case" estimate of $B$ generated by the double circuit, $B_{W C, \text { exact }}$, i.e.:

$$
\begin{equation*}
B_{W C, \text { exact }}=B_{1, \text { exact }}+B_{2, \text { exact }} \tag{3.5}
\end{equation*}
$$

This expression of $B$ is drastically approximated, since it is based on the hypothesis that the two magnetic field vectors have the same direction, but it is also the most conservative one, hence it serves as an upper limit reference for exposure evaluation purposes.

Relationship (3.5) can be further simplified by using, firstly, the approximated formula from the literature (2.17) and, secondly, the simplified innovative formula (2.23), in the $1^{\text {st }}$ and $2^{\text {nd }}$ level of approximation, to express the rms values of field $B$ from circuit 1 and 2 , and then by composing them according to the "worst-case" approach, thereby attaining respectively:

$$
\begin{equation*}
B_{W C, l i t}=B_{1, l i t}+B_{2, l i t} \tag{3.6}
\end{equation*}
$$

and

$$
\begin{gather*}
B_{W C, \text { simp } 1}=B_{1, \text { simp } 1}+B_{2, \text { simp } 1}  \tag{3.7.a}\\
B_{W C, \text { simp } 2}=B_{1, \text { simp } 2}+B_{2, \text { simp } 2} \tag{3.7.b}
\end{gather*}
$$

### 3.2 APPLICATIONS - SIMULATION RESULTS

The final results of the first case-study (see Fig. 3.1) are illustrated in Fig. 3.5, which displays the plots of $B_{D O U B L E}, B_{W C, \text { exact, }} B_{W C, \text { lit }}, B_{W C, \text { simpl },} B_{W C, \text { simp } 2}$ vs. horizontal distance $x$ from the pole at distance $H=0.5 \mathrm{~m}$ from both the conductors for $\Phi_{1}=60^{\circ}$, $\Phi_{2}=60^{\circ}$, and in Table 3.6, which reports the relevant numerical values and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simpl, }} B_{W C, \text { simp } 2}$ with respect to $B_{D O U B L E}$.


Figure 3.5: $B_{\text {Double, }} B_{W C, \text { exact }}, B_{W C, \text { lit }}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. horizontal distance x for the first case-study ( $\left.I_{l}=I_{2}=280 \mathrm{~A}\right)$.

Table 3.6: $B_{\text {DovbLE, }} B_{W C, \text { exact, }} B_{W C, l i i}, B_{W C, \text { simpl }}$, $B_{W C, \text { simp } 2} v s$. distance $x$ and percent errors of $B_{W C, \text { exact, }} B_{W C, \text { lii, }}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {DOUBLE }}$ for the first case study $\left(I_{l}=I_{2}=280 \mathrm{~A}\right)$.

| $\begin{gathered} \mathbf{x} \\ {[\mathrm{m}]} \end{gathered}$ | $\begin{gathered} \text { B }_{\text {DOUBLE }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | B $_{\text {WC,exact }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{wc}, \text {,lit }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\text {WC,simp1 }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \\ \hline \end{gathered}$ | B $_{\text {WC,Simp2 }}$ <br> [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 1.7001 | 1.7019 | 0.10 | 1.3746 | -19.15 | 1.7085 | 2.04 | 1.7349 | -0.03 |
| -0.9 | 2.4866 | 2.4904 | 0.16 | 1.9649 | -20.98 | 2.4720 | 0.75 | 2.5052 | -1.29 |
| -0.8 | 3.4757 | 3.4842 | 0.25 | 2.6829 | -22.81 | 3.4281 | -0.22 | 3.4680 | -2.28 |
| -0.7 | 4.5474 | 4.5659 | 0.41 | 3.4378 | -24.40 | 4.4694 | -0.73 | 4.5144 | -2.83 |
| -0.6 | 5.4336 | 5.4724 | 0.71 | 4.0553 | -25.37 | 5.3479 | -0.69 | 5.3961 | -2.85 |
| -0.5 | 5.8049 | 5.8821 | 1.33 | 4.3394 | -25.25 | 5.7505 | -0.07 | 5.8009 | -2.27 |
| -0.4 | 5.5076 | 5.6510 | 2.60 | 4.2075 | -23.61 | 5.5338 | 1.42 | 5.5859 | -0.81 |
| -0.3 | 4.7131 | 4.9582 | 5.20 | 3.7696 | -20.02 | 4.8755 | 4.57 | 4.9285 | 2.29 |
| -0.2 | 3.7853 | 4.1646 | 10.02 | 3.2512 | -14.11 | 4.1261 | 10.40 | 4.1788 | 8.01 |
| -0.1 | 3.0695 | 3.5788 | 16.59 | 2.8594 | -6.85 | 3.5765 | 18.20 | 3.6282 | 15.64 |
| -0.0 | 2.8136 | 3.3676 | 19.69 | 2.7161 | -3.46 | 3.3789 | 21.91 | 3.4301 | 19.24 |
| 0.1 | 3.1102 | 3.5788 | 15.07 | 2.8594 | -8.06 | 3.5765 | 16.66 | 3.6282 | 14.05 |
| 0.2 | 3.8436 | 4.1646 | 8.35 | 3.2512 | -14.41 | 4.1261 | 8.72 | 4.1788 | 6.29 |
| 0.3 | 4.7687 | 4.9582 | 3.97 | 3.7696 | -20.95 | 4.8755 | 3.35 | 4.9285 | 1.04 |
| 0.4 | 5.5510 | 5.6510 | 1.80 | 4.2075 | -24.20 | 5.5338 | 0.63 | 5.5859 | -1.62 |
| 0.5 | 5.8342 | 5.8821 | 0.82 | 4.3394 | -25.62 | 5.7505 | -0.57 | 5.8009 | -2.78 |
| 0.6 | 5.4511 | 5.4724 | 0.39 | 4.0553 | -25.60 | 5.3479 | -1.01 | 5.3961 | -3.17 |
| 0.7 | 4.5570 | 4.5659 | 0.20 | 3.4378 | -24.56 | 4.4694 | -0.93 | 4.5144 | -3.04 |
| 0.8 | 3.4806 | 3.4842 | 0.11 | 2.6829 | -22.92 | 3.4281 | -0.36 | 3.4680 | -2.42 |
| 0.9 | 2.4889 | 2.4904 | 0.06 | 1.9649 | -21.05 | 2.4720 | 0.65 | 2.5052 | -1.38 |
| 1.0 | 1.7013 | 1.7019 | 0.04 | 1.3746 | -19.20 | 1.7085 | 1.97 | 1.7349 | -0.09 |

It can be noted that the values of the expression $B_{W C, \text { exact }}$, result larger than the exact rms value of the total magnetic field, $B_{\text {DOUBLE }}$, in the mid-span area and the relevant percent errors arrive at $\sim 20 \%$. The explanation for this is that the "worst case estimate" results in an addiction of the rms values of the magnetic field of the singlecircuits in this area, whereas the exact vector analysis results in a subtraction of the magnetic field vectors of the single-circuits in the same area. Therefore, it is logical to obtain greater values with the "worst case" calculus.Additionally, it can be argued that the profiles brought by the expression $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ are practically overlapped to the ones brought by the expression $B_{W C \text { exact. }}$. This is not true for the profiles associated with the expression $B_{W C, l i t}$, that in addition yield much lower - thus less conservative - magnetic field estimates. Consequently the expression $B_{W C, s i m p}$, in both levels of approximation, is more reliable compared to the expression $B_{W C, l i t}$.

The analysis has been completed by reporting three more case-studies of interest:
Second case study: Same as the first case with the only difference that the vertical distance between the field points and the conductors is increased by 0.5 m , therefore resulting $H=1.0 \mathrm{~m}$ (see Fig. 3.6). The results of the second case simulations are shown in Fig. 3.7 and Table 3.7.


Figure 3.6: Geometrical and electrical data concerning the second case-study.


Figure 3.7: $B_{D O U B L E}, B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. horizontal distance $x$ for the second case-study $\left(I_{I}=I_{2}=280 \mathrm{~A}\right)$.

Table 3.7: $B_{D O U B L E}, B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, \text { exact, }} B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {DOUBLE }}$ for the second case study $\left(I_{1}=I_{2}=280 \mathrm{~A}\right)$.

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\mathbf{D O U B L E}}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\mathbf{W C}, \text { exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\mathbf{W C}, \mathbf{l i t}}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\mathbf{W C}, \text { simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\mathbf{W C}, \text { simp2 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 0.1930 | 0.1935 | 0.27 | 0.1696 | -12.10 | 0.2041 | 5.75 | 0.2094 | 8.50 |
| -0.9 | 0.2399 | 0.2408 | 0.39 | 0.2101 | -12.42 | 0.2533 | 5.61 | 0.2597 | 8.26 |
| -0.8 | 0.2869 | 0.2885 | 0.56 | 0.2507 | -12.62 | 0.3029 | 5.56 | 0.3103 | 8.13 |
| -0.7 | 0.3293 | 0.3320 | 0.82 | 0.2876 | -12.65 | 0.3479 | 5.65 | 0.3562 | 8.18 |
| -0.6 | 0.3618 | 0.3662 | 1.22 | 0.3168 | -12.43 | 0.3834 | 5.98 | 0.3925 | 8.49 |
| -0.5 | 0.3809 | 0.3878 | 1.81 | 0.3355 | -11.91 | 0.4061 | 6.61 | 0.4157 | 9.14 |
| -0.4 | 0.3859 | 0.3961 | 2.64 | 0.3433 | -11.04 | 0.4152 | 7.60 | 0.4252 | 10.19 |
| -0.3 | 0.3796 | 0.3937 | 3.72 | 0.3422 | -9.85 | 0.4135 | 8.94 | 0.4236 | 11.60 |
| -0.2 | 0.3677 | 0.3858 | 4.91 | 0.3364 | -8.52 | 0.4061 | 10.42 | 0.4162 | 13.18 |
| -0.1 | 0.3571 | 0.3780 | 5.85 | 0.3305 | -7.45 | 0.3986 | 11.61 | 0.4087 | 14.44 |
| -0.0 | 0.3533 | 0.3749 | 6.12 | 0.3281 | -7.12 | 0.3955 | 11.97 | 0.4056 | 14.83 |
| 0.1 | 0.3581 | 0.3780 | 5.55 | 0.3305 | -7.71 | 0.3986 | 11.30 | 0.4087 | 14.12 |
| 0.2 | 0.3694 | 0.3858 | 4.42 | 0.3364 | -8.94 | 0.4061 | 9.92 | 0.4162 | 12.66 |
| 0.3 | 0.3815 | 0.3937 | 3.19 | 0.3422 | -10.30 | 0.4135 | 8.39 | 0.4236 | 11.04 |
| 0.4 | 0.3877 | 0.3961 | 2.16 | 0.3433 | -11.45 | 0.4152 | 7.10 | 0.4252 | 9.67 |
| 0.5 | 0.3824 | 0.3878 | 1.41 | 0.3355 | -12.25 | 0.4061 | 6.19 | 0.4157 | 8.72 |
| 0.6 | 0.3629 | 0.3662 | 0.91 | 0.3168 | -12.70 | 0.3834 | 5.65 | 0.3925 | 8.16 |
| 0.7 | 0.3300 | 0.3320 | 0.59 | 0.2876 | -12.85 | 0.3479 | 5.41 | 0.3562 | 7.93 |
| 0.8 | 0.2874 | 0.2885 | 0.38 | 0.2507 | -12.78 | 0.3029 | 5.37 | 0.3103 | 7.94 |
| 0.9 | 0.2402 | 0.2408 | 0.26 | 0.2101 | -12.54 | 0.2533 | 5.47 | 0.2597 | 8.12 |
| 1.0 | 0.1932 | 0.1935 | 0.17 | 0.1696 | -12.18 | 0.2041 | 5.64 | 0.2094 | 8.39 |

Third Case-Study: Same as the second case, excepted that cable 2 has a larger cross-section than cable 1 , namely $3 \times 185 \mathrm{~mm}^{2}$, thus larger ampacity $I=360 \mathrm{~A}$, pitch $p=1.5 \mathrm{~m}$ and radius $\alpha=0.022 \mathrm{~m}$ (see Fig. 3.8). The results of the third case simulations are shown in Fig. 3.9 and Table 3.8.


Figure 3.8: Geometrical and electrical data concerning the third case-study.


Figure 3.9: $B_{D O U B L E}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2} v s$. horizontal distance $x$ for the third case-study $\left(I_{1}=280 \mathrm{~A}\right.$ and $\left.I_{2}=360 \mathrm{~A}\right)$.

Table 3.8: $B_{\text {DOUBLE }}, B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{D O U B L E}$ for the third case study $\left(I_{1}=280 \mathrm{~A}\right.$ and $I_{2}=360 \mathrm{~A}$ ).

| $\mathbf{x}$ <br> $[\mathbf{m}]$ | $\mathbf{B}_{\text {DOUBLE }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\mathbf{W C}, \text { exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\mathbf{W C}, \text { lit }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\mathbf{W C}, \text { simp } 1}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\mathbf{W C}, \text { simp } 2}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 0.2016 | 0.2029 | 0.62 | 0.1782 | -11.64 | 0.2138 | 6.03 | 0.2193 | 8.73 |
| -0.9 | 0.2520 | 0.2541 | 0.84 | 0.2221 | -11.85 | 0.2672 | 6.05 | 0.2738 | 8.65 |
| -0.8 | 0.3038 | 0.3073 | 1.16 | 0.2676 | -11.91 | 0.3226 | 6.19 | 0.3303 | 8.71 |
| -0.7 | 0.3525 | 0.3582 | 1.62 | 0.3111 | -11.75 | 0.3756 | 6.54 | 0.3842 | 9.00 |
| -0.6 | 0.3936 | 0.4025 | 2.25 | 0.3491 | -11.31 | 0.4218 | 7.17 | 0.4313 | 9.58 |
| -0.5 | 0.4241 | 0.4372 | 3.10 | 0.3793 | -10.58 | 0.4585 | 8.11 | 0.4687 | 10.51 |
| -0.4 | 0.4443 | 0.4626 | 4.12 | 0.4017 | -9.59 | 0.4858 | 9.34 | 0.4964 | 11.73 |
| -0.3 | 0.4581 | 0.4817 | 5.16 | 0.4189 | -8.56 | 0.5068 | 10.64 | 0.5177 | 13.02 |
| -0.2 | 0.4720 | 0.4998 | 5.89 | 0.4350 | -7.84 | 0.5267 | 11.60 | 0.5378 | 13.94 |
| -0.1 | 0.4931 | 0.5225 | 5.96 | 0.4545 | -7.83 | 0.5511 | 11.77 | 0.5623 | 14.04 |
| -0.0 | 0.5252 | 0.5531 | 5.30 | 0.4799 | -8.63 | 0.5831 | 11.03 | 0.5944 | 13.18 |
| 0.1 | 0.5675 | 0.5912 | 4.18 | 0.5109 | -9.98 | 0.6223 | 9.65 | 0.6338 | 11.68 |
| 0.2 | 0.6139 | 0.6322 | 2.99 | 0.5436 | -11.44 | 0.6639 | 8.15 | 0.6755 | 10.04 |
| 0.3 | 0.6548 | 0.6679 | 2.00 | 0.5717 | -12.69 | 0.6997 | 6.85 | 0.7114 | 8.64 |
| 0.4 | 0.6800 | 0.6888 | 1.28 | 0.5876 | -13.60 | 0.7202 | 5.90 | 0.7318 | 7.62 |
| 0.5 | 0.6815 | 0.6870 | 0.81 | 0.5850 | -14.16 | 0.7177 | 5.31 | 0.7290 | 6.97 |
| 0.6 | 0.6556 | 0.6589 | 0.50 | 0.5611 | -14.41 | 0.6884 | 5.00 | 0.6992 | 6.65 |
| 0.7 | 0.6042 | 0.6061 | 0.32 | 0.5171 | -14.42 | 0.6340 | 4.93 | 0.6439 | 6.58 |
| 0.8 | 0.5339 | 0.5350 | 0.20 | 0.2507 | -14.22 | 0.5607 | 5.02 | 0.5696 | 6.69 |
| 0.9 | 0.4534 | 0.4540 | 0.13 | 0.2101 | -13.89 | 0.4770 | 5.21 | 0.4848 | 6.92 |
| 1.0 | 0.3714 | 0.3717 | 0.09 | 0.1696 | -13.45 | 0.3916 | 5.46 | 0.3982 | 7.22 |

FourthCase-Study: Same overall arrangement as in the third case with the only difference that cable 2 is 0.5 m farther away from the soil with respect to cable 1 . Hence the field points lie on a straight horizontal line whose vertical distance from cable 1 is $H_{l}=1.0 \mathrm{~m}$ and from cable 2 is $H_{2}=1.5 \mathrm{~m}$ (see Fig 3.10). Figure 3.11 and Table 3.9 show the results of the fourth case simulation developed in MATLAB ${ }^{\text {TM }}$ environment.


Figure 3.10: Geometrical and electrical data concerning the fourth case-study.


Figure 3.11: $B_{\text {DOUBLE }}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2} v s$. horizontal distance $x$ for the fourth case-study $\left(I_{1}=280 \mathrm{~A}\right.$ and $\left.I_{2}=360 \mathrm{~A}\right)$.

Table 3.9: $B_{D O U B L E}, B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2} v s$. distance $x$ and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{D O U B L E}$ for the fourth case study ( $I_{1}=280 \mathrm{~A}$ and $I_{2}=360 \mathrm{~A}$ ).

| $\begin{gathered} \mathbf{x} \\ {[\mathrm{m}]} \end{gathered}$ | B $_{\text {Double }}$ <br> [ $\mu \mathrm{T}$ ] | $\mathbf{B}_{\mathrm{WC}, \text { exact }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \mathrm{it}} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\text {WC, simp1 }}$ <br> $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | B wC,simp2 <br> $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.0 | 0.1910 | 0.1911 | 0.08 | 0.1675 | -12.31 | 0.2015 | 5.49 | 0.2066 | 8.19 |
| -0.9 | 0.2365 | 0.2367 | 0.11 | 0.2064 | -12.73 | 0.2489 | 5.25 | 0.2550 | 7.84 |
| -0.8 | 0.2814 | 0.2818 | 0.16 | 0.2446 | -13.06 | 0.2956 | 5.05 | 0.3026 | 7.55 |
| -0.7 | 0.3204 | 0.3211 | 0.22 | 0.2779 | -13.28 | 0.3362 | 4.91 | 0.3440 | 7.36 |
| -0.6 | 0.3480 | 0.3491 | 0.33 | 0.3015 | -13.36 | 0.3650 | 4.89 | 0.3734 | 7.30 |
| -0.5 | 0.3599 | 0.3616 | 0.48 | 0.3122 | -13.26 | 0.3779 | 5.01 | 0.3865 | 7.41 |
| -0.4 | 0.3544 | 0.3569 | 0.70 | 0.3085 | -12.95 | 0.3732 | 5.31 | 0.3818 | 7.72 |
| -0.3 | 0.3333 | 0.3367 | 1.01 | 0.2919 | -12.43 | 0.3526 | 5.80 | 0.3608 | 8.25 |
| -0.2 | 0.3008 | 0.3052 | 1.46 | 0.2657 | -11.67 | 0.3203 | 6.49 | 0.3279 | 9.02 |
| -0.1 | 0.2625 | 0.2678 | 2.04 | 0.2344 | -10.69 | 0.2818 | 7.38 | 0.2887 | 9.99 |
| -0.0 | 0.2236 | 0.2297 | 2.75 | 0.2022 | -9.56 | 0.2424 | 8.40 | 0.2484 | 11.11 |
| 0.1 | 0.1880 | 0.1946 | 3.50 | 0.1722 | -8.37 | 0.2057 | 9.44 | 0.2110 | 12.23 |
| 0.2 | 0.1578 | 0.1643 | 4.14 | 0.1462 | -7.33 | 0.1740 | 10.29 | 0.1785 | 13.15 |
| 0.3 | 0.1334 | 0.1394 | 4.50 | 0.1246 | -6.61 | 0.1477 | 10.75 | 0.1516 | 13.64 |
| 0.4 | 0.1142 | 0.1193 | 4.46 | 0.1069 | -6.35 | 0.1264 | 10.71 | 0.1297 | 13.60 |
| 0.5 | 0.0990 | 0.1030 | 4.05 | 0.0925 | -6.50 | 0.1091 | 10.23 | 0.1119 | 13.08 |
| 0.6 | 0.0864 | 0.0894 | 3.43 | 0.0805 | -6.91 | 0.0946 | 9.47 | 0.0970 | 12.27 |
| 0.7 | 0.0756 | 0.0776 | 2.74 | 0.0700 | -7.42 | 0.0821 | 8.62 | 0.0841 | 11.36 |
| 0.8 | 0.0657 | 0.0670 | 2.11 | 0.0605 | -7.88 | 0.0708 | 7.81 | 0.0726 | 10.50 |
| 0.9 | 0.0564 | 0.0573 | 1.59 | 0.0518 | -8.23 | 0.0604 | 7.09 | 0.0619 | 9.75 |
| 1.0 | 0.0478 | 0.0484 | 1.18 | 0.0438 | -8.46 | 0.0509 | 6.47 | 0.0522 | 9.12 |

A general comment deriving from second, third and fourth case is that the values of the expression $B_{W C, \text { exact }}$ result almost equal to the exact rms value of the total magnetic field, $B_{\text {DOUBLE }}$, with a maximum relevant percent error of $\sim 6 \%$. Additionally, the errors brought by simplified formula (2.23), in both levels of approximation, exhibit the same behavior in all cases, being relatively small- $B_{W C, \text { simpl }}$ not exceeding $12 \%$ and $B_{W C, s i m p 2}$ not exceeding $15 \%$ - and positive, providing, thus a fast, simple and conservative evaluation of the magnetic field generated by doublecircuit three-phase cable lines. To conclude, the error of the expression $B_{W C, l i t}(2.17)$ results larger and always negative, underestimating the total field.

### 3.3 IN SITU MEASUREMENTS

The proposed theoretical approach was experimentally validated by measurements in situ. That required finding easily accessible double-circuit twisted cable lines with known instant values of rms line currents, fact that was proved to be extremely difficult in practice. The two main reasons for this were, firstly, that the electrical utilities were not prone to make their lines accessible and to provide data about the currents and, secondly, that easily-accessible double circuit twisted lines with fully known geometric parameters are very scarce to be found.The only solution found was performing the measurements inside an inspection well of the underground doublecircuit twisted three-phase MV line that connects the wind turbines of the Casoni di Romagna Wind Park - Monterenzio (BO), the largest wind park in Northern Italy with the nearest substation. Fig. 3.12 is an actual photo of the measurement site and Fig. 3.13 is its schematization reporting all the needed geometric information.


Figure 3.12: Photo of the twisted double-circuit power line inside the inspection well, along with the magnetic field measuring probe (the black cable is a signal cable and does not affect the measurements).


Figure 3.13: Schematization of the twisted underground double-circuit power lines.
As Fig. 3.12 shows, the twisted double-circuit cables could be hardly accessed. This required the use of an insulating non-magnetic cane for inserting the probe inside the well, thereby limiting the precision concerning all the distances measured. Moreover, the constancy of the pitch of the single helix and the perfect parallelism of the two helixes were not exactly matched. In addition the instant rms values of the line currents provided by the distribution were relevant to one single phase current per circuit and, of course, there is no certainty that the three-phase currents of each circuit are symmetric all the time. Despite all these inconveniences, using the physical data available the magnetic fields were estimated by the software with acceptable errors.

The two underground circuits are placed at different heights above the well's ground; they are parallel to each other, but not to the ground, forming an angle of $\sim 25^{\circ}$. The two twisted cable circuits are ARE4H5EX [65] cables with cross-section $3 \times 185 \mathrm{~mm}^{2}$, rated voltage 12 (phase-to-ground) $/ 20$ (phase-to-phase) kV , ampacity $I=360 \mathrm{~A}$, pitch $p=1.20 \mathrm{~m}$ (as actually measured) and radius $\alpha=0.022 \mathrm{~m}$.

The rms value of the total magnetic induction field, $B_{\text {DoubLE,meas }}$, has been measured and recorded at subsequent instants of time using a magnetic field measuring device, that consists essentially of a 3D field probe connected with a datalogger system using an optical fiber. The 3D field probe is a cube of side length 0.09 m.

### 3.3.1 Measurements on December $6^{\text {th }} 2011$

Two measurements campaigns were effectuated during this PhD course at the same location, one on December $6^{\text {th }} 2011$ and the other on December $18^{\text {th }}$ 2012. This is due to the low power production on December $6^{\text {th }} 2011$ that resulted in very small current values (see Fig. 3.14) and therefore in small values of the magnetic field generated by the twisted double-circuit cable line under exam. That constrained the repetition of the measurements on a windier day.


Figure 3.14: Instant values of the currents $I_{1}$ and $I_{2}$ for the time interval of the measurements on 06/12/2011

Nonetheless, all measurement data from December $6{ }^{\text {th }} 2011$ are illustrated in Table 3.10, and the comparison between the experimental values of $B_{\text {DOUBLE,meas }}$ and the theoretical calculated values of $B_{D O U B L E}, B_{W C, \text { exact, }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ are presented, highlighting the percent errors of $B_{D O U B L E}, B_{W C, \text { exact, }} B_{W C, l i t}, B_{W C, \text { simpl }}$ and $B_{W C, s i m p 2}$ with respect to $B_{\text {DOUBLE,meas. }}$. Measurements took place $0.0 \mathrm{~m}, 0.2 \mathrm{~m}, 0.5 \mathrm{~m}$ and 1.0 m above each cable and $0.0 \mathrm{~m}, 0.2 \mathrm{~m}, 0.5 \mathrm{~m}, 1.0 \mathrm{~m}, 1.5 \mathrm{~m}$ (at ground level) and 2.5 m ( 1.0 m from ground level) above the midway point between the two cables.

Table 3.10: Measurement results ( 06/12/2011): $B_{\text {Double,meas, }} B_{\text {Double,calc, }} B_{W C, \text { exact, }} B_{W C, \text { li, }}$ $B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2}$ at different field-points and percent errors of $B_{D O U B L E \text { calc, }} B_{W C, \text { exact, }} B_{W C, \text { li, }}$ $B_{W C, \text { simpl }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {Double,meas. }}$.

| $\begin{gathered} \mathbf{B}_{\text {Double,meas }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {DOUBLE, calc }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { exact }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{wc}, \text { lit }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{WC}, \text { simp1 }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{WC}, \text { simp2 }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 9.550 | 11.072 | -15.9 | 12.938 | 35.5 | 4.059 | -57.5 |  |  |  |  |
| 9.458 | 11.067 | -17.0 | 12.932 | 36.7 | 4.058 | -57.1 |  |  |  |  |
| 9.404 | 11.636 | -23.7 | 13.529 | 43.9 | 4.220 | -55.1 |  |  |  |  |
| 6.308 | 11.762 | -86.5 | 13.741 | 118 | 4.310 | -31.7 |  |  |  |  |


| 0.0 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 5.413 | 10.605 | -95.9 | 11.079 | -105 | 3.486 | 35.6 |  |  |  |  |
| 0.0 mm midway between the two cables |  |  |  |  |  |  |  |  |  |  |
| 15.960 | 13.217 | 17.2 | 23.497 | -47.2 | 5.664 | 64.5 |  |  |  |  |
| 0.2 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 2.266 | 2.342 | -3.36 | 2.436 | -7.51 | 1.423 | 37.2 | 2.997 | -32.3 | 2.936 | -29.6 |
| 0.2 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 1.449 | 1.649 | -13.8 | 1.684 | -16.2 | 1.038 | 28.4 | 1.951 | 34.7 | 1.919 | -32.4 |
| 0.2 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 2.916 | 3.516 | -20.6 | 3.709 | -27.2 | 1.925 | 34.0 | 6.016 | -106 | 5.804 | -99.0 |
| 0.5 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 0.378 | 0.357 | 5.50 | 0.360 | 4.73 | 0.277 | 26.7 | 0.355 | 6.09 | 0.355 | 6.18 |
| 0.5 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 0.381 | 0.245 | 35.7 | 0.246 | 35.3 | 0.194 | 49.2 | 0.244 | 36.0 | 0.244 | 36.0 |
| 0.367 | 0.245 | 33.2 | 0.246 | 32.9 | 0.194 | 47.2 | 0.244 | 33.6 | 0.244 | 33.5 |
| 0.5 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.435 | 0.427 | 1.83 | 0.431 | 0.938 | 0.325 | 25.4 | 0.425 | 2.31 | 0.424 | 2.52 |
| 1.0 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 0.108 | 0.0211 | 80.4 | 0.021 | 80.4 | 0.019 | 82.9 | 0.022 | 79.5 | 0.022 | 79.4 |
| 1.0 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 0.102 | 0.0156 | 84.7 | 0.016 | 84.7 | 0014 | 86.6 | 0.016 | 83.9 | 0.017 | 83.8 |
| 1.0 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.108 | 0.034 | 68.3 | 0.034 | 68.2 | 0.030 | 72.7 | 0.036 | 67.2 | 0.036 | 67.0 |
| 0.113 | 0.038 | 66.0 | 0.039 | 65.9 | 0.033 | 70.7 | 0.040 | 64.8 | 0.04 | 64.6 |
| 0.078 | 0.021 | 72.8 | 0.021 | 72.8 | 0.019 | 76.3 | 0.022 | 71.5 | 0.022 | 71.4 |
| At ground level - $1.5 \mathbf{m}$ from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.055 | 0.003 | 95.2 | 0.003 | 95.2 | 0.002 | 95.7 | 0.003 | 94.9 | 0.003 | 94.9 |
| 1.0m from ground level - $2.5 \mathbf{~ m}$ from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.037 | 0.00 | 99.95 | 0.00 | 99.95 | 0.00 | 99.96 | 0.00 | 99.95 | 0.00 | 99.95 |

The percent errors of $B_{\text {DOUBLE, calc }}$ with respect to $B_{\text {DOUBLE,meas }}$ resulted satisfactory concerning the measurements made up to 0.5 m above each cable and above the midway point between the two cables. For distances greater than 1.0 m , the values of $B_{\text {DOUBLE, calc }}$ and $B_{\text {DOUBLE,meas }}$ begin to diverge significantly.

The expression $B_{W C, \text { exact }}$ approximates fairly well the measured values of $B$, in the range of measurements $<1.0 \mathrm{~m}$, as well as its simplifications, $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$. In the parametrical analysis that has resulted the innovative simplified expression (Chapter 2), the approximation of $\ln B$ vs. $r$ with the equation of a straight line plus a hyperbolic term was effectuated for distances $0.3 \mathrm{~m}<r<2.0 \mathrm{~m}$. That is the reason why for the measurements at 0.0 m from each cable and the midway point between the two cables the expressions $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ are omitted from Table 3.10 since they are away from their field of accuracy. The expression $B_{W C, l i t}$ has resulted a higher error with respect to $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ in the range of distances $>0.2 \mathrm{~m}$.

It must be pointed out, that measurements made at ground level and 1.0 m from it, were affected by the presence of a MV overhead single-circuit power line passing nearby.

### 3.3.2 Measurements on December $\mathbf{1 8}^{\text {th }} 2012$

Figure 3.15 illustrates the plots of the currents vs. time for the time interval of the measurements. Measurements took place $0.0 \mathrm{~m}, 0.3 \mathrm{~m}, 0.5 \mathrm{~m}, 0.8 \mathrm{~m}$ and 1.0 m above each cable and the midway point between the two cables, and at ground level, i.e. 1.5 m from cable 2. The measurement data are illustrated in Table 3.11, and the comparison between the experimental values of $B_{\text {DOUBLE,meas }}$ and the theoretical calculated values of $B_{D O U B L E, \text { calc, }}, B_{W C, \text { exact, }}, B_{W C, \text { lit }}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ are presented, highlighting the percent errors of $B_{D O U B L E, \text { calc, }} B_{W C, \text { exact, }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {DOUBLE,meas }}$.


Figure 3.15: Instant values of the currents $I_{1}$ and $I_{2}$ for the time interval of the measurements on 18/12/2012

Table 3.11: Measurement results (18/12/2012): $B_{\text {Double,meas, }} B_{\text {DoUBLE,calc, }} B_{W C, e x a c t,} B_{W C, l i t}$ $B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ at different field-points and percent errors of $B_{D o u b L E, \text { calc, }} B_{W C, \text { exact }} B_{W C, \text { lib }}$ $B_{W C, \text { simpl }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {DoUBLE, meas. }}$.

| B $_{\text {DOUBLE, meas }}$ $[\mu \mathrm{T}]$ | B $_{\text {DOUbLE, calc }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\text {wC, exact }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { lit }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \hline \mathbf{B}_{\mathrm{WC}, \text { simp } 1} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{WC}, \text { simp2 }}$ <br> $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 37.25 | 44.041 | -18.2 | 46.851 | -25.8 | 13.593 | 63.5 |  |  |  |  |
| 40.64 | 46.933 | -15.5 | 49.814 | -22.6 | 14.362 | 64.7 |  |  |  |  |
| 41.91 | 49.824 | -18.9 | 52.775 | -25.9 | 15.129 | 63.9 |  |  |  |  |
| 41.71 | 49.143 | -17.8 | 52.087 | -24.9 | 14.959 | 64.1 |  |  |  |  |
| 0.0 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 31.35 | 53.346 | -0.162 | 55.8 | -78.0 | 14.680 | 53.2 |  |  |  |  |
| 31.65 | 53.077 | -67.7 | 55.431 | -75.4 | 14.490 | 54.2 |  |  |  |  |
| 31.34 | 52.116 | -66.3 | 54.432 | -73.7 | 14.235 | 54.6 |  |  |  |  |


| 0.0 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 95.97 | 36.684 | 61.8 | 96.169 | -0.207 | 21.963 | 77.1 |  |  |  |  |
| 91.16 | 33.499 | 63.3 | 89.255 | 2.09 | 20.437 | 77.6 |  |  |  |  |
| 87.25 | 32.182 | 63.1 | 85.945 | 1.50 | 19.687 | 77.4 |  |  |  |  |
| 88.24 | 33.919 | 61.6 | 87.561 | 0.769 | 19.954 | 77.4 |  |  |  |  |
| 91.05 | 36.124 | 60.3 | 95.840 | -5.26 | 21.929 | 75.9 |  |  |  |  |
| 92.13 | 35.315 | 61.7 | 88.922 | 3.48 | 20.202 | 78.1 |  |  |  |  |
| 115.7 | 40.998 | 64.5 | 98.810 | 14.6 | 22.341 | 80.7 |  |  |  |  |
| 119.7 | 43.724 | 63.5 | 99.821 | 16.6 | 22.45 | 81.2 |  |  |  |  |
| 49.36 | 38.765 | 21.5 | 89.583 | -81.5 | 20.171 | 59.1 |  |  |  |  |
| 62.86 | 39.298 | 37.5 | 91.886 | -46.2 | 20.713 | 67.0 |  |  |  |  |
| 63.6 | 49.5 | 22.1 | 103.101 | -62.1 | 22.993 | 63.8 |  |  |  |  |
| 0.3 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 3.903 | 3.521 | 9.79 | 3.595 | 7.90 | 2.401 | 38.5 | 3.740 | 4.18 | 3.705 | 5.08 |
| 3.754 | 3.521 | 6.21 | 3.595 | 4.25 | 2.401 | 36.0 | 3.740 | 0.372 | 3.705 | 1.31 |
| 3.731 | 3.612 | 3.19 | 3.688 | 1.17 | 2.464 | 34.0 | 3.836 | -2.82 | 3.800 | -1.86 |
| 0.3 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 4.561 | 3.229 | 29.2 | 3.266 | 28.4 | 2.234 | 51.0 | 3.362 | 26.3 | 3.335 | 26.9 |
| 4.184 | 3.099 | 25.9 | 3.136 | 25.1 | 2.146 | 48.7 | 3.226 | 22.9 | 3.200 | 23.5 |
| 3.666 | 2.875 | 21.6 | 2.911 | 20.6 | 1.997 | 45.5 | 2.991 | 18.4 | 2.967 | 19.1 |
| 0.3 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 9.467 | 9.255 | 2.24 | 9.613 | -1.54 | 5.445 | 42.5 | 12.749 | -34.7 | 12.434 | -31.3 |
| 7.723 | 9.163 | -18.7 | 9.519 | -23.3 | 5.394 | 30.2 | 12.613 | -63.3 | 12.302 | -59.3 |
| 7.338 | 8.704 | -18.6 | 9.042 | -23.2 | 5.124 | 30.2 | 11.980 | -63.3 | 11.685 | -59.2 |
| 7.374 | 8.425 | -14.3 | 8.761 | 18.8 | 4.981 | 32.5 | 11.544 | -56.6 | 11.264 | -52.7 |
| 6.944 | 8.479 | -22.1 | 8.818 | -27.0 | 5.016 | 27.8 | 11.609 | -67.2 | 11.327 | -63.1 |
| 7.678 | 8.479 | -10.4 | 8.818 | -14.9 | 5.016 | 34.7 | 11.609 | -51.2 | 11.327 | -47.5 |
| 0.5 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 1.165 | 1.115 | 4.29 | 1.125 | 3.48 | 0.861 | 26.1 | 1.108 | 4.89 | 1.107 | 5.00 |
| 1.163 | 1.096 | 5.77 | 1.105 | 4.98 | 0.846 | 27.3 | 1.089 | 6.37 | 1.087 | 6.48 |
| 1.141 | 1.084 | 4.98 | 1.093 | 4.18 | 0.837 | 26.6 | 1.077 | 5.58 | 1.076 | 5.69 |
| 1.288 | 1.173 | 8.95 | 1.183 | 8.17 | 0.906 | 29.7 | 1.166 | 9.51 | 1.164 | 9.62 |
| 1.252 | 1.194 | 4.62 | 1.204 | 3.80 | 0.923 | 26.3 | 1.187 | 5.21 | 1.185 | 5.32 |
| 1.282 | 1.204 | 6.10 | 1.214 | 5.30 | 0.930 | 27.4 | 1.196 | 6.68 | 1.195 | 6.79 |
| 1.341 | 1.169 | 12.2 | 1.179 | 12.1 | 0.903 | 32.6 | 1.162 | 13.4 | 1.160 | 13.5 |
| 1.371 | 1.190 | 13.2 | 1.200 | 12.5 | 0.920 | 32.9 | 1.183 | 13.7 | 1.181 | 13.8 |
| 0.5 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 1.302 | 0.957 | 26.5 | 0.962 | 26.1 | 0.758 | 41.8 | 0.953 | 26.8 | 0.953 | 26.8 |
| 1.257 | 0.931 | 25.9 | 0.936 | 25.5 | 0.737 | 41.4 | 0.927 | 26.3 | 0.927 | 26.3 |
| 1.224 | 0.897 | 26.8 | 0.901 | 26.4 | 0.709 | 42.1 | 0.892 | 27.1 | 0.892 | 27.1 |
| 1.226 | 0.878 | 28.4 | 0.883 | 28.0 | 0.695 | 43.4 | 0.874 | 28.7 | 0.874 | 28.7 |
| 1.257 | 0.854 | 32.0 | 0.859 | 31.7 | 0.677 | 46.2 | 0.851 | 32.3 | 0.851 | 32.3 |
| 0.5 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 2.604 | 2.587 | 0.654 | 2.620 | -0.618 | 1.891 | 27.4 | 2.609 | -0.177 | 2.597 | 0.277 |
| 2.398 | 2.571 | -7.22 | 2.604 | -8.59 | 1.879 | 21.7 | 2.593 | -8.12 | 2.581 | -7.63 |
| 2.425 | 2.606 | -7.48 | 2.640 | -8.85 | 1.903 | 21.5 | 2.629 | -8.41 | 2.617 | -7.91 |
| 0.8 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 0.450 | 0.240 | 46.6 | 0.241 | 46.4 | 0.203 | 54.8 | 0.246 | 45.3 | 0.247 | 45.1 |
| 0.465 | 0.249 | 46.5 | 0.249 | 46.4 | 0.210 | 54.8 | 0.255 | 45.2 | 0.256 | 45.1 |
| 0.462 | 0.249 | 46.2 | 0.249 | 46.0 | 0.210 | 54.5 | 0.255 | 44.9 | 0.256 | 44.7 |
| 0.8 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 0.363 | 0.179 | 50.8 | 0.179 | 50.6 | 0.153 | 57.9 | 0.184 | 49.2 | 0.185 | 49.0 |
| 0.354 | 0.174 | 50.9 | 0.174 | 50.8 | 0.149 | 58.0 | 0.179 | 49.4 | 0.180 | 49.1 |
| 0.351 | 0.166 | 52.8 | 0.166 | 52.7 | 0.142 | 59.6 | 0.171 | 51.3 | 0.172 | 51.1 |
| 0.347 | 0.163 | 53.0 | 0.164 | 52.9 | 0.139 | 59.9 | 0.168 | 51.6 | 0.169 | 51.3 |
| 0.8 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.534 | 0.389 | 27.1 | 0.391 | 26.8 | 0.322 | 39.7 | 0.394 | 26.3 | 0.395 | 26.1 |
| 0.527 | 0.384 | 27.1 | 0.386 | 26.8 | 0.318 | 39.6 | 0.389 | 26.2 | 0.390 | 26.0 |


| 0.550 | 0.376 | 31.6 | 0.378 | 31.4 | 0.312 | 43.4 | 0.381 | 30.8 | 0.382 | 30.6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.539 | 0.381 | 29.4 | 0.382 | 29.1 | 0.316 | 41.5 | 0.385 | 28.5 | 0.386 | 28.3 |
| 0.521 | 0.368 | 29.3 | 0.370 | 29.1 | 0.305 | 41.5 | 0.373 | 28.5 | 0.374 | 28.3 |
| 1.0 m from cable 1 |  |  |  |  |  |  |  |  |  |  |
| 0.228 | 0.075 | 67.3 | 0.075 | 67.2 | 0.065 | 71.4 | 0.078 | 65.8 | 0.079 | 65.6 |
| 0.240 | 0.062 | 74.0 | 0.063 | 73.9 | 0.055 | 77.3 | 0.065 | 72.8 | 0.066 | 72.6 |
| 0.225 | 0.061 | 73.1 | 0.061 | 73.0 | 0.053 | 76.5 | 0.063 | 71.8 | 0.064 | 71.7 |
| 0.211 | 0.056 | 73.7 | 0.056 | 73.6 | 0.049 | 77.0 | 0.058 | 72.5 | 0.059 | 72.3 |
| 1.0 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 0.165 | 0.038 | 77.2 | 0.038 | 77.1 | 0.033 | 79.9 | 0.040 | 76.0 | 0.040 | 75.9 |
| 0.163 | 0.035 | 78.4 | 0.035 | 78.3 | 0.031 | 81.0 | 0.037 | 77.3 | 0.037 | 77.1 |
| 0.138 | 0.031 | 77.3 | 0.031 | 77.2 | 0.028 | 80.0 | 0.033 | 76.1 | 0.033 | 76.0 |
| 0.137 | 0.031 | 77.5 | 0.031 | 77.5 | 0.027 | 80.2 | 0.032 | 76.4 | 0.033 | 76.2 |
| 0.139 | 0.032 | 77.2 | 0.032 | 77.1 | 0.028 | 79.9 | 0.033 | 76.0 | 0.034 | 75.9 |
| 1.0 m from the midway point between the two cables |  |  |  |  |  |  |  |  |  |  |
| 0.203 | 0.081 | 60.3 | 0.081 | 60.2 | 0.069 | 65.8 | 0.084 | 58.9 | 0.084 | 58.7 |
| 0.218 | 0.084 | 61.7 | 0.084 | 61.6 | 0.072 | 67.0 | 0.087 | 60.3 | 0.087 | 60.1 |
| 0.220 | 0.084 | 61.7 | 0.085 | 61.6 | 0.073 | 67.0 | 0.088 | 60.2 | 0.088 | 60.0 |
| 0.218 | 0.086 | 60.5 | 0.086 | 60.4 | 0.074 | 66.0 | 0.893 | 59.1 | 0.090 | 58.8 |
| 0.217 | 0.088 | 59.5 | 0.088 | 59.4 | 0.076 | 65.1 | 0.091 | 58.0 | 0.092 | 57.8 |
| 0.226 | 0.088 | 61.1 | 0.088 | 61.0 | 0.076 | 66.5 | 0.091 | 59.7 | 0.092 | 59.5 |
| 0.224 | 0.087 | 61.4 | 0.087 | 61.3 | 0.075 | 66.7 | 0.090 | 60.0 | 0.090 | 59.8 |
| At ground level - 1.5 m from cable 2 |  |  |  |  |  |  |  |  |  |  |
| 0.112 | 0.003 | 97.6 | 0.003 | 97.6 | 0.003 | 97.8 | 0.003 | 97.4 | 0.003 | 97.4 |
| 0.110 | 0.003 | 97.6 | 0.003 | 97.6 | 0.002 | 97.8 | 0.003 | 97.5 | 0.003 | 97.5 |
| 0.104 | 0.003 | 97.5 | 0.003 | 97.5 | 0.002 | 97.7 | 0.003 | 97.3 | 0.003 | 97.3 |
| 0.104 | 0.003 | 97.5 | 0.003 | 97.5 | 0.002 | 97.7 | 0.003 | 97.3 | 0.003 | 97.3 |
| 0.102 | 0.003 | 97.5 | 0.003 | 97.5 | 0.002 | 97.7 | 0.003 | 97.3 | 0.003 | 97.3 |

The percent errors of $B_{\text {DOUBLE,calc }}$ with respect to $B_{\text {DOUBLE,meas }}$ resulted satisfactory concerning the measurements made at 0.3 m and 0.5 m above each cable and above the midway point between the two cables. The expression $B_{W C, \text { exact }}$ approximates fairly well the measured values of $B$, in this range of measurements, as well as its simplifications, $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp2 } 2}$. Also here, for measurements at 0.0 m from each cable and from the midway point between the two cables, $B_{W C, s i m p l}$ and $B_{W C, \text { simp } 2}$ are omitted from Table 3.11 since the innovative simplified formula, in both levels of approximation, is out of its application limits. The expression $B_{W C, l i t}$ has resulted a higher error with respect to $B_{W C, s i m p 1}$ and $B_{W C, \text { simp } 2}$.

For distances 0.8 m and 1.0 m above each cable and above the midway point between the two cables and 1.5 m above cable 2, the relative percent errors of $B_{\text {Double,calc }}$ with respect to $B_{\text {Double,meas }}$ seem to increase. The exact, approximated and simplified 'worst case' expressions present the same behaviour as $B_{\text {DOUBLE,calc. }}$.

Another supplementary set of measurements was also effectuated at $0.0 \mathrm{~m}, 0.3 \mathrm{~m}$, $0.5 \mathrm{~m}, 0.8 \mathrm{~m}$ and 1.0 m , placing the cane that holds the probe at the other end of the well opening i.e. 0.57 m left from cable 2, as shown in Fig. 3.16. The results of the second set of measurements effectuated are illustrated in Table 3.12.


Figure 3.16: Photo of the probe at 0.57 m left from cable 2 having positioned the cane at the other end of the well opening.

Table 3.12: Supplementary measurement results (18/12/2012): $B_{\text {DOUBLE,meas }}, B_{\text {DOUBLE,calc, }}$, $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ at different field-points (see Fig. 3.16) and percent errors of $B_{\text {DOUBLE, calc }}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{D O U B L E, \text { meas }}$.

| B $_{\text {DOUBLE,meas }}$ $[\mu \mathrm{T}]$ | $\mathbf{B}_{\text {DOUbLE, calc }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \hline \mathbf{B}_{\mathrm{WC}, \text { exact }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{wc}, \text {,it }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | B $_{\text {WC,simp1 }}$ <br> $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \hline \mathbf{B}_{\mathrm{WC}, \text { simp2 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 ~ m}$ |  |  |  |  |  |  |  |  |  |  |
| 0.340 | 1.053 | -207 | 1.054 | -207 | 0.757 | -121 | 1.056 | -208 | 1.051 | -206 |
| 0.356 | 1.053 | -196 | 1.054 | -196 | 0.757 | -113 | 1.056 | -197 | 1.051 | -195 |
| 0.359 | 1.098 | -206 | 1.100 | -206 | 0.791 | -120 | 1.102 | -207 | 1.097 | -205 |
| 0.354 | 1.037 | -193 | 1.039 | -193 | 0.747 | -111 | 1.040 | -194 | 1.035 | -192 |
| 0.3 m |  |  |  |  |  |  |  |  |  |  |
| 0.281 | 0.410 | -45.9 | 0.412 | -46.4 | 0.325 | -15.8 | 0.408 | -45.2 | 0.408 | -45.2 |
| 0.276 | 0.410 | -48.6 | 0.412 | -49.1 | 0.325 | -17.9 | 0.408 | -47.8 | 0.408 | -47.8 |
| 0.269 | 0.390 | -45.0 | 0.391 | -45.5 | 0.310 | -15.1 | 0.388 | -44.2 | 0.388 | -44.3 |
| 0.267 | 0.385 | -44.2 | 0.386 | -44.7 | 0.306 | -14.4 | 0.383 | -43.4 | 0.383 | -43.5 |
| 0.5m |  |  |  |  |  |  |  |  |  |  |
| 0.209 | 0.152 | 27.5 | 0.152 | 27.3 | 0.126 | 39.7 | 0.154 | 26.5 | 0.154 | 26.3 |
| 0.183 | 0.148 | 19.0 | 0.149 | 18.8 | 0.123 | 32.3 | 0.150 | 17.9 | 0.151 | 17.7 |
| 0.184 | 0.147 | 19.9 | 0.148 | 19.7 | 0.123 | 33.4 | 0.149 | 18.8 | 0.150 | 18.6 |
| 0.8 m |  |  |  |  |  |  |  |  |  |  |
| 0.152 | 0.041 | 73.4 | 0.041 | 73.3 | 0.035 | 76.8 | 0.042 | 72.2 | 0.043 | 72.0 |
| 0.152 | 0.041 | 73.3 | 0.041 | 73.2 | 0.035 | 76.7 | 0.042 | 72.1 | 0.043 | 71.9 |
| 0.142 | 0.040 | 71.9 | 0.04 | 71.8 | 0.035 | 75.5 | 0.042 | 70.6 | 0.042 | 70.5 |
| 1.0 m |  |  |  |  |  |  |  |  |  |  |
| 0.138 | 0.016 | 88.4 | 0.016 | 88.4 | 0.014 | 89.7 | 0.017 | 87.8 | 0.017 | 87.7 |
| 0.131 | 0.016 | 87.8 | 0.016 | 87.8 | 0.014 | 89.2 | 0.017 | 87.1 | 0.017 | 87.0 |
| 0.122 | 0.016 | 86.9 | 0.016 | 86.9 | 0.014 | 88.3 | 0.017 | 86.1 | 0.017 | 86.0 |

Measurements at 0.3 m and 0.5 m had a positive outcome, unlike measurements at distances higher than 0.8 m that have manifested significant errors.

Thus, the experiment results can be retained as satisfactory since the agreement between the code and the measurements is good where the field is higher. Discrepancies are found only where the field is either close to or even below the sensitivity of the measuring instruments. The reasons for these discrepancies are:

- the twisted double-circuit cables could be hardly accessed. The use of an insulating non-magnetic cane for inserting the probe inside the well was required, limiting the precision concerning all the distances measured.
- the constancy of the pitch of the single helix and the perfect parallelism of the two helixes were not exactly matched.
- the instant rms values of the line currents provided by the distribution utility were relevant to one single phase current per circuit and, of course, there is no certainty that the three-phase currents of each circuit are symmetric all the time.
- magnetic field values below $0.1 \mu \mathrm{~T}$ are difficult to be tracked because they are below the sensitivity of the measuring device. Additionally, $0.1 \mu \mathrm{~T}$ is the minimum average magnetic field value found everywhere due to the use of electricity.

For distances higher than 0.5 m , the measured and calculatedmagnetic field values diverge significantly. This issue has been treated in literature, where similar divergent results have been obtained between measured and calculated magnetic fields in the vicinity of twisted cables: Pettersson et al. in [14], [15] concluded that this is due to deviations of the windings from perfect helices. In particular, Karady et al. in [22] returned to using three straight, parallel conductors to model the twisted cables because this method provides conservative results at distances greater than 1 m .

It seems that predictions for twisted cables based on models using ideal helices underestimate the measured magnetic fields and it is also possible that these nonconservative results are due to the well-known first term of $B_{\varphi}$ from (1.b), equal to $\mu_{0} I / 2 \pi r$, implying that at large distances the helical line current acts as a straight line current, in particular when the helix is not perfect. In the exact twisted theory [14] this term is not considered for the composition of the total magnetic field generated by a three wire helix.

### 3.4 CONCLUSIONS

The magnetic field generated by an overhead or underground (the theoretical approach is the same) double-circuit twisted three-phase power cable line was treated following two different approaches both respecting the superposition principle; the exact approach which effectuates a complex vector sum of the two vectors of the magnetic field from each twisted three-phase single-circuit and the 'worst case'
approach which implies the algebraic sum of the two individual values of the magnetic field from each twisted three-phase single-circuit.

The exact vector sum has resulted a rather complex procedure since computational hypotheses were obligated to be set, such as the constancy of the pitch of the single helix and also the constancy of the perfect parallelism of the two helixes, and conventions that would facilitate the analysis since it is impossible to obtain exact information about the geometrical arrangement of the power line in practice were introduced. On the other hand, the 'worst case' assumption for the calculation of the magnetic field is a simple approximated calculus that yields a good interpretation of the reality. Additionally, the 'worst case' approach permits further approximation by using both the approximated formula from the literature and the simplified innovative one, in both levels of approximation. The simulation results have proven the efficiency of such an approach, since the expressions $B_{W C, \text { simp }}$ provide a small and positive relative percent error with respect to the exact vector calculus.

Comparing the calculations with results from measurements relevant to a threephase double-circuit twisted power cable line carried out in situ, it can be deducted that the theoretical approach of the double-circuit twisted cable line is correct and provides - especially when considering the difficult circumstances under which the experiment took place - magnetic field values near to the real ones, particularly for field points where the magnetic induction is large enough for being detected by measurement instruments.

Concluding, expression $B_{W C, \text { exact }}$ exhibits similar behavior to the expression $B_{\text {DoUBLE,calc }}$ and it is a simpler alternative to the exact calculus. Moreover, as expected, $B_{W C, \text { simp }}$ results more accurate than $B_{W C, l i t}$. So, $B_{W C, \text { simp }}$ is a "quick-and-easy" calculation tool that provides magnetic field results close to the actual ones where the magnetic field is large enough for being detected by measurement instruments. In this way, the calculation of the magnetic field generated by any multiple-circuit twisted three-phase power cable lines - standard lines for transporting renewable energy becomes possible.

## CHAPTER 4

## Multiple-Circuit Twisted Three - Phase Cable Lines

### 4.1 THEORY

In Chapter 3, the exact and approximated formulae from the literature along with the simplified innovative formulations were extended and applied to double-circuit twisted cable lines, under the following assumptions:

- the current terns carried by both cable lines are in-phase;
- since the medium is linear, at every field-point the magnetic field can be calculated composing linearly the two magnetic fields generated by each circuit separately (superposition principle);
- the exact treatment of the magnetic field from the double-circuit, $B_{D O U B L E}$, was obtained as the vector sum of the Cartesian components of the magnetic field from circuit 1 and 2 after formulae (3.1), (3.2), (3.3) and (3.4).
- for being conservative, a "worst case" estimate of the magnetic field from the double-circuit, $B_{W C, \text { exact }}$ (eq. 3.5) was obtained as the sum of the two rms values of total field $B$ from circuit 1 and 2 after the exact formula (eq. 2.12), $B_{1, \text { exact }}$ and $B_{2, \text { exact }}$, respectively.
- relationship (3.5) was further simplified by using the approximate formula of the literature (eq. 2.17) and the simplified innovative formula (eq. 2.23), in both levels of approximation, to express the rms values of $B$ from circuit 1 and 2, therefore obtaining the expressions $B_{W C, l i t}$ (eq. 3.6), $B_{W C, \text { simp } 1}$ (eq. 3.7.a) and $B_{W C, s i m p 2}$ (eq. 3.7.b).

Analogously to what done in Chapter 3 for double-circuit lines, the exact and the "worst case" approach is extended here to multiple-circuit twisted cable lines. The need for such an extension comes from the fact that not only single and double circuit three-phase cables are used for the connection of renewable sources to the grid, but also multiple-circuit twisted cable lines in underground configuration, especially when the generators (photovoltaic panels or wind generators) are numerous and/or dispersed over a vast surface, and nevertheless have all to be connected to one single substation.

### 4.1.1 Exact Vector Analysis

The magnetic field from amultiple-circuit, $B_{\text {MULTIPLE }}$, is given by the following formula:

$$
\begin{equation*}
B_{M U L T I P L E}=\sqrt{B_{x}^{2}+B_{y}^{2}+B_{z}^{2}} \tag{4.1}
\end{equation*}
$$

being

$$
\left\{\begin{array}{l}
B_{x}=B_{r 1} \cos \theta_{1}+B_{r 2} \cos \theta_{2}+\ldots+B_{r n} \cos \theta_{\mathrm{n}}+B_{\phi 1} \cos \varphi_{1}+B_{\phi 2} \cos \varphi_{2}+\ldots+B_{\phi n} \cos \varphi_{\mathrm{n}}  \tag{4.2}\\
B_{y}=B_{r 1} \sin \theta_{1}+B_{r 2} \sin \theta_{2}+\ldots+B_{r n} \sin \theta_{\mathrm{n}}+B_{\phi 1} \sin \varphi_{1}+B_{\phi 2} \sin \varphi_{2}+\ldots+B_{\phi n} \sin \varphi_{\mathrm{n}} \\
B_{z}=B_{z 1}+B_{z 2}+\ldots+B_{z n}
\end{array}\right.
$$

where $B_{r 1}, B_{r 2}, \ldots, B_{r n}, B_{\varphi 1}, B_{\varphi 2}, \ldots, B_{\varphi n}$ and $B_{z 1}, B_{z 2}, \ldots, B_{z n}$ are the exact values of the radial, azimuthal and axial components of the magnetic field from $n$ circuits.

Angles $\theta_{1}, \theta_{2}, \ldots, \theta_{n}$ and $\varphi_{1}, \varphi_{2}, \ldots, \varphi_{n}$ are the angles formed by the $x$-axis and the radial, $B_{r}$, and azimuthal components, $B_{\varphi}$, from the $n$ circuits, respectively (measured anticlockwise from the $x$-axis):

$$
\begin{align*}
& \left\{\begin{array}{l}
\theta_{1}=\arctan \left(y_{1} / x_{1}\right) \\
\theta_{2}=\arctan \left(y_{2} / x_{2}\right) \\
\ldots \\
\theta_{\mathrm{n}}=\arctan \left(y_{n} / x_{n}\right)
\end{array}\right.  \tag{4.3}\\
& \left\{\begin{array}{l}
\varphi_{1}=\theta_{1}+\pi / 2 \\
\varphi_{2}=\theta_{2}+\pi / 2 \\
\cdots \\
\varphi_{\mathrm{n}}=\theta_{\mathrm{n}}+\pi / 2
\end{array}\right. \tag{4.4}
\end{align*}
$$

and
where $\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right), \ldots,\left(x_{n}, y_{n}\right)$ are the coordinates of the $n$ cables in the plane $z=0$, where $\varphi=\Phi$.

The above exact vector approach has been converted into a calculation tool for the magnetic field from multiple-circuit twisted three-phase cable lines by implementing relationships (4.1) - (4.4), plus exact formulae (2.6) for $B_{r 1}, B_{r 2}, \ldots, B_{r n}, B_{\varphi 1}, B_{\varphi 2}, \ldots$, $B_{\varphi n}$ and $B_{z 1}, B_{z 2}, \ldots, B_{z n}$, respectively, in a script in Matlab ${ }^{\mathrm{TM}}$ environment. The value $\Phi_{1}=\Phi_{2}=\ldots=\Phi_{n}=60^{\circ}$ has been set and kept constant also in the twisted multiple-circuit case considering the results of the angle Фanalysis made in the previous Chapters.

### 4.1.2 "Worst Case" Approach

For being conservative, the "worst case" estimate of the magnetic field from the multiple-circuit is obtained as the sum of all rms values of total field $B$ from $n$ circuits after the exact formula (2.12). Also in the multiple-circuit case, $B_{W C, \text { exact }}$ has the following expression:

$$
\begin{equation*}
B_{W C, \text { exact }}=B_{1, \text { exact }}+B_{2, \text { exact }}+\ldots+B_{n, \text { exact }} \tag{4.5}
\end{equation*}
$$

Relationship (4.5) can be further simplified by using the approximate formula of the literature (eq. 2.17) and the simplified innovative formula (eq. 2.23), in both levels of approximation, to express the rms values of field $B$ from $n$ circuits and by composing them according to the "worst-case" approach. Also in the multiple-circuit case $B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ have the following expressions:

$$
\begin{equation*}
B_{W C, l i t}=B_{1, l i t}+B_{2, l i t}+\ldots+B_{n, l i t} \tag{4.6}
\end{equation*}
$$

and

$$
\begin{gather*}
B_{W C, \text { simp } 1}=B_{1, \text { simp } 1}+B_{2, \text { simp } 1}+\ldots+B_{n, \text { simp } 1}  \tag{4.7.a}\\
B_{W C, \text { simp } 2}=B_{1, \text { simp } 2}+B_{2, \text { simp } 2}+\ldots+B_{n, \text { simp } 2} \tag{4.7.b}
\end{gather*}
$$

### 4.2 APPLICATIONS - SIMULATION RESULTS

By implementing this theoretical approach in Matlab ${ }^{\mathrm{TM}}$ environment, for a triplecircuit twisted three-phase power line a script is obtained which - given the coordinates of the first, $\left(x_{1}, y_{1}\right)$, second, $\left(x_{2}, y_{2}\right)$, and third, $\left(x_{3}, y_{3}\right)$, twisted cable in a 2dimensional Cartesian reference system orthogonal to line axis, where line axis is located at $x=0$ and ground level is located at $y=0$ - calculates the value of the magnetic field generated by the triple-circuit twisted three-phase cable line along a horizontal line parallel to the to the ground and placed at a certain vertical distance $H$ from the cables. This script uses eqs. (4.1) - (4.4) for the exact calculus, $B_{\text {TRIPLE }}$, eqs. (4.5) - (4.7) for the "worst-case" calculus, $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$, plus the exact formula (2.12), the approximate formula of the literature (2.17) and the simplified formula (2.23) for $B_{1 \text { exact }}, B_{2 \text { exact, }} B_{3 \text { exact }}, B_{1 l i t}, B_{2 l i t}, B_{3 l i t}, B_{1 \text { simpl }}, B_{2 \text { simpl }}$, $B_{3 \text { simp } 1}$ and $B_{1 \text { simp } 2}, B_{2 \text { simp } 2}, B_{3 \text { simp } 2}$, respectively, and has been applied to some casestudies taking into account different cross-sections of different cable types and different geometric configurations. Also in these applicative case-studies the percent errors of $B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{T R I P L E}$ have been computed at the various field points.

### 4.2.1 Underground MV Triple-Circuit Three-Phase Twisted Cable Line

First Case: the three twisted cable circuits are ARE4H1RX cables [64], each with cross-section $3 \times 240 \mathrm{~mm}^{2}$, rated voltage 12 (phase-to-ground)/ 20 (phase-to-phase) kV , ampacity $I=423 \mathrm{~A}$, pitch $p=1.52 \mathrm{~m}$, radius $\alpha=0.022 \mathrm{~m}$. The coordinates of the three twisted cables are $(-0.5,-1.0),(0,-1.0),(0.5,-1.0)$, i.e. the triple-circuit twisted cable
line is laid underground at a burying depth $H=1.0 \mathrm{~m}$ below the soil, with the single circuits horizontally spaced of 0.5 m from each other. The field points lie on the soil along a straight horizontal coordinate axis $x$ whose vertical distance from the laying plane of the triple-circuit twisted cable line is $H=1.0 \mathrm{~m}$ and are spread from -1.5 m to 1.5 m at both ends of the central circuit - having coordinates $(0,-1.0)$ - that is taken as the line axis.The geometry of the line section is depicted in Fig. 4.1 and the results of the first case-study are plotted in Fig. 4.2, which displays the curves of $B_{\text {TRIPLE }}$, $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ vs. abscissa $x$, and in Table 4.1, which reports the relevant numerical values and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$.


Figure 4.1: Geometrical and electrical data concerning the first case-study.


Figure 4.2: $B_{\text {TRIPLE }}, B_{W C, \text { exact }}, B_{W C, \text { lit }}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2} v$ s. distance $x$ for the first case-study $\left(I_{1}=I_{2}=I_{3}=423 \mathrm{~A}\right)$.

Table 4.1: Numerical values of $B_{\text {TRIPLE }}, B_{W C, \text { exact, }}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, e x a c t,} B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the first case-study $\left(I_{1}=I_{2}=I_{3}=423 \mathrm{~A}\right)$.

| $\begin{gathered} \mathrm{x} \\ {[\mathrm{~m}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {TRIPLE }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\mathbf{B}_{\text {wC,exact }}$ <br> [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | B $_{\text {wC, lit }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\text {WC, simp1 }}$ <br> [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\overline{\mathbf{B}_{\mathrm{WC}, \text { simp2 }}}$ $[\mu \mathbf{T}]$ | $\begin{gathered} \hline \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.5 | 0.1360 | 0.1362 | 0.13 | 0.1218 | -10.47 | 0.1438 | 5.70 | 0.1463 | 7.54 |
| -1.4 | 0.1882 | 0.1885 | 0.17 | 0.1676 | -10.93 | 0.1993 | 5.90 | 0.2027 | 7.68 |
| -1.3 | 0.2566 | 0.2571 | 0.22 | 0.2273 | -11.39 | 0.2720 | 6.01 | 0.2764 | 7.72 |
| -1.2 | 0.3437 | 0.3447 | 0.28 | 0.3030 | -11.84 | 0.3645 | 6.04 | 0.3701 | 7.68 |
| -1.1 | 0.4514 | 0.4531 | 0.38 | 0.3960 | -12.27 | 0.4785 | 6.02 | 0.4857 | 7.60 |
| -1.0 | 0.5792 | 0.5821 | 0.50 | 0.5060 | -12.64 | 0.6138 | 5.98 | 0.6227 | 7.51 |
| -0.9 | 0.7243 | 0.7291 | 0.67 | 0.6307 | -12.92 | 0.7675 | 5.97 | 0.7783 | 7.46 |
| -0.8 | 0.8805 | 0.8884 | 0.89 | 0.7652 | -13.10 | 0.9336 | 6.03 | 0.9464 | 7.48 |
| -0.7 | 1.039 | 1.051 | 1.18 | 0.9025 | -13.14 | 1.103 | 6.19 | 1.118 | 7.61 |
| -0.6 | 1.190 | 1.208 | 1.52 | 1.035 | -13.04 | 1.267 | 6.46 | 1.284 | 7.87 |
| -0.5 | 1.324 | 1.350 | 1.91 | 1.154 | -12.83 | 1.415 | 6.82 | 1.433 | 8.23 |
| -0.4 | 1.436 | 1.469 | 2.31 | 1.256 | -12.54 | 1.540 | 7.24 | 1.5601 | 8.64 |
| -0.3 | 1.523 | 1.563 | 2.68 | 1.336 | -12.24 | 1.639 | 7.65 | 1.6604 | 9.05 |
| -0.2 | 1.584 | 1.631 | 2.97 | 1.394 | -11.98 | 1.710 | 7.98 | 1.7323 | 9.38 |
| -0.1 | 1.620 | 1.671 | 3.17 | 1.429 | -11.81 | 1.753 | 8.20 | 1.7756 | 9.61 |
| -0.0 | 1.631 | 1.685 | 3.27 | 1.440 | -11.72 | 1.767 | 8.32 | 1.7900 | 9.73 |
| 0.1 | 1.618 | 1.671 | 3.29 | 1.429 | -11.71 | 1.753 | 8.32 | 1.7756 | 9.73 |
| 0.2 | 1.580 | 1.631 | 3.20 | 1.394 | -11.79 | 1.710 | 8.22 | 1.7323 | 9.63 |
| 0.3 | 1.518 | 1.563 | 3.01 | 1.336 | -11.95 | 1.639 | 8.00 | 1.6604 | 9.40 |
| 0.4 | 1.430 | 1.469 | 2.73 | 1.256 | -12.18 | 1.540 | 7.68 | 1.5601 | 9.08 |
| 0.5 | 1.318 | 1.350 | 2.37 | 1.154 | -12.44 | 1.415 | 7.30 | 1.4332 | 8.71 |
| 0.6 | 1.185 | 1.208 | 1.98 | 1.035 | -12.65 | 1.267 | 6.94 | 1.2836 | 8.36 |
| 0.7 | 1.035 | 1.051 | 1.61 | 0.9025 | -12.76 | 1.103 | 6.65 | 1.1182 | 8.08 |
| 0.8 | 0.8771 | 0.8884 | 1.29 | 0.7652 | -12.76 | 0.9336 | 6.44 | 0.9464 | 7.90 |
| 0.9 | 0.7218 | 0.7291 | 1.01 | 0.6307 | -12.63 | 0.7675 | 6.33 | 0.7783 | 7.82 |
| 1.0 | 0.5775 | 0.5821 | 0.79 | 0.5060 | -12.38 | 0.6138 | 6.29 | 0.6227 | 7.82 |
| 1.1 | 0.4503 | 0.4531 | 0.62 | 0.3960 | -12.06 | 0.4785 | 6.27 | 0.4857 | 7.86 |
| 1.2 | 0.3431 | 0.3447 | 0.48 | 0.3030 | -11.67 | 0.3645 | 6.25 | 0.3701 | 7.90 |
| 1.3 | 0.2561 | 0.2571 | 0.38 | 0.2273 | -11.25 | 0.2720 | 6.19 | 0.2764 | 7.89 |
| 1.4 | 0.1880 | 0.1885 | 0.30 | 0.1676 | -10.81 | 0.1993 | 6.05 | 0.2027 | 7.83 |
| 1.5 | 0.1359 | 0.1362 | 0.24 | 0.1218 | -10.37 | 0.1438 | 5.82 | 0.1463 | 7.67 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $3.5 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simp } I}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 8 \%$ for the $1^{\text {st }}$ approximation level and maximum $\sim 10 \%$ for the $2^{\text {nd }}$ approximation level) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6), arrive at $\sim-13 \%$ being always negative.

Second Case: same as the first case, apart that the central cable has a larger crosssection than the other two cables, namely $3 \times 300 \mathrm{~mm}^{2}$, thus larger ampacity $I=478 \mathrm{~A}$, pitch $p=1.66 \mathrm{~m}$ and radius $\alpha=0.024 \mathrm{~m}$ (see Fig. 4.3). The results of the second case simulations are shown in Fig. 4.4 and Table 4.2.


Figure 4.3: Geometrical and electrical data concerning the second case-study.


Figure 4.4: $B_{T R I P L E}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simpl } 1}, B_{W C, \text { simp } 2} v s$. distance $x$ for the second case-study $\left(I_{1}=I_{3}=423 \mathrm{~A}\right.$ and $\left.I_{2}=478 \mathrm{~A}\right)$.

Table 4.2: Numerical values of $B_{\text {TRIPLE, }} B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the second case-study $\left(I_{1}=I_{3}=423 \mathrm{~A}\right.$ and $\left.I_{2}=478 \mathrm{~A}\right)$.

| $\begin{gathered} \mathrm{x} \\ {[\mathrm{~m}]} \end{gathered}$ | $\begin{gathered} \hline \mathbf{B}_{\text {TRIPLE }} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | $\mathbf{B}_{\text {WC,exact }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \mathrm{lit}} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { simp1 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { simp2 } 2} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.5 | 0.1564 | 0.1567 | 0.19 | 0.1402 | -10.36 | 0.1650 | 5.51 | 0.1664 | 6.40 |
| -1.4 | 0.2160 | 0.2165 | 0.24 | 0.1927 | -10.82 | 0.2286 | 5.81 | 0.2304 | 6.65 |
| -1.3 | 0.2942 | 0.2951 | 0.30 | 0.2610 | -11.28 | 0.3119 | 6.00 | 0.3142 | 6.79 |
| -1.2 | 0.3943 | 0.3958 | 0.38 | 0.3480 | -11.73 | 0.4184 | 6.11 | 0.4213 | 6.85 |
| -1.1 | 0.5185 | 0.5211 | 0.49 | 0.4555 | -12.16 | 0.5505 | 6.16 | 0.5540 | 6.84 |
| -1.0 | 0.6674 | 0.6716 | 0.63 | 0.5837 | -12.54 | 0.7087 | 6.19 | 0.7129 | 6.81 |
| -0.9 | 0.8387 | 0.8455 | 0.81 | 0.7308 | -12.86 | 0.8908 | 6.22 | 0.8956 | 6.79 |
| -0.8 | 1.027 | 1.037 | 1.03 | 0.8922 | -13.10 | 1.091 | 6.30 | 1.096 | 6.80 |
| -0.7 | 1.223 | 1.238 | 1.30 | 1.061 | -13.24 | 1.301 | 6.43 | 1.307 | 6.87 |
| -0.6 | 1.416 | 1.439 | 1.60 | 1.228 | -13.29 | 1.510 | 6.62 | 1.515 | 7.00 |
| -0.5 | 1.597 | 1.627 | 1.90 | 1.385 | -13.27 | 1.706 | 6.85 | 1.711 | 7.16 |
| -0.4 | 1.756 | 1.794 | 2.17 | 1.523 | -13.24 | 1.880 | 7.06 | 1.884 | 7.32 |
| -0.3 | 1.887 | 1.932 | 2.39 | 1.638 | -13.22 | 2.024 | 7.23 | 2.027 | 7.44 |
| -0.2 | 1.986 | 2.036 | 2.52 | 1.723 | -13.23 | 2.131 | 7.32 | 2.135 | 7.49 |
| -0.1 | 2.047 | 2.100 | 2.59 | 1.776 | -13.25 | 2.198 | 7.36 | 2.201 | 7.50 |
| -0.0 | 2.068 | 2.122 | 2.61 | 1.794 | -13.26 | 2.221 | 7.38 | 2.223 | 7.51 |
| 0.1 | 2.047 | 2.100 | 2.62 | 1.776 | -13.23 | 2.198 | 7.39 | 2.201 | 7.53 |
| 0.2 | 1.984 | 2.036 | 2.59 | 1.723 | -13.17 | 2.131 | 7.40 | 2.135 | 7.57 |
| 0.3 | 1.885 | 1.932 | 2.52 | 1.638 | -13.11 | 2.024 | 7.37 | 2.027 | 7.58 |
| 0.4 | 1.752 | 1.794 | 2.37 | 1.523 | -13.07 | 1.880 | 7.27 | 1.884 | 7.53 |
| 0.5 | 1.593 | 1.627 | 2.16 | 1.385 | -13.05 | 1.706 | 7.12 | 1.711 | 7.44 |
| 0.6 | 1.412 | 1.439 | 1.89 | 1.228 | -13.03 | 1.510 | 6.93 | 1.515 | 7.31 |
| 0.7 | 1.219 | 1.238 | 1.61 | 1.061 | -12.97 | 1.301 | 6.75 | 1.307 | 7.20 |
| 0.8 | 1.024 | 1.037 | 1.33 | 0.8922 | -12.84 | 1.091 | 6.61 | 1.096 | 7.12 |
| 0.9 | 0.8364 | 0.8455 | 1.09 | 0.7308 | -12.62 | 0.8908 | 6.51 | 0.8956 | 7.08 |
| 1.0 | 0.6658 | 0.6716 | 0.87 | 0.5837 | -12.33 | 0.7087 | 6.45 | 0.7129 | 7.07 |
| 1.1 | 0.5174 | 0.5211 | 0.70 | 0.4555 | -11.97 | 0.5505 | 6.38 | 0.5540 | 7.07 |
| 1.2 | 0.3936 | 0.3958 | 0.56 | 0.3480 | -11.57 | 0.4184 | 6.30 | 0.4213 | 7.04 |
| 1.3 | 0.2938 | 0.2951 | 0.45 | 0.2610 | -11.14 | 0.3119 | 6.16 | 0.3142 | 6.95 |
| 1.4 | 0.2158 | 0.2165 | 0.36 | 0.1927 | -10.71 | 0.2286 | 5.94 | 0.2304 | 6.79 |
| 1.5 | 0.1563 | 0.1567 | 0.30 | 0.1402 | -10.27 | 0.1650 | 5.62 | 0.1664 | 6.52 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $3 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simp } 1}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 7.5 \%$ for both approximation levels) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6), arrive at $\sim-13 \%$ being always negative.

Third Case: same overall arrangement as in the second case, apart that the central cable, with the larger cross-section, is at a burying depth $H_{2}=1.5 \mathrm{~m}$, thus the field points lie on a straight horizontal line whose vertical distance from the central cable is 1.5 m , while their vertical distance from the other two cables is still 1.0 m (see Fig. 4.5). Figure 4.6 and Table 4.3 show the results of the third case simulation.


Figure 4.5: Geometrical and electrical data concerning the third case-study.


Figure 4.6: $B_{\text {TRIPLE }}, B_{W C, \text { exact, }} B_{W C, \text { li, }}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance x for the third case-study $\left(I_{1}=I_{3}=423 \mathrm{~A}\right.$ and $\left.I_{2}=478 \mathrm{~A}\right)$.

Table 4.3: Numerical values of $B_{\text {TRIPLE, }}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the third case-study $\left(I_{1}=I_{3}=423 A\right.$ and $\left.I_{2}=478 \mathrm{~A}\right)$.

| $\begin{gathered} \mathbf{x} \\ {[\mathbf{m}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {TRIPLE }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $B_{\text {wC,exact }}$ <br> [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \mathrm{it}} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\mathbf{B}_{\mathrm{WC}, \text { simp1 }}$ $[\mu \mathrm{T}]$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { simp2 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.5 | 0.1274 | 0.1274 | 0.05 | 0.1139 | -10.61 | 0.1343 | 5.41 | 0.1361 | 6.88 |
| -1.4 | 0.1745 | 0.1746 | 0.06 | 0.1551 | -11.12 | 0.1842 | 5.59 | 0.1867 | 7.00 |
| -1.3 | 0.2352 | 0.2354 | 0.08 | 0.2078 | -11.63 | 0.2485 | 5.66 | 0.2517 | 7.02 |
| -1.2 | 0.3111 | 0.3114 | 0.11 | 0.2733 | -12.14 | 0.3286 | 5.62 | 0.3327 | 6.94 |
| -1.1 | 0.4025 | 0.4031 | 0.15 | 0.3517 | -12.62 | 0.4248 | 5.53 | 0.4298 | 6.79 |
| -1.0 | 0.5076 | 0.5087 | 0.22 | 0.4413 | -13.06 | 0.5350 | 5.40 | 0.5412 | 6.62 |
| -0.9 | 0.6215 | 0.6235 | 0.32 | 0.5381 | -13.42 | 0.6544 | 5.30 | 0.6617 | 6.47 |
| -0.8 | 0.7363 | 0.7398 | 0.47 | 0.6359 | -13.64 | 0.7751 | 5.27 | 0.7835 | 6.40 |
| -0.7 | 0.8419 | 0.8478 | 0.70 | 0.7266 | -13.69 | 0.8872 | 5.37 | 0.8965 | 6.48 |
| -0.6 | 0.9278 | 0.9375 | 1.05 | 0.8024 | -13.52 | 0.9805 | 5.68 | 0.9906 | 6.77 |
| -0.5 | 0.9863 | 1.002 | 1.55 | 0.8574 | -13.06 | 1.048 | 6.26 | 1.059 | 7.33 |
| -0.4 | 1.015 | 1.038 | 2.24 | 0.8902 | -12.32 | 1.088 | 7.12 | 1.098 | 8.19 |
| -0.3 | 1.019 | 1.051 | 3.11 | 0.9037 | -11.32 | 1.103 | 8.24 | 1.114 | 9.32 |
| -0.2 | 1.008 | 1.049 | 4.07 | 0.9048 | -10.22 | 1.103 | 9.48 | 1.114 | 10.56 |
| -0.1 | 0.9936 | 1.042 | 4.90 | 0.9013 | -9.29 | 1.098 | 10.53 | 1.109 | 11.62 |
| -0.0 | 0.9864 | 1.039 | 5.34 | 0.8993 | -8.83 | 1.096 | 11.06 | 1.106 | 12.16 |
| 0.1 | 0.9905 | 1.042 | 5.22 | 0.9013 | -9.01 | 1.098 | 10.87 | 1.109 | 11.97 |
| 0.2 | 1.003 | 1.049 | 4.62 | 0.9048 | -9.75 | 1.103 | 10.05 | 1.114 | 11.14 |
| 0.3 | 1.013 | 1.051 | 3.76 | 0.9037 | -10.77 | 1.103 | 8.92 | 1.114 | 10.01 |
| 0.4 | 1.009 | 1.038 | 2.88 | 0.8902 | -11.77 | 1.088 | 7.79 | 1.098 | 8.87 |
| 0.5 | 0.9808 | 1.002 | 2.12 | 0.8574 | -12.57 | 1.048 | 6.85 | 1.059 | 7.94 |
| 0.6 | 0.9234 | 0.9375 | 1.53 | 0.8024 | -13.11 | 0.9805 | 6.18 | 0.9906 | 7.28 |
| 0.7 | 0.8387 | 0.8478 | 1.09 | 0.7266 | -13.36 | 0.8872 | 5.78 | 0.8965 | 6.89 |
| 0.8 | 0.7341 | 0.7398 | 0.77 | 0.6359 | -13.38 | 0.7751 | 5.58 | 0.7835 | 6.73 |
| 0.9 | 0.6200 | 0.6235 | 0.55 | 0.5381 | -13.21 | 0.6544 | 5.54 | 0.6617 | 6.72 |
| 1.0 | 0.5067 | 0.5087 | 0.40 | 0.4413 | -12.91 | 0.5350 | 5.59 | 0.5412 | 6.81 |
| 1.1 | 0.4020 | 0.4031 | 0.29 | 0.3517 | -12.51 | 0.4248 | 5.67 | 0.4298 | 6.93 |
| 1.2 | 0.3108 | 0.3114 | 0.21 | 0.2733 | -12.05 | 0.3286 | 5.73 | 0.3327 | 7.04 |
| 1.3 | 0.2350 | 0.2354 | 0.15 | 0.2078 | -11.56 | 0.2485 | 5.73 | 0.2517 | 7.10 |
| 1.4 | 0.1744 | 0.1746 | 0.11 | 0.1551 | -11.07 | 0.1842 | 5.65 | 0.1867 | 7.06 |
| 1.5 | 0.1273 | 0.1274 | 0.09 | 0.1139 | -10.58 | 0.1343 | 5.45 | 0.1361 | 6.92 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $5 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simpl }}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 11 \%$ for the $1^{\text {st }}$ level of approximation and maximum $\sim 12 \%$ for the $2^{\text {nd }}$ level of approximation) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6), arrive at $\sim-13.5 \%$ being always negative.

### 4.2.2 Overhead MV Triple-Circuit Three-Phase Twisted Cable Line

First Case: the three twisted cable circuits are ARG7H1RX cables [63], each with cross-section $3 \times 120 \mathrm{~mm}^{2}$, rated voltage 12 (phase-to-ground)/ 20 (phase-to-phase) kV , ampacity $I=280 \mathrm{~A}$, pitch $p=1.37 \mathrm{~m}$, radius $\alpha=0.020 \mathrm{~m}$. The coordinates of the three twisted cables are $(-0.4,10),(0,10),(0.4,10)$, i.e. the triple-circuit twisted cable
line is situated at a height of 10 m above the ground, with the single circuits horizontally spaced of 0.4 m from each other. The field points lie on a straight horizontal line whose vertical distance from the laying plane of the triple-circuit twisted cable line is $H=0.5 \mathrm{~m}$ and are spread from -1.5 m to 1.5 m at both ends of the central circuit - having coordinates $(0,10)$ - that is taken as the line axis. The geometry of the line section is depicted in Fig. 4.7 and the results of the first casestudy are plotted in Fig. 4.8, which displays the curves of $B_{\text {TRIPLE }}, B_{W C, \text { exact }}, B_{W C, l i t}$, $B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp2 }}$ vs. abscissa $x$, and in Table 4.4, which reports the relevant numerical values and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$.


Figure 4.7: Geometrical and electrical data concerning the first case-study.


Figure 4.8: $B_{\text {TRIPLE }}, B_{W C, \text { exact, }} B_{W C, \text { lit }} B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ for the first case-study ( $\left.I_{1}=I_{2}=I_{3}=280 \mathrm{~A}\right)$.

Table 4.4: Numerical values of $B_{\text {TRIPLE, }}, B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the first case-study $\left(I_{1}=I_{2}=I_{3}=280 \mathrm{~A}\right)$.

| $\mathbf{x}$ |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[\mathbf{m}]$ | $\mathbf{B}_{\text {TRIPLE }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | $\mathbf{B}_{\text {WC,exact }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {WC,lit }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\text {WC,simp1 }}$ <br> $[\boldsymbol{\mu} \mathbf{T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {WC,simp } 2}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ |
| -1.5 | 0.1387 | 0.1388 | 0.08 | 0.1233 | -11.08 | 0.1469 | 5.90 | 0.1512 | 8.99 |
| -1.4 | 0.2206 | 0.2208 | 0.11 | 0.1944 | -11.89 | 0.2330 | 5.63 | 0.2393 | 8.49 |
| -1.3 | 0.3495 | 0.3500 | 0.15 | 0.3047 | -12.82 | 0.3674 | 5.15 | 0.3766 | 7.76 |
| -1.2 | 0.5506 | 0.5517 | 0.20 | 0.4742 | -13.87 | 0.5751 | 4.46 | 0.5882 | 6.84 |
| -1.1 | 0.8604 | 0.8628 | 0.28 | 0.7308 | -15.05 | 0.8913 | 3.60 | 0.9099 | 5.76 |
| -1.0 | 1.328 | 1.334 | 0.40 | 1.111 | -16.37 | 1.364 | 2.65 | 1.389 | 4.58 |
| -0.9 | 2.015 | 2.027 | 0.59 | 1.656 | -17.80 | 2.050 | 1.72 | 2.084 | 3.43 |
| -0.8 | 2.976 | 3.002 | 0.88 | 2.403 | -19.26 | 3.005 | 0.97 | 3.049 | 2.48 |
| -0.7 | 4.226 | 4.283 | 1.34 | 3.356 | -20.60 | 4.251 | 0.60 | 4.307 | 1.93 |
| -0.6 | 5.679 | 5.797 | 2.08 | 4.457 | -21.52 | 5.724 | 0.79 | 5.791 | 1.98 |
| -0.5 | 7.108 | 7.335 | 3.20 | 5.566 | -21.69 | 7.223 | 1.63 | 7.301 | 2.72 |
| -0.4 | 8.236 | 8.622 | 4.69 | 6.512 | -20.94 | 8.479 | 2.94 | 8.567 | 4.01 |
| -0.3 | 8.965 | 9.519 | 6.18 | 7.202 | -19.66 | 9.349 | 4.29 | 9.447 | 5.37 |
| -0.2 | 9.442 | 10.10 | 6.98 | 7.664 | -18.83 | 9.917 | 5.03 | 10.02 | 6.14 |
| -0.1 | 9.797 | 10.48 | 6.93 | 7.950 | -18.85 | 10.29 | 5.04 | 10.40 | 6.15 |
| -0.0 | 9.950 | 10.62 | 6.71 | 8.053 | -19.07 | 10.44 | 4.88 | 10.55 | 5.98 |
| 0.1 | 9.822 | 10.48 | 6.66 | 7.950 | -19.05 | 10.29 | 4.78 | 10.40 | 5.89 |
| 0.2 | 9.510 | 10.10 | 6.22 | 7.664 | -19.41 | 9.917 | 4.28 | 10.02 | 5.38 |
| 0.3 | 9.071 | 9.519 | 4.93 | 7.202 | -20.60 | 9.349 | 3.06 | 9.447 | 4.14 |
| 0.4 | 8.347 | 8.622 | 3.30 | 6.512 | -21.99 | 8.479 | 1.58 | 8.567 | 2.63 |
| 0.5 | 7.193 | 7.335 | 1.97 | 5.566 | -22.62 | 7.223 | 0.42 | 7.301 | 1.50 |
| 0.6 | 5.732 | 5.797 | 1.12 | 4.457 | -22.26 | 5.724 | -0.15 | 5.791 | 1.03 |
| 0.7 | 4.255 | 4.283 | 0.64 | 3.356 | -21.14 | 4.251 | -0.10 | 4.307 | 1.22 |
| 0.8 | 2.991 | 3.002 | 0.37 | 2.403 | -19.67 | 3.005 | 0.46 | 3.049 | 1.96 |
| 0.9 | 2.022 | 2.027 | 0.23 | 1.656 | -18.10 | 2.050 | 1.35 | 2.084 | 3.06 |
| 1.0 | 1.332 | 1.334 | 0.14 | 1.111 | -16.59 | 1.364 | 2.39 | 1.389 | 4.31 |
| 1.1 | 0.8620 | 0.8628 | 0.09 | 0.7308 | -15.21 | 0.8913 | 3.41 | 0.9099 | 5.56 |
| 1.2 | 0.5513 | 0.5517 | 0.06 | 0.4742 | -13.99 | 0.5751 | 4.31 | 0.5882 | 6.69 |
| 1.3 | 0.3498 | 0.3500 | 0.04 | 0.3047 | -12.91 | 0.3674 | 5.03 | 0.3766 | 7.65 |
| 1.4 | 0.2208 | 0.2208 | 0.03 | 0.1944 | -11.96 | 0.2330 | 5.55 | 0.2393 | 8.40 |
| 1.5 | 0.1388 | 0.1388 | 0.02 | 0.1233 | -11.13 | 0.1469 | 5.84 | 0.1512 | 8.92 |
|  |  |  |  |  |  |  |  |  |  |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $7 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simp } 1}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 5 \%$ for the $1^{\text {st }}$ level of approximation and maximum $\sim 6 \%$ for the $2^{\text {nd }}$ level of approximation) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6), arrive at $\sim-23 \%$ being always negative.

Second Case: same as the first case, apart thatthe field points lie on a straight horizontal line whose vertical distance from the laying plane of the triple-circuit twisted cable line is, now, $H=1.0 \mathrm{~m}$ (see Fig. 4.9).The results of the second case simulations are shown in Fig. 4.10 and Table 4.5.


Figure 4.9: Geometrical and electrical data concerning the second case-study.


Figure 4.10: $B_{T R I P L E}, B_{W C, \text { exact, }} B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ for the second casestudy $\left(I_{1}=I_{2}=I_{3}=280 \mathrm{~A}\right)$.

Table 4.5: Numerical values of $B_{\text {TRIPLE }}, B_{W C, \text { exact, }} B_{W C, \text { li, }} B_{W C, \text { simpl }}$, $B_{W C, \text { simp } 2} v s$. distance $x$ and percent errors of $B_{W C, \text { exact, }} B_{W C, l i t}, B_{W C, \text { simp } 1}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the second case-study ( $\left.I_{1}=I_{2}=I_{3}=280 \mathrm{~A}\right)$.

| $\begin{gathered} \mathbf{x} \\ {[\mathbf{m}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {TRIPLE }} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { exact }} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | error <br> [\%] | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \mathrm{lit}} \\ {[\boldsymbol{\mu T}]} \end{gathered}$ | error $[\%]$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { simp1 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | error $[\%]$ | $\begin{gathered} \mathbf{B}_{\text {WC,simp2 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.5 | 0.0361 | 0.0362 | 0.17 | 0.0329 | -8.98 | 0.0381 | 5.53 | 0.0395 | 9.31 |
| -1.4 | 0.0524 | 0.0525 | 0.21 | 0.0475 | -9.37 | 0.0554 | 5.88 | 0.0573 | 9.49 |
| -1.3 | 0.0748 | 0.0750 | 0.26 | 0.0675 | -9.77 | 0.0794 | 6.11 | 0.0820 | 9.57 |
| -1.2 | 0.1052 | 0.1056 | 0.33 | 0.0945 | -10.17 | 0.1118 | 6.26 | 0.1153 | 9.56 |
| -1.1 | 0.1452 | 0.1458 | 0.41 | 0.1299 | -10.56 | 0.1544 | 6.32 | 0.1590 | 9.49 |
| -1.0 | 0.1960 | 0.1971 | 0.53 | 0.1746 | -10.92 | 0.2085 | 6.35 | 0.2144 | 9.37 |
| -0.9 | 0.2581 | 0.2598 | 0.67 | 0.2291 | -11.23 | 0.2745 | 6.35 | 0.2820 | 9.26 |
| -0.8 | 0.3305 | 0.3333 | 0.85 | 0.2925 | -11.48 | 0.3515 | 6.36 | 0.3608 | 9.17 |
| -0.7 | 0.4104 | 0.4148 | 1.06 | 0.3627 | -11.64 | 0.4368 | 6.42 | 0.4480 | 9.15 |
| -0.6 | 0.4936 | 0.5000 | 1.31 | 0.4358 | -11.70 | 0.5258 | 6.54 | 0.5390 | 9.20 |
| -0.5 | 0.5744 | 0.5835 | 1.58 | 0.5073 | -11.68 | 0.6130 | 6.71 | 0.6280 | 9.33 |
| -0.4 | 0.6475 | 0.6595 | 1.85 | 0.5725 | -11.59 | 0.6924 | 6.93 | 0.7091 | 9.51 |
| -0.3 | 0.7082 | 0.7230 | 2.08 | 0.6270 | -11.47 | 0.7588 | 7.14 | 0.7770 | 9.71 |
| -0.2 | 0.7534 | 0.7705 | 2.26 | 0.6678 | -11.36 | 0.8085 | 7.31 | 0.8278 | 9.87 |
| -0.1 | 0.7813 | 0.7997 | 2.35 | 0.6930 | -11.31 | 0.8391 | 7.40 | 0.8591 | 9.95 |
| -0.0 | 0.7910 | 0.8095 | 2.34 | 0.7015 | -11.32 | 0.8495 | 7.40 | 0.8697 | 9.95 |
| 0.1 | 0.7821 | 0.7997 | 2.25 | 0.6930 | -11.39 | 0.8391 | 7.29 | 0.8591 | 9.85 |
| 0.2 | 0.7549 | 0.7705 | 2.07 | 0.6678 | -11.53 | 0.8085 | 7.11 | 0.8278 | 9.66 |
| 0.3 | 0.7101 | 0.7230 | 1.82 | 0.6270 | -11.70 | 0.7588 | 6.86 | 0.7770 | 9.42 |
| 0.4 | 0.6495 | 0.6595 | 1.54 | 0.5725 | -11.86 | 0.6924 | 6.60 | 0.7091 | 9.18 |
| 0.5 | 0.5763 | 0.5835 | 1.25 | 0.5073 | -11.96 | 0.6130 | 6.37 | 0.6280 | 8.98 |
| 0.6 | 0.4951 | 0.5000 | 0.99 | 0.4358 | -11.98 | 0.5258 | 6.20 | 0.5390 | 8.85 |
| 0.7 | 0.4116 | 0.4148 | 0.77 | 0.3627 | -11.90 | 0.4368 | 6.11 | 0.4480 | 8.83 |
| 0.8 | 0.3313 | 0.3333 | 0.58 | 0.2925 | -11.71 | 0.3515 | 6.09 | 0.3608 | 8.89 |
| 0.9 | 0.2587 | 0.2598 | 0.44 | 0.2291 | -11.43 | 0.2745 | 6.11 | 0.2820 | 9.01 |
| 1.0 | 0.1964 | 0.1971 | 0.34 | 0.1746 | -11.09 | 0.2085 | 6.14 | 0.2144 | 9.17 |
| 1.1 | 0.1454 | 0.1458 | 0.25 | 0.1299 | -10.70 | 0.1544 | 6.16 | 0.1590 | 9.31 |
| 1.2 | 0.1054 | 0.1056 | 0.19 | 0.0945 | -10.29 | 0.1118 | 6.12 | 0.1153 | 9.42 |
| 1.3 | 0.0749 | 0.0750 | 0.15 | 0.0675 | -9.87 | 0.0794 | 6.00 | 0.0820 | 9.45 |
| 1.4 | 0.0524 | 0.0525 | 0.12 | 0.0475 | -9.45 | 0.0554 | 5.78 | 0.0573 | 9.39 |
| 1.5 | 0.0362 | 0.0362 | 0.09 | 0.0329 | -9.05 | 0.0381 | 5.45 | 0.0395 | 9.23 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $2.5 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simpl }}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 7.5 \%$ for the $1^{\text {st }}$ level of approximation and maximum $\sim 10 \%$ for the $2^{\text {nd }}$ level of approximation) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6 ), arrive at $\sim-11 \%$ being always negative.

Third Case: same overall arrangement as in the second case, apart that the central cable has a larger cross-section than the other two cables, namely $3 \times 185 \mathrm{~mm}^{2}$, thus larger ampacity $I=360 \mathrm{~A}$, pitch $p=1.5 \mathrm{~m}$ and radius $\alpha=0.022 \mathrm{~m}$ (see Fig. 4.11). Figure 4.12 and Table 4.6 show the results of the third case simulation.


Figure 4.11: Geometrical and electrical data concerning the third case-study.


Figure 4.12: $B_{\text {TRIPLE, }}, B_{W C, \text { exact, }}, B_{W C, \text { lit }} B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ for the third case-study $\left(I_{1}=I_{3}=280 \mathrm{~A}\right.$ and $\left.I_{2}=360 \mathrm{~A}\right)$.

Table 4.6: Numerical values of $B_{T R I P L E}, B_{W C, e x a c t,} B_{W C, \text { lii, }} B_{W C, \text { simpl }}, B_{W C, \text { simp } 2} v$ s. distance $x$ and percent errors of $B_{W C, \text { exact, }} B_{W C, \text { lit }}, B_{W C, \text { simpl }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the third case-study $\left(I_{1}=I_{3}=280\right.$ A and $\left.I_{2}=360 \mathrm{~A}\right)$.

| $\mathbf{x}$ |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $[\mathbf{m}]$ | $\mathbf{B}_{\text {TRIPLE }}$ <br> $[\boldsymbol{\mu T}]$ | $\mathbf{B}_{\text {WC,exact }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\text {WC,lit }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\%]$ | $\mathbf{B}_{\text {WC,simp1 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\boldsymbol{\%}]$ | $\mathbf{B}_{\text {WC,simp2 }}$ <br> $[\boldsymbol{\mu T}]$ | error <br> $[\mathbf{\%}]$ |
| -1.5 | 0.0455 | 0.0456 | 0.19 | 0.0414 | -8.94 | 0.0479 | 5.22 | 0.0494 | 8.54 |
| -1.4 | 0.0657 | 0.0658 | 0.24 | 0.0595 | -9.33 | 0.0694 | 5.66 | 0.0715 | 8.84 |
| -1.3 | 0.0936 | 0.0938 | 0.29 | 0.0844 | -9.74 | 0.0992 | 5.99 | 0.1020 | 9.02 |
| -1.2 | 0.1314 | 0.1318 | 0.36 | 0.1180 | -10.15 | 0.1395 | 6.21 | 0.1433 | 9.11 |
| -1.1 | 0.1813 | 0.1821 | 0.45 | 0.1621 | -10.55 | 0.1928 | 6.35 | 0.1978 | 9.11 |
| -1.0 | 0.2452 | 0.2465 | 0.55 | 0.2183 | -10.94 | 0.2609 | 6.42 | 0.2674 | 9.06 |
| -0.9 | 0.3241 | 0.3263 | 0.68 | 0.2875 | -11.31 | 0.3451 | 6.45 | 0.3532 | 8.97 |
| -0.8 | 0.4178 | 0.4212 | 0.83 | 0.3692 | -11.63 | 0.4447 | 6.46 | 0.4548 | 8.87 |
| -0.7 | 0.5236 | 0.5288 | 1.00 | 0.4612 | -11.91 | 0.5575 | 6.47 | 0.5696 | 8.79 |
| -0.6 | 0.6370 | 0.6445 | 1.17 | 0.5597 | -12.13 | 0.6783 | 6.49 | 0.6926 | 8.72 |
| -0.5 | 0.7516 | 0.7617 | 1.35 | 0.6591 | -12.30 | 0.8006 | 6.52 | 0.8168 | 8.68 |
| -0.4 | 0.8598 | 0.8727 | 1.50 | 0.7528 | -12.44 | 0.9161 | 6.54 | 0.9342 | 8.65 |
| -0.3 | 0.9540 | 0.9694 | 1.61 | 0.8342 | -12.56 | 1.017 | 6.56 | 1.036 | 8.62 |
| -0.2 | 1.027 | 1.045 | 1.69 | 0.8973 | -12.65 | 1.095 | 6.55 | 1.116 | 8.59 |
| -0.1 | 1.074 | 1.092 | 1.72 | 0.9372 | -12.72 | 1.144 | 6.54 | 1.166 | 8.55 |
| -0.0 | 1.090 | 1.109 | 1.72 | 0.9509 | -12.76 | 1.161 | 6.52 | 1.183 | 8.53 |
| 0.1 | 1.074 | 1.092 | 1.69 | 0.9372 | -12.75 | 1.144 | 6.51 | 1.166 | 8.52 |
| 0.2 | 1.028 | 1.045 | 1.62 | 0.8973 | -12.71 | 1.095 | 6.49 | 1.116 | 8.52 |
| 0.3 | 0.9550 | 0.9694 | 1.51 | 0.8342 | -12.65 | 1.017 | 6.45 | 1.036 | 8.51 |
| 0.4 | 0.8610 | 0.8727 | 1.36 | 0.7528 | -12.56 | 0.9161 | 6.40 | 0.9342 | 8.50 |
| 0.5 | 0.7528 | 0.7617 | 1.18 | 0.6591 | -12.45 | 0.8006 | 6.34 | 0.8168 | 8.50 |
| 0.6 | 0.6382 | 0.6445 | 0.99 | 0.5597 | -12.29 | 0.6783 | 6.30 | 0.6926 | 8.52 |
| 0.7 | 0.5245 | 0.5288 | 0.81 | 0.4612 | -12.07 | 0.5575 | 6.27 | 0.5696 | 8.58 |
| 0.8 | 0.4185 | 0.4212 | 0.65 | 0.3692 | -11.79 | 0.4447 | 6.27 | 0.4548 | 8.68 |
| 0.9 | 0.3247 | 0.3263 | 0.52 | 0.2875 | -11.45 | 0.3451 | 6.28 | 0.3532 | 8.79 |
| 1.0 | 0.2455 | 0.2465 | 0.41 | 0.2183 | -11.07 | 0.2609 | 6.27 | 0.2674 | 8.90 |
| 1.1 | 0.1815 | 0.1821 | 0.32 | 0.1621 | -10.66 | 0.1928 | 6.22 | 0.1978 | 8.98 |
| 1.2 | 0.1315 | 0.1318 | 0.26 | 0.1180 | -10.24 | 0.1395 | 6.10 | 0.1433 | 9.00 |
| 1.3 | 0.0936 | 0.0938 | 0.20 | 0.0844 | -9.82 | 0.0992 | 5.90 | 0.1020 | 8.93 |
| 1.4 | 0.0657 | 0.0658 | 0.16 | 0.0595 | -9.40 | 0.0694 | 5.58 | 0.0715 | 8.76 |
| 1.5 | 0.0455 | 0.0456 | 0.13 | 0.0414 | -9.00 | 0.0479 | 5.16 | 0.0494 | 8.48 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $2 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, \text { simp } 1}$ and $B_{W C}$ simp2, are fairly small (maximum $\sim 6.5 \%$ for the $1^{\text {st }}$ level of approximation and maximum $\sim 9 \%$ for the $2^{\text {nd }}$ level of approximation) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6 ), arrive at $\sim-13 \%$ being always negative.

Fourth Case: same overall arrangement as in the third case, apart that the central cable, with the larger cross-section, is situated 0.5 m higher than the other two ( $H_{2}=1.5 \mathrm{~m}$ ), thus the field points lie on a straight horizontal line whose vertical distance from the central cable is 1.5 m , while their vertical distance from the other two cables is still 1.0 m (see Fig. 4.13). Figure 4.14 and Table 4.7 show the results of the fourth case simulation.


Figure 4.13: Geometrical and electrical data concerning the fourth case-study.


Figure 4.14: $B_{T R I P L E}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simp } 1}, B_{W C, \text { simp } 2}$ vs. distance $x$ for the fourth case-study $\left(I_{1}=I_{3}=280 \mathrm{~A}\right.$ and $\left.I_{2}=360 \mathrm{~A}\right)$.

Table 4.7: Numerical values of $B_{T R I P L E}, B_{W C, \text { exact }}, B_{W C, l i t}, B_{W C, \text { simpl }}, B_{W C, \text { simp } 2}$ vs. distance $x$ and percent errors of $B_{W C, e x a c t}, B_{W C, l i t}, B_{W C, \text { simp1 }}$ and $B_{W C, \text { simp } 2}$ with respect to $B_{\text {TRIPLE }}$ for the fourth case-study $\left(I_{1}=I_{3}=280 \mathrm{~A}\right.$ and $\left.I_{2}=360 \mathrm{~A}\right)$.

| $\begin{gathered} \mathbf{x} \\ {[\mathrm{m}]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {TRIPLE }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | B $_{\text {WC,exact }}$ [ $\mu \mathrm{T}$ ] | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \mathrm{lit}} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\text {WC,simp1 }} \\ {[\mu \mathrm{T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ | $\begin{gathered} \mathbf{B}_{\mathrm{WC}, \text { simp } 2} \\ {[\mu \mathrm{~T}]} \end{gathered}$ | $\begin{gathered} \text { error } \\ {[\%]} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| -1.5 | 0.0338 | 0.0338 | 0.08 | 0.0307 | -9.09 | 0.0355 | 5.13 | 0.0367 | 8.73 |
| -1.4 | 0.0484 | 0.0484 | 0.10 | 0.0438 | -9.51 | 0.0510 | 5.48 | 0.0527 | 8.92 |
| -1.3 | 0.0682 | 0.0683 | 0.12 | 0.0615 | -9.95 | 0.0721 | 5.71 | 0.0744 | 8.99 |
| -1.2 | 0.0946 | 0.0948 | 0.16 | 0.0848 | -10.39 | 0.1001 | 5.83 | 0.1031 | 8.97 |
| -1.1 | 0.1285 | 0.1288 | 0.22 | 0.1146 | -10.82 | 0.1360 | 5.86 | 0.1399 | 8.87 |
| -1.0 | 0.1704 | 0.1709 | 0.30 | 0.1513 | -11.22 | 0.1803 | 5.84 | 0.1852 | 8.72 |
| -0.9 | 0.2197 | 0.2206 | 0.41 | 0.1943 | -11.56 | 0.2325 | 5.81 | 0.2386 | 8.58 |
| -0.8 | 0.2746 | 0.2762 | 0.57 | 0.2422 | -11.81 | 0.2906 | 5.80 | 0.2979 | 8.48 |
| -0.7 | 0.3315 | 0.3342 | 0.80 | 0.2919 | -11.94 | 0.3510 | 5.88 | 0.3597 | 8.49 |
| -0.6 | 0.3856 | 0.3899 | 1.10 | 0.3397 | -11.90 | 0.4091 | 6.09 | 0.4190 | 8.65 |
| -0.5 | 0.4318 | 0.4384 | 1.52 | 0.3815 | -11.66 | 0.4598 | 6.48 | 0.4708 | 9.02 |
| -0.4 | 0.4666 | 0.4760 | 2.03 | 0.4142 | -11.22 | 0.4995 | 7.06 | 0.5114 | 9.60 |
| -0.3 | 0.4887 | 0.5015 | 2.62 | 0.4368 | -10.63 | 0.5267 | 7.77 | 0.5393 | 10.34 |
| -0.2 | 0.5002 | 0.5161 | 3.18 | 0.4501 | -10.01 | 0.5427 | 8.49 | 0.5557 | 11.10 |
| -0.1 | 0.5049 | 0.5229 | 3.56 | 0.4566 | -9.56 | 0.5504 | 9.00 | 0.5636 | 11.63 |
| -0.0 | 0.5063 | 0.5247 | 3.63 | 0.4585 | -9.45 | 0.5525 | 9.12 | 0.5659 | 11.76 |
| 0.1 | 0.5059 | 0.5229 | 3.34 | 0.4566 | -9.75 | 0.5504 | 8.78 | 0.5636 | 11.40 |
| 0.2 <br> 0.3 | 0.5020 | 0.5161 | 2.80 | 0.4501 | -10.34 | 0.5427 | 8.10 | 0.5557 | 10.69 |
| 0.3 | 0.4908 | 0.5015 | 2.17 | 0.4368 | -11.01 | 0.5267 | 7.31 | 0.5393 | 9.87 |
| 0.4 | 0.4686 | 0.4760 | 1.59 | 0.4142 | -11.60 | 0.4995 | 6.60 | 0.5114 | 9.13 |
| 0.5 | 0.4335 | 0.4384 | 1.12 | 0.3815 | -12.01 | 0.4598 | 6.07 | 0.4708 | 8.60 |
| 0.6 | 0.3869 | 0.3899 | 0.77 | 0.3397 | -12.19 | 0.4091 | 5.74 | 0.4190 | 8.30 |
| 0.7 | 0.3324 | 0.3342 | 0.53 | 0.2919 | -12.18 | 0.3510 | 5.60 | 0.3597 | 8.20 |
| 0.8 | 0.2752 | 0.2762 | 0.36 | 0.2422 | -12.00 | 0.2906 | 5.58 | 0.2979 | 8.26 |
| 0.9 | 0.2201 | 0.2206 | 0.25 | 0.1943 | -11.70 | 0.2325 | 5.64 | 0.2386 | 8.41 |
| 1.0 | 0.1706 | 0.1709 | 0.18 | 0.1513 | -11.32 | 0.1803 | 5.72 | 0.1852 | 8.60 |
| 1.1 | 0.1286 | 0.1288 | 0.13 | 0.1146 | -10.89 | 0.1360 | 5.77 | 0.1399 | 8.77 |
| 1.2 | 0.0947 | 0.0948 | 0.10 | 0.0848 | -10.44 | 0.1001 | 5.76 | 0.1031 | 8.90 |
| 1.3 | 0.0683 | 0.0683 | 0.08 | 0.0615 | -9.99 | 0.0721 | 5.66 | 0.0744 | 8.95 |
| 1.4 | 0.0484 | 0.0484 | 0.07 | 0.0438 | -9.54 | 0.0510 | 5.45 | 0.0527 | 8.89 |
| 1.5 | 0.0338 | 0.0338 | 0.06 | 0.0307 | -9.11 | 0.0355 | 5.12 | 0.0367 | 8.71 |

The error computed according to the "worst case" approach, $B_{W C, \text { exact }}$ (eq. 4.5), is below $3.5 \%$ and always positive. The errors computed by the application of the simplified formula in the "worst case" estimate calculation (eqs. 4.7), $B_{W C, s i m p 1}$ and $B_{W C \text { simp2 }}$, are fairly small (maximum $\sim 9 \%$ for the $1^{\text {st }}$ level of approximation and maximum $\sim 12 \%$ for the $2^{\text {nd }}$ level of approximation) and always positive. The errors computed by the approximate formula of the literature in the "worst case" estimate calculation, $B_{W C, l i t}$ (eq. 4.6), arrive at $\sim-12 \%$ being always negative.

### 4.3 CONCLUSIONS

Theexact calculation of the magnetic field generated by multiple-circuit twisted three-phase power cable lines that are the standard lines for transporting renewable energy has been calculated as a generalization of the vector theory brought about in Chapter 3 concerning a double-circuit twisted three-phase power cable lines. In the framework of the'worst-case' approach, the analysis has also extended to multiplecircuit twisted three-phase cables both the exact and approximated formula derived from the literature about single-circuits and the simple innovative formula developed in both levels of approximation.

As shown by the simulations regarding a triple-circuit twisted three-phase cable lines, underground and overhead, the error computed by the 'worst case' estimate is fairly small making the relevant expression (eq. 4.5) an efficient alternative way of calculating the magnetic field. The use of the innovative simplified formula (eqs. 4.7) in the 'worst case' approach further simplifies the field calculation with an inevitable positive error of maximum $\sim 9 \%$ for the first and $\sim 12 \%$ for the second level of approximation. On the other hand, the error computed by the use of the approximated formula of the literature in the 'worst case' approach, (eq. 4.6), is always higher in absolute terms and negative, underestimating the field. It is thus confirmed once more that the innovative formula results definitely much simpler than the exact one and provides a much smaller relative error compared to the approximated one from the literature and it can be used for the calculus of the magnetic field from multiple twisted field sources without problems of accuracy.

Concluding, expressions (4.7) result an efficient alternative way of calculating the magnetic field of a multiple-circuit twisted three-phase cable lines. This is due to its advantages which lie on its simplicity - an algebraic sum of the rms values of the total field $B$ generated by each twisted cable, calculated without using sophisticated mathematics - and on its reliability - the maximum error presented is below $12 \%$ for both levels of approximation and positive providing a conservative estimation of the magnetic field. Unfortunately, experimental validation by measurements in situ of the above used formulas for multiple-circuit twisted three-phase cable lines was not effectuated due to the extreme difficulty of finding easily accessible multiple-circuit twisted cable lines with known instant values of rms line currents.

## CHAPTER 5

## Development of Numerical Codes for the Calculation of the magnetic Field Generated by Overhead Lines

### 5.1 MAGNETIC FIELD CALCULUS

The magnetic field generated by an electrical current is given by the Biot-Savart law that relates the magnetic field to the magnitude, direction, length, and proximity of the electric current. The Biot-Savart law is used to compute the resultant magnetic field $B$ at position $r$ generated by a current $i(t)$ (see Fig. 5.1).


Figure 5.1: Representation of a filamentary conductor in close path C carrying current $i(t)$.
The law is a physical example of a line integrale valuated over the path $C$ where the electric currents flow:

$$
\begin{equation*}
\vec{B}(t)=\frac{\mu_{0}}{4 \pi} i(t) \int_{c} \frac{d \vec{l} \times \vec{r}}{|\vec{r}|^{2}} \tag{5.1}
\end{equation*}
$$

In case of a straight current-carrying wire of finite length AB (see Fig. 5.2) the BiotSavart law (eq. 5.1) becomes:

$$
\begin{equation*}
\Delta \vec{B}(t)=\frac{\mu_{0}}{2 \pi R} i(t) \frac{\cos \alpha+\cos \beta}{2} \vec{u}_{n} \tag{5.2}
\end{equation*}
$$

where $R$ is the distance between the field-point $P$ and the finite length segment $A B, \alpha$ and $\beta$ are the angles indicated in Fig. 5.2, and $\vec{u}_{n}$ is a vector orthogonal to the plane determined by points $\mathrm{A}, \mathrm{B}$ and P with the direction of the vector product $d \vec{l} \times \vec{r}$.


Figure 5.2: Representation of a filamentary conductor segment AB carrying current i(t).
One can notice that by increasing the length of the segment AB the term $\frac{\cos \alpha+\cos \beta}{2}$ tends to 1 , and so the classic formula of Biot-Savart for filamentary, straight, infinitely long wires derives:

$$
\begin{equation*}
\overrightarrow{\mathrm{B}}(t)_{\infty}=\frac{\mu_{0}}{2 \pi r} \cdot i(t) \cdot \vec{u}_{n} \tag{5.3}
\end{equation*}
$$

### 5.2 2D MAGNETIC FIELD CALCULATION CODE

### 5.2.1 Theory

As far as an overhead or underground power line with $n$ conductors is concerned, making the following assumptions and simplifications:

- the algorithm considers all conductors constituting the line to be straight, horizontal, of infinite length and parallel to each other;
- the currents on each conductor are considered to be in-phase;
- the soil is considered flat and free of irregularities -perfectly transparent from the magnetic point of view;
- the presence of towers, utility poles, buildings, vegetation and any other object is in the area concerned is neglected.
the calculation of the magnetic field is reduced into a simple plane problem, easily implemented into a software tool, since the situation is exactly the same on every section of the line, meaning the vertical plane orthogonal to the longitudinal axis of the line (i.e. the direction of the conductors that constitute it) that passes through the field-point of interest. In this case, the magnetic field components are given by the following formulas, as also defined by the Italian law [48]:

$$
\begin{gather*}
B_{x}(t)=\frac{\mu_{0}}{2 \pi} \sum_{n} i_{n}(t)\left[\frac{\left(y_{n}-y_{p}\right)}{\left(x_{p}-x_{n}\right)^{2}+\left(y_{p}-y_{n}\right)^{2}}\right]  \tag{5.4.a}\\
B_{y}(t)=\frac{\mu_{0}}{2 \pi} \sum_{n} i_{n}(t)\left[\frac{\left(x_{p}-x_{n}\right)}{\left(x_{p}-x_{n}\right)^{2}+\left(y_{p}-y_{n}\right)^{2}}\right] \tag{5.4.b}
\end{gather*}
$$

In order to fully characterize the given power line from both a geometric and an electric point of view, a set of input data is required that includes the number of conductors that constitute the power line, $n$, the position of the conductors given in Cartesian coordinates $\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right),\left(x_{3}, y_{3}\right), \ldots,\left(x_{n}, y_{n}\right)$, the coordinates of the fieldpoints of interest, $\left(x_{p}, y_{p}\right)$, the rms currents that flow in these conductors, $I_{1}, I_{2}, I_{3}, \ldots, I_{n}$, and the electrical phase of the currents, $\varphi_{1}, \varphi_{2}, \varphi_{3}, \ldots, \varphi_{n}$. The field-points lay on the $x$ axis that is situated at a distance $H$ from the ground (see Fig. 5.3).


Figure 5.3:Vertical section (x, y plane) of a (a) single- and a(b) double-circuit three-phase overhead line.

One can notice that expressions (5.4) use the instantaneous values of the currents $i(t)$. In the case of currents that vary sinusoidally with time it is possible, and preferable, to use the phasors (complex numbers) of the currents, providing thus the phasors of the magnetic induction components. Programming in MatLab ${ }^{\text {TM }}$ environment consents the direct treatment of complex numbers and eqs. (5.4) were implemented using the currents phasors. So, after having calculated the magnetic field components phasors, the program provides their rms values to be subsequently used for the composistion of the total bidimensional magnetic field from the equation:

$$
\begin{equation*}
B=\sqrt{B_{x}^{2}+B_{y}^{2}} \tag{5.5}
\end{equation*}
$$

### 5.2.2 Applications - Simulation Results

### 5.2.2.1 132 kV Single-Circuit Three-Phase Overhead Line

Figures 5.4 show the geometrical characteristics of the 132 kV single-circuit threephase overhead line used for the simulations, along with the vertical section under exam, corresponding at the mid-span between two adjacent poles.


Figure 5.4: 132 kV single-circuit three-phase overhead line:
(a) geometrical characteristics and (b) vertical mid-span section (z-y plane).

The profiles of the field components $B_{z}$ and $B_{y}$ along with the total magnetic field $B$ were calculated for $I=385 \mathrm{~A}$ and $z \in[-50 \mathrm{~m} ; 50 \mathrm{~m}]$ at different heights $H$ from the ground and the results are depicted in Figs. 5.5-5.8.


Figure 5.5: $B_{z}, B_{y}, B$ vs. $z$ at $2 m$ from the ground calculated at the mid-span section of the 132 $k V$ single-circuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.6: $B_{z}, B_{y}, B$ vs. $z$ at 6 m from the ground calculated at the mid-span section of the 132 $k V$ single-circuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.7: $B_{z}, B_{y}, B$ vs. $z$ at 10 m from the ground calculated at the mid-span section of the 132 $k V$ single-circuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.8: $B_{z}, B_{y}, B$ vs. $z$ at 12.5 m from the ground calculated at the mid-span section of the 132 kV single-circuit overhead line with $I=385 \mathrm{~A}$.

### 5.2.2.2 380 kV Double-Circuit Three-Phase Overhead Line

Figures 5.9 show the geometrical characteristics of a typical Italian 380 kV doublecircuit three-phase overhead line used for the simulations, along with the vertical section under exam, corresponding at the mid-span between two adjacent poles. The simulation has been made considering the reverse arrangement of the conductors phases with $I_{1}=I_{2}$ $=2040$ A, i.e. the winter season ampacity for Italy, zone B, according to [49].


Figure 5.9: 380 kV double-circuit three-phase overhead line:
(a) geometrical characteristics and (b) vertical mid-span section (not in scale).

The profiles of the field components $B_{z}$ and $B_{y}$ along with the total magnetic field $B$ were calculated for $I_{l}=I_{2}=2040 \mathrm{~A}$ and $z \in[-100 \mathrm{~m} ; 100 \mathrm{~m}]$ at different heights $H$ from the ground and the results are depicted in Figs. 5.10-5.13.


Figure 5.10: $B_{z}, B_{y}, B$ vs. $z$ at $5 m$ from the ground calculated at the mid-span section of the 380 $k V$ double-circuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.11: $B_{z}, B_{y}, B v s . z$ at 10 m from the ground calculated at the mid-span section of the 380 kV double-circuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.12: $B_{z}, B_{y}, B v s . z$ at 20 m from the ground calculated at the mid-span section of the 380 kV double-circuit overhead power with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.13: $B_{z}, B_{y}, B v s . z$ at $28 m$ from the ground calculated at the mid-span section of the 380 kV double-circuit overhead line with $I_{I}=I_{2}=2040 \mathrm{~A}$.

### 5.3 3D MAGNETIC FIELD CALCULATION CODE

### 5.3.1 Theory

The exact calculation of the magnetic field generated by real catenary-form overhead lines has been treated in [34] - [50] and, in particular, the Italian law [48] defines two methods of calculating the total tridimensional magnetic field generated by $n$ conductors, considering that the each conductor is divided in infinite straight segments carrying current $i(t)$. Doing so, it is possible to reconstruct the catenary form of the conductor line by simulating it with a sequence of a large number of straight segments. Calculating the contribution of each segment and performing the summation over the entire conductor line and repeating this procedure for every conductor present in the power line, the three magnetic field components, $B_{x}, B_{y}, B_{z}$, at the given field point will be also calculated.

## Method 1

- Determination of the direction parameters $(a, b, c)$ of the line that passes from points A $\left(x_{A}, y_{A}, z_{A}\right)$ and $\mathrm{B}\left(x_{B}, y_{B}, z_{B}\right)$ (see Fig. 5.2):

$$
\begin{gather*}
a=x_{B}-x_{A}  \tag{5.6.a}\\
b=y_{B}-y_{A}  \tag{5.6.b}\\
c=z_{B}-z_{A} \tag{5.6.c}
\end{gather*}
$$

- Determination of the coefficients $\left(a_{1}, b_{1}, c_{l}, d_{1}\right)$ of the plane equation that passes from point $\mathrm{P}\left(x_{P}, y_{P}, z_{P}\right)$ and is orthogonal to the line (see Fig. 5.2):

$$
\begin{gather*}
a_{1}=a  \tag{5.7.a}\\
b_{1}=b  \tag{5.7.b}\\
c_{1}=c  \tag{5.7.c}\\
d_{1}=-\left(a \cdot x_{P}+b \cdot y_{P}+c \cdot z_{P}\right) \tag{5.8}
\end{gather*}
$$

- Determination of the intersection point $\mathrm{H}\left(x_{H}, y_{H}, z_{H}\right)$ of the orthogonal plane with the line (see Fig. 5.2):

$$
\begin{align*}
& x_{H}=x_{A}+a \cdot t  \tag{5.9.a}\\
& y_{H}=y_{A}+b \cdot t  \tag{5.9.b}\\
& z_{H}=z_{A}+c \cdot t \tag{5.9.c}
\end{align*}
$$

with

$$
\begin{equation*}
t=-\frac{a_{1} \cdot x_{A}+b_{1} \cdot y_{A}+c_{1} \cdot z_{A}+d_{1}}{a \cdot a_{1}+b \cdot b_{1}+c \cdot c_{1}} \tag{5.10}
\end{equation*}
$$

- Determination of the segment PH length $L_{P H}$ (distance $R$ between the line and point P ) (see Fig. 5.2):

$$
\begin{equation*}
R=L_{P H}=\sqrt{\left(x_{P}-x_{H}\right)^{2}+\left(y_{P}-y_{H}\right)^{2}+\left(y_{P}-y_{H}\right)^{2}} \tag{5.11}
\end{equation*}
$$

## Method 2

- Determination of the segment PH length $L_{P H}$ (distance $R$ between the line and point P ) from the rectangle PHA (see Fig. 5.2):
with

$$
\begin{equation*}
R=L_{P H}=\sqrt{L_{P A}^{2}-L_{A H}^{2}} \tag{5.12}
\end{equation*}
$$

Once calculated the distance $R$, the magnetic field magnitude can be then valuated from the following formula:

$$
\begin{equation*}
\Delta B(t)=\frac{i(t)}{4 \cdot \pi \cdot R}\left(\frac{\left(L_{A B}-L_{A H}\right)}{L_{P B}}+\frac{L_{A H}}{L_{P A}}\right) \tag{5.14}
\end{equation*}
$$

with

$$
\begin{equation*}
L_{P B}=\sqrt{\left(x_{P}-x_{B}\right)^{2}+\left(y_{P}-y_{B}\right)^{2}+\left(y_{P}-y_{B}\right)^{2}} \tag{5.15}
\end{equation*}
$$

In order to calculate the total magnetic field generated by all the segments of all conductors constituting the line, it is necessary determine the three spatial components of the magnetic field by means of the following expressions:

$$
\begin{align*}
& \Delta B_{x}(t)=\Delta B(t) \cdot \frac{\cos \alpha}{\sqrt{\cos \alpha^{2}+\cos \beta^{2}+\cos \gamma^{2}}}  \tag{5.16.a}\\
& \Delta B_{y}(t)=\Delta B(t) \cdot \frac{\cos \beta}{\sqrt{\cos \alpha^{2}+\cos \beta^{2}+\cos \gamma^{2}}}  \tag{5.16.b}\\
& \Delta B_{z}(t)=\Delta B(t) \cdot \frac{\cos \gamma}{\sqrt{\cos \alpha^{2}+\cos \beta^{2}+\cos \gamma^{2}}} \tag{5.16.c}
\end{align*}
$$

where the direction cosines are functions of the spatial coordinates of the points $\mathrm{A}, \mathrm{B}$ and P , as follows:

$$
\begin{align*}
& \cos \alpha=y_{A} \cdot z_{B}+z_{A} \cdot y_{P}+y_{B} \cdot z_{P}-z_{B} \cdot y_{P}-z_{A} \cdot y_{B}-y_{A} \cdot z_{P}  \tag{5.17.a}\\
& \cos \beta=z_{B} \cdot x_{P}+z_{P} \cdot x_{A}+z_{A} \cdot x_{B}-x_{A} \cdot z_{B}-z_{A} \cdot x_{P}-x_{B} \cdot z_{P}  \tag{5.17.b}\\
& \cos \gamma=x_{A} \cdot y_{B}+x_{P} \cdot y_{A}+x_{B} \cdot y_{P}-x_{P} \cdot y_{B}-x_{B} \cdot y_{A}-x_{A} \cdot y_{P} \tag{5.17.c}
\end{align*}
$$

Finally, the data needed for the definition of power lines also include the parameters of the catenary formula in Cartesian coordinates:

$$
\begin{equation*}
y=\alpha \cdot \cosh \left(\frac{x-\left(L / 2-X_{1}\right)}{\alpha}\right)+(F-a) \tag{5.18}
\end{equation*}
$$

where X 1 and X 2 are the catenary suspension points positioned at height $F$ and having length $L$ (see Fig. 5.14 ) and $a$ is the catenary parameter given by the following expression as a function of the $\operatorname{sag} f$ :

$$
\begin{equation*}
a=\frac{\left(L^{2}\right)}{8 f} \tag{5.19}
\end{equation*}
$$



Figure 5.14: Schematization of the catenary parameters.

### 5.3.2 Applications - Simulation Results

### 5.3.2.1 132 kV Single-Circuit Three-Phase Overhead Line

Considering the same 132 kV single-circuit three-phase overhead line of Fig. 5.4 suspended at six poles, thus, forming five catenaries, each with length of 200 m for a total of 1000 m , the following Figs. 5.15-5.26 were plotted as the results of the 3D simulations of the magnetic field generated by such a line carrying balanced three-phase current $I=385 \mathrm{~A}$.


Figure 5.15: $3 D$ plot of $B$ vs. $z$ and $x$ at $2 m$ from the ground for the 132 kV single-circuit overhead line with $I=385$ A.


Figure 5.16: B isolines ( $1 \mu T$ and $3 \mu T$ ) at $2 m$ from the ground for the 132 kV single-circuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.17: Details of Fig. 5.16 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 132 kV single-circuit overhead lines $(I=385 \mathrm{~A})$.


Figure 5.18: 3D plot of B vs. $z$ and $x$ at 6 m from the ground for the 132 kV single-circuit overhead line with $I=385$ A..


Figure 5.19: $B$ isolines ( $1 \mu \mathrm{~T}$ and $3 \mu \mathrm{~T}$ ) at 6 m from the ground for the 132 kV single-circuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.20: Details of Fig. 5.19 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 132 kV single-circuit overhead line ( $I=385 \mathrm{~A}$ ).


Figure 5.21: $3 D$ plot of $B$ vs. $z$ and $x$ at 10 m from the ground for the 132 kV single-circuit overhead line with $I=385$ A.


Figure 5.22: B isolines ( $3 \mu T$ and $10 \mu T)$ at 10 m from the ground for the 132 kV single-circuit overhead line with $I=385$ A.


(c)

Figure 5.23: Details of Fig. 5.22 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and $(\boldsymbol{b})$ the middle of the 132 kV single-circuit overhead line ( $I=385 \mathrm{~A}$ ).


Figure 5.24: $3 D$ plot of $B$ vs. $z$ and $x$ at 12.5 m from the ground for the 132 kV single-circuit overhead line with $I=385$ A.


Figure 5.25: B isolines ( $3 \mu T, 10 \mu T$ and $100 \mu T$ ) at 12.5 m from the ground for the 132 kV singlecircuit overhead line with $I=385 \mathrm{~A}$.


Figure 5.26: Details of Fig. 5.25 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 132 kV single-circuit overhead line ( $I=385 \mathrm{~A}$ ).

### 5.3.2.2 380 kV Double-Circuit Three-Phase Overhead Line

Considering the 380 kV single-circuit three-phase overhead line of Fig. 5.9 carrying balanced three-phase currents $I_{1}=I_{2}=2040 \mathrm{~A}$, suspended at six poles, thus, forming five catenaries, each with length of 400 m for a total of 2000 m with $\operatorname{sag} f=7$ m, Figs. 5.27-5.38 were plotted as the results of the 3 D simulations of the magnetic field generated by such a line.


Figure 5.27: $3 D$ plot of $B$ vs. $z$ and $x$ at 5 m from the ground for the 380 kV double-circuit overhead line with $I_{I}=I_{2}=2040 \mathrm{~A}$.


Figure 5.28: B isolines ( $1 \mu T$ and $3 \mu T$ ) at $5 m$ from the ground for the $380 k V$ double-circuit overhead line with $I_{I}=I_{2}=2040 \mathrm{~A}$.


Figure 5.29: Details of Fig. 5.28 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 380 kV double-circuit overhead line ( $I_{l}=$ $\left.I_{2}=2040 \mathrm{~A}\right)$.


Figure 5.30: $3 D$ plot of $B v s . z$ and $x$ at 10 m from the ground for the 380 kV double-circuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.31: B isolines ( $1 \mu T, 3 \mu T$ and $10 \mu T$ ) at 10 m from the ground for the 380 kV doublecircuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.32: Details of Fig. 5.31 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 380 kV double-circuit overhead line ( $I_{l}=$ $\left.I_{2}=2040 \mathrm{~A}\right)$.


Figure 5.33: 3 D plot of $B \mathrm{vs} . z$ and $x$ at 20 m from the ground for the 380 kV double-circuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.34: $B$ isolines ( $1 \mu T, 3 \mu T$ and $10 \mu T$ ) at 20 m from the ground for the 380 kV doublecircuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.35: Details of Fig. 5.34 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 380 kV double-circuit overhead line ( $I_{l}=$ $\left.I_{2}=2040 \mathrm{~A}\right)$.


Figure 5.36: 3 D plot of $B \mathrm{vs}$.z and $x$ at 28 m from the ground for the 380 kV double-circuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.37: B isolines ( $3 \mu T, 10 \mu T$ and $100 \mu T)$ at $28 m$ from the ground for the 380 kV doublecircuit overhead line with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.38: Details of Fig. 5.37 concerning (a) the beginning (in equal-scaled axes), (b) the end (in equal-scaled axes) and (c) the middle of the 380 kV double-circuit overhead line ( $I_{l}=$ $\left.I_{2}=2040 \mathrm{~A}\right)$.

### 5.4 CONCLUSIONS

The results of all 2D simulations made were compared with magnetic field calculation programs commercially available. In particular, the software used, that calculates the electric and magnetic field generated by overhead and underground power lines, is "CAMPI" [51], in C++ language, developed and distributed by the "NelloCarrara" Institute of Applied Physics (IFAC), part of the National Research Council (CNR), which is the main public organization pursuing research and innovation in Italy.

The results of all 3D simulations made were compared with magnetic field calculation programs not commercially available and classified. Nevertheless, some
critical observations concerning the correctness of the script were made and are depicted in Figs. 5.39-5.70.

Three sections of the 132 kV single-circuit overhead line relative to the middle catenary (for $x$ values from 400 m to 600 m ) were further examined, for the purpose of comparing the values of the magnetic field calculated according to the bi-dimensional approach, $B_{2 D}$, with that calculated according to the tri-dimensional approach, $B_{3 D}$, and the relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ has been estimated, at different heights from the ground. More precisely, the first section considered is at $x=400 \mathrm{~m}$ (under the pole), the second is at $x=450 \mathrm{~m}$ ( $1 / 4$ of the catenary) and the third section is at $x=500 \mathrm{~m}$ (at the mid-span) and the simulation results are depicted in Figs. 5.39-5.46.


Figure 5.39: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=400 m, x=450$ $m$ and $x=500 m$ for the 132 kV single-circuit overhead line at $2 m$ from the ground with $I=385 \mathrm{~A}$.



Figure 5.40: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.39.


Figure 5.41: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=400 m, x=450$ $m$ and $x=500 m$ for the 132 kV single-circuit overhead line at 6 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.42: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.41.


Figure 5.43: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=400 m, x=450 m$ and $x=500 \mathrm{~m}$ for the 132 kV single-circuit overhead line at 10 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.44: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.43.


Figure 5.39: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=400 m, x=450 \mathrm{~m}$ and $x=500 \mathrm{~m}$ for the 132 kV single-circuit overhead line at 12.5 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.46: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.45.

The same considerations were made also for the 380 kV power line case where also three vertical sections relative to the middle catenary (for $x$ values from 800 m to 1200 m ) were further examined, at different heights from the ground, for the purpose of comparing the values of $B_{2 D}$ with that of $B_{3 D}$ and estimating the relative percent error of $B_{2 D}$ with respect to $B_{3 D}$. The first section considered is at $x=800 \mathrm{~m}$ (under the pole), the second is at $x=900 \mathrm{~m}(1 / 4$ of the catenary) and the third section is at $x=1000 \mathrm{~m}$ (at the mid-span) and the results are depicted in Figs. 5.47-5.54.


Figure 5.47: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=800 m, x=900 \mathrm{~m}$ and $x=1000 \mathrm{~m}$ for the 380 kV double-circuit overhead line at 5 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.48: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.47.


Figure 5.49: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=800 m, x=900 \mathrm{~m}$ and $x=1000 \mathrm{~m}$ for the 380 kV double-circuit overhead line at 10 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.50: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.49.


Figure 5.51: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=800 m, x=900 m$ and $x=1000 \mathrm{~m}$ for the 380 kV double-circuit overhead line at 20 m from the ground with $I_{l}=I_{2}=2040 \mathrm{~A}$.


Figure 5.52: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.51.


Figure 5.53: Comparison between $B_{2 D}$ and $B_{3 D}$ at the three vertical sections of $x=800 m, x=900 m$ and $x=1000 \mathrm{~m}$ for the 380 kV double-circuit overhead line at 28 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.54: Relative percent error of $B_{2 D}$ with respect to $B_{3 D}$ for the sections considered at Fig. 5.53.

One can notice the growing discrepancy between the results obtained at the mid-span between the poles (third section) and those under the pole (first section). For reasons of providing conservative results, the bi-dimensional approach calculates the field in correspondence to the mid-span between two poles where the conductors are closer to the ground and the field results higher. That explains why the field curves of obtained at the third section are almost coincident unlike the field curves obtained at the second and, particularly, the first section, where the maximum divergence between $B_{2 D}$ and $B_{3 D}$ is presented. Another reason for this deviation of the results, is the fact that in the area under the pole, the line segments are not horizontal and parallel to each other as hypothesized by the bi-dimensional calculus.

Additionally, it must be pointed out, that for the 132 kV single-circuit overhead line, the curves of $B_{3 D}$ become closer to those of $B_{2 D}$ as the distance $H$ from the ground rises. This is logical and expected since the effect of the catenary-form of the line is compensated by the small distance from the field sources. This is not true, though, for the 380 kV double-circuit overhead line case due, firstly, to the greater value of the currents and, secondly, to the geometry of the line that is more similar to the considered geometry of the line in the bi-dimensional case (i.e. a straight horizontal line parallel to the ground), due to the greater length $L$ of every catenary and the greater value of $a$ in eq. (5.18).

Another interesting point to be examined is the overall behavior of the $B_{3 D}$ field along the whole overhead line (from $x=0 \mathrm{~m}$ up to $x=1000 \mathrm{~m}$ for the 132 kV single-circuit overhead line and from $x=0 \mathrm{~m}$ up to $x=2000 \mathrm{~m}$ for the 380 kV double-circuit overhead line) under every pole and under every mid-span area. For this reason, Figs. 5.55-5.62 are plotted that depict the magnetic field $B_{3 D}$ as a function of the coordinate $z$ calculated at the vertical sections of $x=0 \mathrm{~m}$ (first pole), $x=100 \mathrm{~m}$ (first mid-span), $x=$ 200 m (second pole), $x=300 \mathrm{~m}$ (second mid-span), $x=400 \mathrm{~m}$ (third pole), $x=500 \mathrm{~m}$ (third mid-span), $x=600 \mathrm{~m}$ (fourth pole), $x=700 \mathrm{~m}$ (fourth mid-span), $x=800 \mathrm{~m}$ (fifth pole), $x=900 \mathrm{~m}$ (fifth mid-span) and $x=1000 \mathrm{~m}$ (sixth and last pole) for the 132 kV single-circuit overhead line and at the vertical sections of $x=0 \mathrm{~m}$ (first pole), $x=200 \mathrm{~m}$ (first mid-span), $x=400 \mathrm{~m}$ (second pole), $x=600 \mathrm{~m}$ (second mid-span), $x=800 \mathrm{~m}$ (third pole), $x=1000 \mathrm{~m}$ (third mid-span), $x=1200 \mathrm{~m}$ (fourth pole), $x=1400 \mathrm{~m}$ (fourth mid-span), $x=1600 \mathrm{~m}$ (fifth pole), $x=1800 \mathrm{~m}$ (fifth mid-span) and $x=2000 \mathrm{~m}$ (sixth and last pole) for the 380 kV double-circuit overhead line.


Figure 5.55: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 132 kV single-circuit overhead line, at 2 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.56: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 132 kV single-circuit overhead line, at 6 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.57: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 132 kV single-circuit overhead line, at 10 m from the ground with $I=385$ A.


Figure 5.58: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 132 kV single-circuit overhead line, at 12.5 m from the ground with $I=385 \mathrm{~A}$.

In Figs. 5.55-5.58 the magnetic field profiles are categorized in three groups since their curves are practically coincident. There is the first group of $B_{x=100 \mathrm{~m}}, B_{x=300 \mathrm{~m}}$, $B_{x=500 m}, B_{x=700 \mathrm{~m}}$ and $B_{x=900 m}$ (under the mid-span areas), the second group of the $B_{x=400 \mathrm{~m}}$, $B_{x=600 m}$ and $B_{x=800 m}$ (under the second, third and fourth pole) and the third group of $B_{x=100 m}$ and $B_{x=1000 m}$ (under the first and fifth pole). The results, at all heights from the ground considered, have had a positive outcome since the first group of the $B_{3 D}$ profiles provide the highest values of the field and the second group provides two times the values of the third group. These considerations were expected: the first group provides the magnetic field values under the mid-span areas where the field is maximum; the third group calculates the field under the poles situated at the two ends of the line, meaning that the magnetic field contribute of half the line is missing, resulting unavoidably at being $\sim 50 \%$ of the field calculated under the intermediate poles of the line, i.e. second, third and fourth pole (second group).


Figure 5.59: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 380 kV double-circuit overhead line, at 5 m from the ground with $I_{l}=I_{2}=2040 \mathrm{~A}$.


Figure 5.60: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 380 kV double-circuit overhead line, at 10 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.61: Comparison between the magnetic field profiles at the vertical sections of interest. i.e. under the poles and the mid-spans, for the 380 kV double-circuit overhead line, at 20 m from the ground with $I_{I}=I_{2}=2040 \mathrm{~A}$.


Figure 5.62: Comparison between the magnetic field profiles at the vertical sections of interest, i.e. under the poles and the mid-spans, for the 380 kV double-circuit overhead line, at 28 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.

Same considerations hold for the results depicted in Figs. 5.59-5.62 about the 380 kV double-circuit overhead line. There are three groups of coincident magnetic field profiles: the first group of $B_{x=200 m}, B_{x=600 m}, B_{x=1000 m}, B_{x=1400 m}$ and $B_{x=1800 m}$ (under the mid-span areas), the second group of $B_{x=800 \mathrm{~m}}, B_{x=1200 \mathrm{~m}}$ and $B_{x=1600 \mathrm{~m}}$ (under the second, third and fourth pole) and the third group of $B_{x=200 m}$ and $B_{x=2000 m}$ (under the first and fifth pole). The results, at all heights from the ground considered, have also proven that the first group of the $B_{3 D}$ profiles provide the highest values of the field and the second group provides two times the values of the third group. The same considerations as in the single-circuit case hold: the first group of field profiles is the highest since it effectuates the calculation under the mid-span areas and the second group of field profiles results two times the third group since it calculates the field under the intermediate poles of the line.

Finally, the case where $B_{3 D}$ tends at $B_{2 D}$ is examined, in the sense that the tridimensional calculation is made without considering the catenary effect of the line. This is achieved by increasing the parameter $a$ of the catenary in eq. 5.18 and Figs. 5.635.70 are the results of the simulations made for the sake of comparison between $B_{2 D}$ and $B_{3 D}$ at the same vertical sections considered previously, i.e. at $x=0 \mathrm{~m}, x=100 \mathrm{~m}, x=$ $200 \mathrm{~m}, x=300 \mathrm{~m}, x=400 \mathrm{~m}, x=500 \mathrm{~m}, x=600 \mathrm{~m}, x=700 \mathrm{~m}, x=800 \mathrm{~m}, x=900 \mathrm{~m}$ and $x=1000 \mathrm{~m}$ for the 132 kV single-circuit overhead line and at $x=0 \mathrm{~m}, x=200 \mathrm{~m}, x$ $=400 \mathrm{~m}, x=600 \mathrm{~m}, x=800 \mathrm{~m}, x=1000 \mathrm{~m}, x=1200 \mathrm{~m}, x=1400 \mathrm{~m}, x=1600 \mathrm{~m}, x=$

1800 m and $x=2000 \mathrm{~m}$ for the 380 kV double-circuit overhead line. The simulations according to the tri-dimensional approach have been also made in correspondence to the mid-span area, similar to the bi-dimensional approach, in order to have comparable results.


Figure 5.63: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 132 kV single-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 2 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.64: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 132 kV single-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 6 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.65: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 132 kV single-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 10 m from the ground with $I=385 \mathrm{~A}$.


Figure 5.66: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 132 kV single-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 12.5 m from the ground with $I=385 \mathrm{~A}$.

From Figs. 5.63-5.66, dealing the 132 kV single-circuit overhead line, is verified, as expected, that the magnetic field profiles $\left(B_{x=100 \mathrm{~m}}, B_{x=200 \mathrm{~m}}, B_{x=300 \mathrm{~m}}, B_{x=400 \mathrm{~m}}, B_{x=500 \mathrm{~m}}\right.$, $B_{x=600 m}, B_{x=700 m}, B_{x=800 m}$ and $B_{x=900 \mathrm{~m}}$ ) along the straight, finite, parallel to the ground line, calculated at different heights from the ground are the same, except for the two extreme cases of $B_{x=0 m}$ and $B_{x=1000 m}$ that result half times smaller, as explained previously. Another consequence of the fact that the line conductors have not infinite length is that in the cases of the field calculation at 2 m and 6 m from the ground, the field $B_{2 D}$, derived by the bi-dimensional calculation code, still results higher than the field $B_{3 D}$ derived by the tri-dimensional calculation code. Nevertheless, the influence of the finite length of the line is compensated by the smaller distance of the conductors from the field-points in the cases where the conductors distant 10 m and 12.5 m from the ground resulting at coincident curves of the $B_{2 D}$ and $B_{3 D}$ profiles.


Figure 5.67: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 380 kV double-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 5 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.68: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 380 kV double-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 10 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.69: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 380 kV double-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 20 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.


Figure 5.70: Comparison between the magnetic field profiles of $B_{2 D}$ and $B_{3 D}$ calculated without considering the catenary form of the 380 kV double-circuit overhead line, at the vertical sections of interest, i.e. under the poles and the mid-spans, at 28 m from the ground with $I_{1}=I_{2}=2040 \mathrm{~A}$.

Same considerations hold for the 380 kV double-circuit overhead line as proven by Figs. 5.67-5.70. The magnetic field profiles ( $B_{x=200 \mathrm{~m}}, B_{x=400 \mathrm{~m}}, B_{x=600 \mathrm{~m}}, B_{x=800 \mathrm{~m}}, B_{x=1000 \mathrm{~m}}$, $B_{x=1200 m}, B_{x=1400 m}, B_{x=1600 m}$ and $B_{x=1800 m}$ ) along the straight, finite, parallel to the ground line, calculated at different heights from the ground are the same, except for the two extreme cases of $B_{x=0 m}$ and $B_{x=2000 m}$ that result half times smaller, because they are calculated at the two extreme points of the line. In these cases, the influence of the finite length of the line is drastically compensated by the great value of the currents, resulting thus, at coincident curves of the $B_{2 D}$ and $B_{3 D}$ profiles for all calculated heights of the ground.

Concluding, all the above mentioned observations regarding the tri-dimensional numerical code for the calculation of the magnetic field generated by overhead power lines, can serve as proves for the proper implementation of the magnetic field theory and for the correct functioning of the software created.

## CHAPTER 6

## "Smart" Measurement and Evaluation System of the Magnetic Field

### 6.1 GENERAL IDEA

The growing concern among the population for the impact of electromagnetic fields from power systems, unavoidably leads to the construction of sophisticated measurement and evaluation systems of the magnetic field that can provide reliable and accurate indications even in complex 3D arrangements of the field sources. Such arrangements occur more and more frequently in the vicinity of residential and industrial areas, where overhead and underground lines of different voltage rating and geometry, as well as more or less wide substations, often lie close to each other. In this framework, an innovative measurement and evaluation system capable of matching this need has been constructed, calibrated and tested by actual measurements in situ.

The innovation of this measurement and evaluation system is twofold:

1) the "smart" measurement device - that consists essentially of a 3 D field probe plus proper signal conditioning circuits plus a data-logger system - not only measures the components of the magnetic induction field in a Cartesian coordinate system, $B_{x}, B_{y}$ and $B_{z}$ [53] - [56], but also records and stores the rms value of the three field components as a function of time, for subsequent analytical processing. Usually magnetic field measuring devices record and store only the rms value of total magnetic induction field, $B$;
2) a 3D post-processing of the three field components stored over the monitoring period is performed, where the values of the field components are correlated to the corresponding rms values of source currents - that vary randomly with time as the load changes - via a multi-linear regression technique. This enables the extrapolation of the rms value of the magnetic field to any combination of source currents of interest. Indeed, when the number and the geometrical complexity of the sources is increased, a linear relationship between the rms value of total magnetic induction and source currents can be hardly found, whereas one can be pretty confident that such a linear relationship exists between the rms values of magnetic induction components and source currents. Moreover, it must be pointed out that in cases of multiple and complex 3D field sources, in practice the exact geometry of the sources is often unknown or not accessible: in such cases, even powerful 3D codes might fail in forecasting accurately the field value at a given point of interest. For this reason, where a 3D code fails an empirically-derived multi-linear relationship can succeed.

The investigation is divided in three stages:

1) the first stage is carried out in the laboratory and has the purpose of creating an appropriate three-dimensional measurement device of the magnetic field, that is based on an isotropic 3D coil employed as a magnetic field probe (see Fig. 6.1);


Figure 6.1: An isotropic 3D coil employed as a magnetic field probe.
2) the second stage is a theoretical study, where analytical multilinear regression algorithms are developed in Matlab ${ }^{\circledR}$ environment for correlating measurement results of $B_{x}, B_{y}$ and $B_{z}$ at a given field point to the relevant values of time-varying currents of a certain batch of sources. An increasing number of sources can be considered, thereby increasing the complexity and the effectiveness of the algorithm at the same time. The theoretical basis for this development is contained in [57] - [59]. Of course, this stage can be performed at the same time or even before the first stage;
3) the third part consists in an experimental campaign in the field in the presence of multiple current sources. The field components monitored and recorded are processed later on via the algorithm developed at stage 2), in order to derive the numerical values of the coefficients of the multilinear relationship between field components and source currents, thereby validating the evaluation algorithm.

### 6.2 CONSTRUCTION, CALIBRATION AND CHARACTERIZATION OF THE MEASUREMENT DEVICE

The block diagram of a single channel of the 3D measurement device - i.e. the field probe (in this case an one-dimensional induction coil), plus the relevant signal conditioning circuit, plus a data-logger system - is displayed in Fig. 6.2; a picture of the electronic signal conditioning circuit of a single channel is shown in Fig. 6.3.


Figure 6.2: Block diagram of a single channel of the measurement device of the magnetic induction field ( $X, Y, Z$ denotes that this channel is repeated over all the three Cartesian components of the field) isotropic 3D coil employed as a magnetic field probe.


Figure 6.3: Picture of the electronic signal conditioning circuit of a single channel of the 3D measurement device

The rms magnetic induction B under exam is converted into a voltage signal by the induction coil $L$. According to the European legislation, that imposes constraints on the coil surface limits, a rectangular coil with 100 turns and an area of $\sim 0.01 \mathrm{~m}^{2}$ has been selected.

The transfer function $V=f(B)$ results from the fundamental Faraday's law of induction:

$$
\begin{gather*}
v(t)=-N \cdot \frac{\mathrm{~d} \Phi(t)}{\mathrm{d} t}=-N \cdot \frac{\mathrm{~d}}{\mathrm{~d} t}\left[\int_{A} \vec{B}(t) \cdot \vec{n} \cdot \mathrm{~d} A\right]=-N \cdot \frac{\mathrm{~d}}{\mathrm{~d} t}[B(t) \cdot A]  \tag{6.1}\\
V_{R M S}=N \cdot \frac{2 \pi}{\sqrt{2}} \cdot f \cdot \hat{B} \cdot A=N \cdot 2 \pi \cdot f \cdot B \cdot A \tag{6.2}
\end{gather*}
$$

where $\Phi$ is the magnetic flux passing through the coil with an area $A$ and a number of turns $N, B$ is the rms value of magnetic induction and $f$ is the current frequency.

By substituting $A=0.01 \mathrm{~m}^{2}, N=100, f=50 \mathrm{~Hz}$ and by assuming the minimum measured value of magnetic induction $B=0.1 \mu \mathrm{~T}$, equation (6.1) gives a voltage of $V$
$=31 \mu \mathrm{~V}$ : therefore, the integrated electronic circuit must be able to accept a minimum input voltage of $31 \mu \mathrm{~V}$.

This signal is amplified by the first stage A1 of Fig. 6.2, which has two possible gain settings - 1 and 100 - and also includes a simple circuit for overvoltage protection.

In order to restore the waveform of the magnetic field and eliminate the dependence of equation (6.2) on frequency, the inverse of the mathematical operation carried out in (6.2) should be performed, namely integration, by introducing an integrator with variable gain (to permit calibration of the system). The integrator circuit ensures that the signal frequency response is flat in the frequency range concerned (from a few Hz to 400 Hz for low frequency magnetic field measurements according to [50]). A True RMS to DC Converter follows, that generates a dc output equal to the rms value of the input signal and finally, since the maximum expected amplitude of the output signal is approximately 200 mV , a second amplifier stage A2 has been added with gain 10 .

Each channel - shown in Fig. 6.2 - of the 3D measurement device of the magnetic induction field here developed has to be tested and calibrated in order to estimate the device proportionality coefficient between applied current and generated magnetic field, by checking for linearity and comparing the reading of the instrument with other commercial instrumentation available. After all these tests have given a positive outcome, the whole instrument is finally assembled (see Fig. 6.4), achieving in this way a three-axial probe that performs 3D field measurements and stores the three field components values.


Figure 6.4: $3 D$ 'smart' measurement system.

The circuit for testing the performance of the proposed measurement device consists of a Keithley AC current source plus a "Helmholtz" system, made of a square coil with a side length of 1 m that ensures an internal region of nearly uniform magnetic field; the field points of interest are located in the coil center. The value and the direction of the magnetic induction field in the center of the Helmholtz coil are known, since the value of the current applied is known, and are used for the calibration of the measuring instrument. Thus, by adjusting the current value through the circuit shown in Fig. 6.5, it is possible to adjust the value and direction of magnetic induction at the center of the Helmholtz coil.


Figure 6.5: Testing performance circuit of the measurement device.
Applying the Biot-Savart law to the magnetic circuit represented by the "Helmholtz" system, the relationship between the current flowing in the system and the value of the magnetic induction generated was obtained. Taking into account the uncertainty due to the mechanical design, we have:

$$
\begin{equation*}
B_{\text {ref }}=1.14 \mu T \cdot A \pm 0.4 \% \tag{6.3}
\end{equation*}
$$

This result was verified with measurements made using as a reference instrument the Wandel \& Goltermann mod. EFA 3.

In order to calibrate each channel of the 3D measurement device proposed, different measurements were carried out by setting different rms values of 50 Hz current in the "Helmholtz" system, ranging from 0.175 A to 3.51 A in steps of 0.175 A , thereby yielding rms values of magnetic field ranging from $0.2 \mu \mathrm{~T}$ to $4.0 \mu \mathrm{~T}$ (the so-called reference magnetic field values, $B_{r e f}$ ). The quantities to be measured (coil current and output voltage of the signal conditioning system applied to the coil that act as probe) were acquired using a National Instruments board (model NI 9239, with 4 simultaneous input channels of $24 \mathrm{bit}, 50 \mathrm{kS} / \mathrm{s}$ ) managed by an appropriate software for automatic processing of the acquired data, thereby achieving the measured magnetic field values, $B_{\text {meas }}$. The results were then verified by a further series of
measures carried out by varying the magnetic field from $0.1 \mu \mathrm{~T}$ to $2.0 \mu \mathrm{~T}$ in steps of $0.1 \mu \mathrm{~T}$. The results obtained are depicted in Figs. 6.6 and 6.7 and have shown that the measurement uncertainty (with coverage factor $k=2$ ) is $0.5 \%$ maximum for all three systems.


Figure 6.6: Comparison between reference (solid line) and measured (crosses) values of the magnetic field.


Figure 6.7: Measurement uncertainty from the comparison between reference and measured values of the magnetic field.

To improve the characterization of the measurement system, an estimate $S$ of the standard deviation of the measurements performed using the proposed measurement device has also been attained as follows:

1) several field measurements have been carried out in sequence with the same value of pre-set current in the "Helmholtz" system and the relevant standard deviation has been estimated (see standard deviation of measured values of the magnetic field, red dots in Fig. 6.8);
2) these field measurements have been compared with the field values calculated simultaneously from synchronous current measurements, and the relevant standard deviation has been estimated (see standard deviation of reference values of the magnetic field, blue dots in Fig. 6.8);
3) points 1) and 2) have been repeated by varying the pre-set magnetic induction values in the range from $0.25 \mu \mathrm{~T}$ to $4 \mu \mathrm{~T}$.


Figure 6.8: Standard deviation of reference (blue points) and measured (red points) values of the magnetic field.

As it can be argued from Fig. 6.8, very similar amplitude values of estimated standard deviation have been obtained for reference and measured field values. This means that the variation in the measurement results does not depend on the measuring system itself, but on the power supply of the "Helmholtz" system. Indeed, since the power supply depends directly on the instant value of the grid voltage, that in turn varies randomly with the load demand, it is unable to keep the current strictly constant over time in the "Helmholtz" system, that acts as reference magnetic field source.

### 6.3 FIRST TESTING OF THE MEASUREMENT DEVICE

### 6.3.1 Measurements in Situ

The performance of the proposed measurement device has been tested in situ. Figure 6.9 is an actual photo of the area of interest and Fig. 6.10 is its schematization.


Figure 6.9: Photo of the double-circuit power line close to the measurement point.


Figure 6.10: Schematization of the double-circuit power line tower and of the measurement point nearby.

The measurements were performed in the proximity of a double-circuit overhead MV power line of known geometry, in correspondence of the half span between one pole and the adjacent transformation cabin. This fact is depicted in the map of the area of interest in Fig. 6.11, and is the main reason why this particular point was chosen since it can guarantee the presence of a significant contribute of the magnetic field on the $z$-axis. The two circuits consist of bare conductors with ampacity $I_{1}=285 \mathrm{~A}$ and $I_{2}=230 \mathrm{~A}$, attached on opposite sides of the poles and placed at the same height from the ground $(\sim 9 \mathrm{~m})$. Measurements took place at 6 m from the ground and at 5 m distance from the pole.


Figure 6.11: Map of the area of interest.

The magnetic field components have been monitored and recorded every 5 min for a period of 24 h . The instant values of the line currents have been provided by the distribution utility. The measurement data are illustrated in Fig. 6.12, which displays the plots of the currents vs. time, and in Fig. 6.13, which displays the plots of the magnetic field components vs. time. In Fig. 6.13 the value of $B_{\text {tot }}$ is calculated by the composition of the three components, $B_{x}, B_{y}$ and $B_{z}$ using the relationship:

$$
\begin{equation*}
B=\sqrt{B_{x}^{2}+B_{y}^{2}+B_{z}^{2}} \tag{6.4}
\end{equation*}
$$



Figure 6.12: 24 h monitored instant values of the currents $I_{1}$ and $I_{2}$.


Figure 6.13: 24 h measurements of the magnetic field components.
$B_{\text {tot }}$ calculated via the three components according to eq. (6.4) is also displayed.

### 6.3.2 Data Processing

The handling and storage of measurement results are managed by a PLC in order to be subsequently used for extrapolating the functional relationship between the applied current and the generated magnetic field components.

Coming to this stage of the investigation, as previously argued, in the presence of multiple (say, $M$ ) source currents the vector composition of the total magnetic field makes the linear dependence of the total magnetic field $B$ on the currents unlikely to happen in general, whereas a linear functional relationships can be expected to hold between source currents and the components of the magnetic field. Following this procedure, expressions like the following can be derived:

$$
\begin{align*}
& B_{x}=\sum_{i=1}^{M} k_{x, i} I_{i}  \tag{6.5.a}\\
& B_{y}=\sum_{i=1}^{M} k_{y, i} I_{i}  \tag{6.5.b}\\
& B_{z}=\sum_{i=1}^{M} k_{z, i} I_{i} \tag{6.5.c}
\end{align*}
$$

Equations (6.5) have been validated experimentally using the measured data of $B$ field components (see Fig. 6.13) in Matlab ${ }^{\circledR}$ environment and the numerical values of coefficients $k_{x, 1}, k_{x, 2}, k_{y, 1}, k_{y, 2}, k_{z, 1}$ and $k_{z, 2}$, (since $i=1,2$ in our case), have been derived by the procedure of multiple linear regression. Doing so, the following expressions have been established:

$$
\begin{gather*}
B_{x}=k_{x, 1} I_{1}+k_{x, 2} I_{2} \Rightarrow B_{x}=0.0008 \cdot I_{1}+0.0012 \cdot I_{2}  \tag{6.6.a}\\
B_{y}=k_{y, 1} I_{1}+k_{y, 2} I_{2} \Rightarrow B_{y}=0.0059 \cdot I_{1}+0.0091 \cdot I_{2}  \tag{6.6.b}\\
B_{z}=k_{z, 1} I_{1}+k_{z, 2} I_{2} \Rightarrow B_{z}=0.0005 \cdot I_{1}+0.0008 \cdot I_{2} \tag{6.6.c}
\end{gather*}
$$

Figs. 6.14 are the results of the correlation analysis and show the trends of the components $B_{x, \text { measured }}, B_{y, \text { measured }}$ and $B_{z, \text { measured }}$ Vs. $B_{x}, B_{y}$ and $B_{z}$ calculated.


Figure 6.14.a: Correlation analysis of $B_{x}$ measured vs. $B_{x}$ calculated.


Figure 6.14.b: Correlation analysis of $B_{y}$ measured vs. $B_{y}$ calculated.


Figure 6.14.c: Correlation analysis of $B_{z}$ measured vs. $B_{z}$ calculated.

The knowledge of these coefficients enables the estimation of the magnetic field under load conditions different from those encountered during the measurement, in particular allows the estimation of the three components of the magnetic field vector at the point of maximum exposure and also allows the comparison of that value with the legal limits. The method of estimating the magnetic field described above is called "indirect method" and the Italian legislation has set boundaries in its use: the maximum value of the current should be the ampacity of conductors, and the correlation index between current and magnetic field should not less than 0.9 . In the case of multiple linear regression, the Italian law has not yet been updated since this is an innovative approach. However, the correlation coefficients of the multiple linear regression applied to the field components in this paper, result greater than $96 \%$ for all three components, fact that indicates that the regression lines approximate satisfactorily the real data points.

The relations used are listed below:

$$
\begin{align*}
& B_{x, \text { max }}=k_{x, 1} I_{1, \text { max }}+k_{x, 2} I_{2, \text { max }} \Rightarrow B_{x, \text { max }}=0.504 \mu T  \tag{6.7.a}\\
& B_{y, \text { max }}=k_{y, 1} I_{1, \text { max }}+k_{y, 2} I_{2, \text { max }} \Rightarrow B_{y, \text { max }}=3.775 \mu T  \tag{6.7.b}\\
& B_{z, \text { max }}=k_{z, 1} I_{1, \text { max }}+k_{z, 2} I_{2, \max } \Rightarrow B_{z, \text { max }}=0.327 \mu T  \tag{6.7.c}\\
& B_{\max }=\sqrt{B_{x, \text { max }}^{2}+B_{y, \text { max }}^{2}+B_{z, \text { max }}^{2}} \Rightarrow B_{\max }=3.823 \mu T \tag{6.8}
\end{align*}
$$

with $I_{1, \max }=285 \mathrm{~A}, I_{2, \max }=230 \mathrm{~A}, k_{x, 1}=0.0008, k_{x, 2}=0.0012, k_{y, 1}=0.0059, k_{y, 2}=0.0091$, $k_{z, 1}=0.0005, k_{z, 2}=0.0008$.

### 6.4 SECOND TESTING OF THE MEASUREMENT DEVICE

The novelty here is that the measurement device and the multilinear regression algorithm are applied to a three-phase single-circuit power line where the timedependent rms values of each phase are known. In practice, power systems rarely have perfectly balanced loads, currents, voltages and impedances in all three phases, so the particular interest of this case is to examine if the three-phase currents are equilibrated and in case that they are not, which are the consequences in terms of the magnetic field.

### 6.4.1 Measurements in Situ

A second campaign of measurements in situ has been effectuated for testing the proposed measurement device performance. Fig. 6.15 is an actual photo of the area of interest and Fig. 6.16 is its schematization.


Figure 6.15: Photo of the single-circuit power line close to the measurement point.


Figure 6.16: Schematization of the single-circuit power line tower and of the measurement point nearby.

The measurements were performed in the proximity of a single-circuit overhead HV power line of known geometry, near the half span between two adjacent towers
(see Fig. 6.17). The single-circuit consists of 3 bare conductors (phases A, B, C) with ampacity $I=385$ A and attached on opposite sides of the pole (see Fig. 6.16). Phase A is placed at 10 m height from the ground and at 4.45 m distance from the pole, phase B is placed at 12.25 m height from the ground and at 3.60 m distance from the pole and phase C is placed at 14.5 m height from the ground and at 3.45 m distance from the pole. Measurements took place at 6 m from the ground and at 8.5 m distance from the line.


Figure 6.17: Map of the area of interest.
The time-dependent rms values of the line currents have been provided by the distribution utility (see Fig. 6.18). The magnetic field components have been monitored and recorded every 1 min for a period of 24 h and in Fig. 6.19 their plots vs. time are displayed along with the plot of $B_{\text {tot }}$ vs. time, calculated by the composition of the three components, $B_{x}, B_{y}$ and $B_{z}$ using the relationship (6.4).


Figure 6.18: $24 h$ monitored instant values of the currents $I_{A}, I_{B}$ and $I_{C}$.


Figure 6.19: 24 h measurements of the magnetic field components. $B_{\text {tot }}$ calculated via the three components according to eq. (6.4) is also displayed.

Elaborating the currents data in Matlab ${ }^{\mathrm{TM}}$ environment it was seen that the threephase currents are not equilibrated since the vector sum of the line currents does not result zero. The estimated values of that non-zero current are shown in Fig. 6.20 and its median value results $I_{\mathrm{N}, \text { median }}=0.919 \mathrm{~A}$.


Figure 6.20: 24 h calculated neutral current.

### 6.4.2 Data Processing

The handling and storage of measurement results are managed by a Laptop with a data acquisition card NI USB 9239 ( 4 simultaneous channels of 24 bits, $50 \mathrm{kS} / \mathrm{s}$ per channel). Subsequently, the functional relationships between the applied currents and the generated magnetic field components, eqs. (6.5), have been validated experimentally using the measured data of $B$-field components (see Fig. 6.19) in Matlab ${ }^{\circledR}$ environment. The numerical values of coefficients $k_{x, 1}, k_{x, 2}, k_{x, 3}, k_{y, 1}, k_{y, 2}, k_{y, 3}$, $k_{z, 1}, k_{z, 2}$, and $k_{z, 3}$, (since $i=1,2,3$ in our case), have been derived by the procedure of multiple linear regression and the following expressions have been established:

$$
\begin{equation*}
B_{x}=k_{x, 1} I_{A}+k_{x, 2} I_{B}+k_{x, 3} I_{C} \Rightarrow B_{x}=0.0014 \cdot I_{A}+0.0206 \cdot I_{B}-0.0115 \cdot I_{C} \tag{6.9.a}
\end{equation*}
$$

$$
\begin{equation*}
B_{y}=k_{y, 1} I_{A}+k_{y, 2} I_{B}+k_{y, 3} I_{C} \Rightarrow B_{y}=0.0080 \cdot I_{A}+0.0162 \cdot I_{B}-0.0154 \cdot I_{C} \tag{6.9.b}
\end{equation*}
$$

$$
\begin{equation*}
B_{z}=k_{z, 1} I_{A}+k_{z, 2} I_{B}+k_{z, 2} I_{C} \Rightarrow B_{z}=0.0009 \cdot I_{A}+0.0024 \cdot I_{B}-0.0021 \cdot I_{C} \tag{6.9.c}
\end{equation*}
$$

Figs. 6.21 are the results of the correlation analysis and show the trends of the components $B_{x, \text { measured }}, B_{y, \text { measured }}$ and $B_{z, \text { measured }}$ Vs. $B_{x}, B_{y}$ and $B_{z}$ calculated.


Figure 6.21.a: Correlation analysis of $B_{x}$ measured vs. $B_{x}$ calculated.


Figure 6.21.b: Correlation analysis of $B_{y}$ measured vs. $B_{y}$ calculated.


Figure 6.21.c: Correlation analysis of $B_{z}$ measured vs. $B_{z}$ calculated.
Using the "indirect method", described previously, the estimation of the three components of the magnetic field vector at the point of maximum exposure for the comparison of that value with the legal limits, has been achieved. The results are depicted in the following relationships:

$$
\begin{align*}
& B_{x, \text { max }}=k_{x, 1} I_{A, \text { max }}+k_{x, 2} I_{B, \text { max }}+k_{x, 3} I_{C, \text { max }} \Rightarrow B_{x, \text { max }}=4.04 \mu T  \tag{6.10.a}\\
& B_{y, \text { max }}=k_{y, 1} I_{A, \text { max }}+k_{y, 2} I_{B, \text { max }}+k_{y, 3} I_{C, \text { max }} \Rightarrow B_{y, \text { max }}=3.39 \mu T \tag{6.10.b}
\end{align*}
$$

$$
\begin{gather*}
B_{z, \text { max }}=k_{z, 1} I_{A, \text { max }}+k_{z, 2} I_{B, \text { max }}+k_{z, 3} I_{C, \text { max }} \Rightarrow B_{z, \text { max }}=0.46 \mu T  \tag{6.10.c}\\
B_{\max }=\sqrt{B_{x, \text { max }}^{2}+B_{y, \text { max }}^{2}+B_{z, \text { max }}^{2}} \Rightarrow B_{\max }=5.29 \mu T \tag{6.11}
\end{gather*}
$$

with $I_{\mathrm{A}, \max }=385 \mathrm{~A}, I_{\mathrm{B}, \max }=385 \mathrm{~A}, I_{\mathrm{C}, \max }=385 \mathrm{~A}, k_{x, I}=0.0014, k_{x, 2}=0.0206, k_{x, 3}=-$ $0.0115, k_{y, 1}=0.0080, k_{y, 2}=0.0162, k_{y, 3}=-0.0154, k_{z, 1}=0.0009, k_{z, 2}=0.0024$ and $k_{z, 3}=$ -0.0021.

### 6.4 CONCLUSIONS

The construction, calibration and characterization of the innovative 3D measurement and evaluation system of the magnetic field has enabled the measuring, recording and storing of the rms values of the magnetic field components in a Cartesian reference system, $B_{x}, B_{y}$ and $B_{z}$, and its performance has been successfully tested in laboratory and in situ.

Through 24 h continuous measurements and acquiring the load curves of the double-circuit three-phase line under exam, the coefficients that describe the proportionality relationship between magnetic induction components and intensity of the currents flowing in the two lines were estimated. The same procedure was followed also for the case of a single-circuit three-phase line, where the coefficients that describe the proportionality relationship between magnetic induction components and intensity of the each phase current flowing in the line were estimated.

The correlation coefficients of the multiple linear regression applied to the field components in the case of the double-circuit three-phase line, result greater than $96 \%$ for all three components, fact that indicates that the regression lines approximate satisfactorily the real data points. On the other hand, the correlation coefficients of the multiple linear regression applied to each field component in the case of the singlecircuit three-phase line, result approximately $66 \%$, fact that indicates that the regression lines approximate satisfactorily enough the real data points but not as well as in the first case. Such a result was expected since the three-phase line was not balanced and a non-zero current flowing in the neutral was present.

The knowledge of these coefficients enable us to estimate the magnetic field under load conditions different from those encountered during the measurement, in particular allows the estimation of the three components of the magnetic field vector at the point of maximum exposure and also allows the comparison of that value with the legal limits. In the case of the double-circuit three-phase line, the magnetic field can arrive up to maximum $3.82 \mu \mathrm{~T}$ and in the case of the single-circuit three-phase line, the magnetic field can arrive up to maximum $5.29 \mu \mathrm{~T}$. In both cases, the maximum magnetic field value is within the legal limits (in particular below the attention value of $10 \mu \mathrm{~T}$ ) as set by the Italian law for existing lines.

The method of estimating the magnetic field described above is called "indirect method" and the Italian legislation has set boundaries in its use: the maximum value of the current should be the ampacity of conductors, and the correlation index between current and magnetic field should not less than 0.9 . In the case of multiple linear regression, the Italian law has not yet been updated since this is an innovative approach. However, the correlation coefficients of the multiple linear regression applied to the field components in the case of the double-circuit three-phase line, result greater than $96 \%$ for all three components. This is not true though for the single-circuit three-phase line, where the correlation coefficients of the multiple linear regression applied to the field components result approximately $66 \%$ for all three components. Even if this value is less than $90 \%$, as defined by the Italian law, still remains an acceptable value as far as the results multiple linear regression are concerned.

Finally, it can be argued, that the "indirect method" of estimating the magnetic field has proven to be more accurate and reliable for multiple and complex 3D field sources, where the exact geometry of the sources is unknown or not accessible and even powerful 3D codes might not work: where a 3D code fails an empirically-derived multi-linear relationship can succeed.

## CHAPTER 7

## Conclusions

The calculation of the human exposure to magnetic field generated by complex configurations of the field sources was the subject of this PhD Thesis. The need of precise evaluation of the magnetic field arises from the public concern about its impact near sensitive receptors and also for the determination of the safety distances corresponding to a maximum limit value of the rms magnetic induction, as defined by the national and international legislation.

More precisely, the main object of the scientific research made included the calculation of the magnetic field generated by twisted configurations of the conductors, since this is the main trend in the electric power distribution networks. This helical configuration provides a drastic reduction of the magnetic field generated and - both for this and for other practical reasons - twisted three-phase cables are more and more often the preferred solution for LV and MV power lines. This holds not only close to residential areas, but also in rural and wilderness areas, where twisted three-phase cables are often chosen as a replacement of existing traditional bare conductors overhead lines, since their external insulation protects them against external faults. Additionally, it must be pointed out that twisted threephase cables are the environmental friendly solution used for connecting "green energy" production systems, such as photovoltaic systems and wind-generators, to the grid.

The exact and approximate theory for the calculation of the magnetic field generated by a twisted three-phase cable configuration was presented. The literature approximation provides acceptable results for large distances from the helix axis only, whereas for distances close to the twisted cable it exhibits large errors compared to the exact rigorous expression of the magnetic field. For this reason, a parametric heuristic analysis was performed that results in an innovative simplified expression of the rms magnetic induction as a function of the distance from the helix axis. This innovative expression approximates the logarithm of $B$ with the equation of a straight line plus a hyperbolic term for simulating the deviation from linearity close to the helix axis. The coefficients of the straight line equation plus the hyperbolic term depend on the pitch and radius of the helix, and can be expressed as pure numbers $-1^{\text {st }}$ level of approximation - or as analytical functions of the pitch and radius $-2^{\text {nd }}$ level of approximation. The great advantage of the $2^{\text {nd }}$ level of approximation is the total explanation of the dependence of $B$ on pitch and radius, without the need to know the exact values of those coefficients, but its disadvantage
is that the higher level of approximation has inevitably brought a higher error. The effectiveness of the innovative expression is evaluated by some numerical simulations relevant to typical MV cable designs and the comparison between the results provided by the exact and the approximate formulae from the literature with those obtained via the innovative simplified formula has shown that the error involved by the innovative simplified formula in $1^{\text {st }}$ level of approximation is always smaller than that brought about by the approximate formula from the literature, both for small and for large distances from the helix axis and always positive, providing a conservative estimation of the magnetic field. It can be said that the innovative formula in the $1^{\text {st }}$ level of approximation results definitely much simpler than the exact one and provides a much smaller relative error compared to the approximated one from the literature, especially for small distances from the helix axis. As to the $2^{\text {nd }}$ level of approximation, similar considerations hold, apart some field point ranges where the absolute error of the simplified innovative formula in the $2^{\text {nd }}$ level of approximation is practically the same in absolute terms as the error of the approximate formula from the literature: this is the price paid to the further level of approximation. However, in these ranges (and mostly elsewhere) the $2^{\text {nd }}$ level of approximation overestimates $B$, contrarily to the approximated one from the literature; hence - as the $1^{\text {st }}$ level of approximation - also the $2^{\text {nd }}$ level provides a conservative estimate with respect to the literature approximation.

Subsequently, the magnetic field generated by an overhead or underground double-circuit twisted three-phase power cable line was treated following two different approaches; the exact calculus which effectuates a complex vector sum of the two vectors of the magnetic field from each twisted three-phase single-circuit and the 'worst case' calculus which implies the algebraic sum of the two individual values of the magnetic field from each twisted three-phase single-circuit. The exact vector sum has resulted a rather complex procedure since computational hypotheses were set - such as the constancy of the pitch of the single helix, as well as the constancy of the perfect parallelism of the two helixes. Moreover, conventions that would facilitate the analysis were introduced, since it is impossible to obtain exact information about the geometrical arrangement of the power line in practice. On the other hand, the 'worst case' assumption for the calculation of the magnetic field is a simple approximated calculus that yields a good interpretation of the reality. Additionally, the 'worst case' approach permits further approximation by using both the approximated formula from the literature and the simplified innovative one, in both levels of approximation. The simulations made using a typical MV cable have concluded that the relative percent error of the innovative simplified formula with respect to the approximated one from the literature is always smaller and positive, confirming the better performances of the innovative simplified formula in the 'worst case' calculus.

Comparing the calculations with results from actual measurements relevant to a three-phase double-circuit twisted power cable line carried out in situ, it can be
deducted that the theoretical approach of the double-circuit twisted cable line is correct and provides - especially when considering the difficult circumstances under which the experiment took place - magnetic field values near to the real ones, particularly for field points where the magnetic induction is large enough for being detected by measurement instruments. Moreover, at these same field points the expression $B_{W C, \text { simp }}$, calculated using the innovative simplified formula, can provide a "quick-and-easy" way of calculating the generated magnetic field. The simplicity of this innovative expression lies on the fact that it is the algebraic sum of the two rms values of the total field $B$ generated by each twisted cable that are calculated without using sophisticated mathematics. In this way, the calculation of the magnetic field generated by any multiple-circuit twisted three-phase power cable line becomes possible in the framework of a worst-case approach.

By generalizing the vector theory concerning a double-circuit twisted three-phase power cable line to a multiple-circuit twisted three-phase power cable line - the standard lines for transporting renewable energy - the exact calculation of the magnetic field generated by any number of circuits has been calculated. Also, in the framework of the 'worst-case' approach, the analysis has been extended to multiplecircuit twisted three-phase cables using both the exact and approximated formula derived from the literature about single-circuits and the simple innovative formula developed in both levels of approximation. The simulation results have proven, once more, that the smaller and positive relative error of the simplified innovative formula compared to the approximated one from the literature, make expressions $B_{W C, \text { simp }}$ an efficient alternative way of calculating the magnetic field of a multiple-circuit twisted three-phase cable line. This is due to its advantages which lie on its simplicity - an algebraic sum of the rms values of the total field $B$ generated by each twisted cable, calculated without using sophisticated mathematics - and on its reliability - the maximum error presented is small for both levels of approximation and positive, providing a conservative estimation of the magnetic field.

Dealing now with complex configurations of overhead power lines, an appropriate software was created for the exact three-dimensional calculation of the generated magnetic field, following the guidelines of the Italian law, taking into consideration the real geometrical configuration of the lines. This has permitted the exact calculation of the magnetic field at field points of interest near the overhead power lines without overestimating the field as existing two-dimensional magnetic field calculating programs do. For the sake of comparison a two-dimensional calculating program has also been created, following the guidelines of the Italian law, and the results of all the simulations made were satisfactorily compared with 2D magnetic field calculation programs commercially available. The results of all 3D simulations made concerning single- and double-circuit overhead power lines by the threedimensional code created were compared with a powerful 3D magnetic field calculation program (unfortunately not commercially available and classified).

Nevertheless, critical observations concerning the correctness of the script were made, by comparing the two-dimensional results with the three-dimensional ones.

Finally, for complex three-dimensional configurations of the field sources, where their exact geometry is unknown or not accessible and even powerful 3D codes might not work, the magnetic field has been successfully estimated via the so-called "indirect method" and via the use of an innovative 'smart' measurement instrument. The construction, calibration and characterization of the innovative 3D measurement and evaluation system of the magnetic field has enabled measuring and recording the rms values of the magnetic field components in a Cartesian reference system, $B_{x}, B_{y}$ and $B_{z}$, and its performance has been successfully tested in laboratory and in situ. The post-processing of the measurement results has permitted the calculation of the coefficients that describe the proportionality relationship between magnetic induction components and intensity of the currents via multilinear regression techniques appropriately created. The knowledge of these coefficients enables the estimation of the magnetic field under any load conditions and in particular allows the estimation of the three components of the magnetic field vector at the point of maximum exposure for controlling whether this value lies within the legal limits. This method of estimating the magnetic field described above is called "indirect method" and has proven to be more accurate and reliable for multiple and complex 3D field sources: where a 3D code fails an empirically-derived multi-linear relationship may succeed.

Concluding, the purpose of this PhD Thesis, that was the analysis and calculation of the magnetic field generated by complex configurations of electric power systems, has been successfully achieved by:
$\checkmark$ firstly, presenting a complete exact theory regarding a single-, double- and multiple-circuit twisted three-phase power cable lines along with a simplified approach for the magnetic field calculation generated by these configurations proposing an innovative formula derived from a parametrical analysis;
$\checkmark$ secondly, creating appropriate 3D numerical codes for simulating real existing overhead power lines and calculating the magnetic field in their vicinity;
$\checkmark$ finally proposing an innovative 'smart' measurement and evaluation system, which measures and records the rms magnetic field components values for finding the proportionality coefficients between them and the intensity of the currents via multilinear regression techniques for the final extrapolation of the maximum value of the field at the point of interest ("indirect method").

## APPENDIX 1

## CALCULATION OF COEFFICIENTS $a_{0}, a_{1}, a_{2}, a_{3}$

The unknown parameters of the innovative approximate expression (eq. 2.23) discussed in Chapter 2, are $a_{0}, a_{1}, a_{2}$ and $a_{3}$. These parameters are $\alpha$ and $p$ functions, and they have been derived by best-fitting techniques developed in Matlab ${ }^{\mathrm{TM}}$.

The helix radius depends from the geometrical arrangement of the three conductors as shown in Fig. A1.1, where the section of a three-phase twisted cable is depicted along with the quantities $D_{\text {out }}, R_{\text {out }}$ (cable outer diameter and radius), $D_{\text {in }}, R_{\text {in }}$ (single conductor diameter and radius) and the helix radius $\alpha$.


Figure A1.1: Three-phase twisted cable section.
The helix radius derives from the relationship:

$$
\begin{equation*}
\alpha=D_{i n} \cdot \sqrt{3} \tag{A1.1}
\end{equation*}
$$

Even though the cable manufacturing companies provide the value of $D_{i n}$, in their datasheets, it is preferable use the value of $D_{\text {out }}$ for calculating helix radius $\alpha$ : this is why, $D_{i n}$ is usually an approximate value without taking into consideration the conductor insulation thickness and, in cases of missing cable datasheet, it is easier acquiring the $D_{\text {out }}$ value simply by measuring it. So, after some elementary mathematical processes, the following expression has resulted that gives the value $D_{\text {in }}$ as a function of $D_{\text {out }}$ :

$$
\begin{equation*}
D_{\text {in }}=0.4641 \cdot D_{\text {out }} \tag{A1.2}
\end{equation*}
$$

As far as the helix pitch value is concerned, there is a general empirical rule that says:

$$
\begin{equation*}
p=40 \cdot D_{i n} \tag{A1.3}
\end{equation*}
$$

but is not always applied and it is better if every case is treated using the real, measured value of the helix pitch.

Subsequently, the cases of the cables ARE4H1RX $12 / 20 \mathrm{kV}$ [64], ARG7H1RX $12 / 20 \mathrm{kV}$ [63] and RG7H1OR $12 / 20 \mathrm{kV}$ [66] are treated, providing two Tables, for every existing cross-section of these cables, one with their construction and electrical characteristics, and the other with the values of the coefficients $a_{0}, a_{1}, a_{2}$, and $a_{3}$ with respect to some practical helix pitch values. The calculated helix radius value is also provided.

## 1) ARE4H1RX 12/20 kV

Table A1.1.a: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 \times 3 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 7.0 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 56.1 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 1930 |
| min bending radius | $[\mathrm{mm}]$ | 520 |
| open air installation | $[\mathrm{A}]$ | 154 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 147 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 112 |

Table A1.1.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{I}}$ | coefficient <br> $\boldsymbol{a}_{\mathbf{2}}$ | coefficient <br> $\boldsymbol{a}_{\mathbf{3}}$ |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 0}$ | -0.6204 | -6.6444 | 0.1 | 1.69 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 2}$ | -0.5014 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 4}$ | -0.3957 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 6}$ | -0.2997 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 8}$ | -0.2111 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{2 . 0}$ | -0.1283 | -3.5512 | 0.1 | 1.93 |

Table A1.2.a: Cable ARE4HIRX 12/20 kV $3 x 50 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{\mathbf{2}}\right]$ | $\mathbf{3 \times 5 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 8.2 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 58.8 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 2140 |
| min bending radius | $[\mathrm{mm}]$ | 540 |
| open air installation | $[\mathrm{A}]$ | 185 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 174 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 131 |

Table A1.2.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 50 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\mathbf{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{1 . 0}$ | -0.6199 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{1 . 2}$ | -0.5011 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{1 . 4}$ | -0.3955 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{1 . 6}$ | -0.2995 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{1 . 8}$ | -0.2109 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 6}$ | $\mathbf{2 . 0}$ | -0.1282 | -3.5512 | 0.1 | 1.93 |

Table A1.3.a: Cable ARE4H1RX 12/20 kV $3 x 70 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{\mathbf{2}}\right]$ | $\mathbf{3 x 7 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 9.7 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 62.6 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 2490 |
| min bending radius | $[\mathrm{mm}]$ | 580 |
| open air installation | $[\mathrm{A}]$ | 230 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 212 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 160 |

Table A1.3.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 70 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 0}$ | -0.6194 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 2}$ | -0.5008 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 4}$ | -0.3952 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 6}$ | -0.2993 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 8}$ | -0.2108 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{2 . 0}$ | -0.1280 | -3.5513 | 0.1 | 1.93 |

Table A1.4.a: Cable ARE4HIRX 12/20 kV $3 x 95 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 x 9 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 11.4 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 66.3 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 2860 |
| min bending radius | $[\mathrm{mm}]$ | 610 |
| open air installation | $[\mathrm{A}]$ | 280 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 253 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 190 |

Table A1.4.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 95 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient | $\boldsymbol{a}_{\boldsymbol{2}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |

Table A1.5.a: Cable ARE4H1RX 12/20 kV $3 x 120 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 2 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 12.9 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 70.2 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 3260 |
| min bending radius | $[\mathrm{mm}]$ | 650 |
| open air installation | $[\mathrm{A}]$ | 323 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 288 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 216 |

Table A1.5.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 120 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius | helix pitch |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | coefficient | $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 0}$ | -0.6184 | -6.6444 | 0.1 | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 2}$ | -0.5000 | -5.6066 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 4}$ | -0.3947 | -4.8682 | 0.1 | 1.76 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 6}$ | -0.2989 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 8}$ | -0.2104 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{2 . 0}$ | -0.1277 | -3.5513 | 0.1 | 1.93 |

Table A1.6.a: Cable ARE4H1RX 12/20 kV $3 x 150 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 \times 1 5 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 14.0 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 72.7 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 3560 |
| min bending radius | $[\mathrm{mm}]$ | 670 |
| open air installation | $[\mathrm{A}]$ | 365 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 322 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 241 |

Table A1.6.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 150 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 0}$ | -0.6184 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 2}$ | -0.5000 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 4}$ | -0.3947 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 6}$ | -0.2989 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 8}$ | -0.2104 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{2 . 0}$ | -0.1277 | -3.5513 | 0.1 | 1.93 |

Table A1.7.a: Cable ARE4H1RX 12/20 kV $3 x 185 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 8 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 15.8 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 77.2 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 4100 |
| min bending radius | $[\mathrm{mm}]$ | 720 |
| open air installation | $[\mathrm{A}]$ | 421 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 365 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 272 |

Table A1.7.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | | helix pitch |
| :---: |
| $\boldsymbol{p}[\mathbf{m}]$ | | coefficient |
| :---: |
| $\boldsymbol{a}_{\boldsymbol{0}}$ | | coefficient |
| :---: |
| $\boldsymbol{a}_{\boldsymbol{1}}$ | | coefficient |
| :---: | | $\boldsymbol{a}_{\boldsymbol{2}}$ |
| :---: |

Table A1.8.a: Cable ARE4H1RX 12/20 kV $3 x 240 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 x 2 4 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 18.2 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 82.6 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 4830 |
| min bending radius | $[\mathrm{mm}]$ | 770 |
| open air installation | $[\mathrm{A}]$ | 498 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 423 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 314 |

Table A1.8.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 240 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, a $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 0}$ | -0.6166 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 2}$ | -0.4987 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 4}$ | -0.3937 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 6}$ | -0.2982 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 8}$ | -0.2098 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{2 . 0}$ | -0.1272 | -3.5513 | 0.1 | 1.93 |

Table A1.9.a: Cable ARE4H1RX 12/20 kV $3 x 300 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 \times 3 0 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 20.8 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 89.8 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 5720 |
| min bending radius | $[\mathrm{mm}]$ | 740 |
| open air installation | $[\mathrm{A}]$ | 576 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 478 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 354 |

Table A1.9.b: Cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 300 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 0}$ | -0.6152 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 2}$ | -0.4978 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 4}$ | -0.3930 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 6}$ | -0.2976 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 8}$ | -0.2094 | -3.8904 | 0.1 | 1.90 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{2 . 0}$ | -0.1268 | -3.5513 | 0.1 | 1.93 |

## 2) ARG7H1RX $12 / 20 \mathrm{kV}$

Table A1.10.a: Cable ARG7H1RX 12/20 kV $3 x 70 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 x 7 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 9.8 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 67 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 3000 |
| min bending radius | $[\mathrm{mm}]$ | 720 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 200 |

Table A1.10.b: Cable ARG7H1RX 12/20 $\mathrm{kV} 3 x 70 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius | helix pitch <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{p}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient | $\boldsymbol{a}_{\boldsymbol{1}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | | coefficient |
| :---: | | $\boldsymbol{a}_{\boldsymbol{2}}$ |
| :---: |

Table A1.11.a: Cable ARG7H1RX 12/20 kV $3 x 120 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 2 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 13.1 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 74 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 4000 |
| min bending radius | $[\mathrm{mm}]$ | 790 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 280 |

Table A1.11.b: Cable ARG7H1RX $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{3}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{1 . 0}$ | -0.6178 | -6.6444 | 0.1 | 1.69 |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{1 . 2}$ | -0.4996 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{1 . 4}$ | -0.3944 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{1 . 6}$ | -0.2987 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{1 . 8}$ | -0.2102 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 2 0}$ | $\mathbf{2 . 0}$ | -0.1276 | -3.5513 | 0.1 | 1.93 |

Table A1.12.a: Cable ARG7H1RX 12/20 kV $3 x 185 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 8 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 16.1 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 81 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 4800 |
| min bending radius | $[\mathrm{mm}]$ | 860 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 360 |

Table A1.12.b: Cable ARG7H1RX $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 0}$ | -0.6166 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 2}$ | -0.4987 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 4}$ | -0.3937 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 6}$ | -0.2982 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 8}$ | -0.2098 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{2 . 0}$ | -0.1272 | -3.5513 | 0.1 | 1.93 |

## 3) RG7H1OR $12 / 20 \mathrm{kV}$

Table A1.13.a: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 25 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{\mathbf{2}}\right]$ | $\mathbf{3 \times 2 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 6.0 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 54.5 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 3550 |
| min bending radius | $[\mathrm{mm}]$ | 720 |
| open air installation | $[\mathrm{A}]$ | 153 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 150 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 118 |

Table A1.13.b: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 x 25 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\mathbf{2}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{3}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 0}$ | -0.6204 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 2}$ | -0.5014 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 4}$ | -0.3957 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 6}$ | -0.2997 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 8}$ | -0.2111 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{2 . 0}$ | -0.1283 | -3.5512 | 0.1 | 1.93 |

Table A1.14.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 x 3 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 7.0 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 54.7 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 3790 |
| min bending radius | $[\mathrm{mm}]$ | 720 |
| open air installation | $[\mathrm{A}]$ | 177 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 175 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 138 |

Table A1.14.b: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius $\alpha[\mathrm{m}]$ | $\begin{aligned} & \text { helix pitch } \\ & p[\mathrm{~m}] \end{aligned}$ | $\begin{gathered} \text { coefficient } \\ a_{0} \end{gathered}$ | $\begin{gathered} \text { coefficient } \\ a_{1} \end{gathered}$ | $\begin{gathered} \text { coefficient } \\ a_{2} \end{gathered}$ | $\begin{gathered} \text { coefficient } \\ a_{3} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.015 | 1.0 | -0.6204 | -6.6444 | 0.1 | 1.70 |
| 0.015 | 1.2 | -0.5014 | -5.6066 | 0.1 | 1.75 |
| 0.015 | 1.4 | -0.3957 | -4.8682 | 0.1 | 1.80 |
| 0.015 | 1.6 | -0.2997 | -4.3169 | 0.1 | 1.85 |
| 0.015 | 1.8 | -0.2111 | -3.8904 | 0.1 | 1.89 |
| 0.015 | 2.0 | -0.1283 | -3.5512 | 0.1 | 1.93 |

Table A1.15.a: Cable RG7H1OR $12 / 20$ kV $3 x 50 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 \times 5 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 8.2 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 57.4 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 4360 |
| min bending radius | $[\mathrm{mm}]$ | 760 |
| open air installation | $[\mathrm{A}]$ | 209 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 207 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 162 |

Table A1.15.b: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 50 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}, a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{3}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 0}$ | -0.6204 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 2}$ | -0.5014 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 4}$ | -0.3957 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 6}$ | -0.2997 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{1 . 8}$ | -0.2111 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 5}$ | $\mathbf{2 . 0}$ | -0.1283 | -3.5512 | 0.1 | 1.93 |

Table A1.16.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 x 70 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 x 7 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 9.9 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 62.0 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 5390 |
| min bending radius | $[\mathrm{mm}]$ | 820 |
| open air installation | $[\mathrm{A}]$ | 260 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 253 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 198 |

Table A1.16.b: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 70 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\mathbf{2}}$ | $\boldsymbol{a}_{\mathbf{3}}$ |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 0}$ | -0.6194 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 2}$ | -0.5008 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 4}$ | -0.3952 | -4.8682 | 0.1 | 1.80 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 6}$ | -0.2993 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{1 . 8}$ | -0.2108 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 7}$ | $\mathbf{2 . 0}$ | -0.1280 | -3.5513 | 0.1 | 1.93 |

Table A1.17.a: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 x 95 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{\mathbf{2}}\right]$ | $\mathbf{3 x 9 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 11.6 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 65.9 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 6470 |
| min bending radius | $[\mathrm{mm}]$ | 870 |
| open air installation | $[\mathrm{A}]$ | 315 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 300 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 234 |

Table A1.17.b: Cable RG7HIOR 12/20 $\mathrm{kV} 3 x 95 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{3}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{1 . 0}$ | -0.6189 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{1 . 2}$ | -0.5004 | -5.6066 | 0.1 | 1.75 |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{1 . 4}$ | -0.3949 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{1 . 6}$ | -0.2991 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{1 . 8}$ | -0.2106 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 8}$ | $\mathbf{2 . 0}$ | -0.1279 | -3.5513 | 0.1 | 1.93 |

Table A1.18.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 x 120 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 2 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 13.1 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 69.3 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 7470 |
| min bending radius | $[\mathrm{mm}]$ | 920 |
| open air installation | $[\mathrm{A}]$ | 362 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 342 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 266 |

Table A1.18.b: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 0}$ | -0.6184 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 2}$ | -0.5000 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 4}$ | -0.3947 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 6}$ | -0.2989 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 8}$ | -0.2104 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{2 . 0}$ | -0.1277 | -3.5513 | 0.1 | 1.93 |

Table A1.19.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 x 150 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 5 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 14.4 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 72.6 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 8540 |
| min bending radius | $[\mathrm{mm}]$ | 960 |
| open air installation | $[\mathrm{A}]$ | 408 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 381 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 296 |

Table A1.19.b: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 x 150 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius <br> $\boldsymbol{\alpha}[\mathbf{m}]$ | helix pitch <br> $\boldsymbol{p}[\mathbf{m}]$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{0}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ | coefficient <br> $\boldsymbol{a}_{\mathbf{3}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 0}$ | -0.6184 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 2}$ | -0.5000 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 4}$ | -0.3947 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 6}$ | -0.2989 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{1 . 8}$ | -0.2104 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 1 9}$ | $\mathbf{2 . 0}$ | -0.1277 | -3.5513 | 0.1 | 1.93 |

Table A1.20.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 x 185 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathbf{m m}^{2}\right]$ | $\mathbf{3 x 1 8 5}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 16.1 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 76.8 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 10020 |
| min bending radius | $[\mathrm{mm}]$ | 1020 |
| open air installation | $[\mathrm{A}]$ | 468 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 431 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 335 |

Table A1.20.b: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{1 . 0}$ | -0.6172 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{1 . 2}$ | -0.4992 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{1 . 4}$ | -0.3941 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{1 . 6}$ | -0.2984 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{1 . 8}$ | -0.2100 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 2 1}$ | $\mathbf{2 . 0}$ | -0.1274 | -3.5513 | 0.1 | 1.93 |

Table A1.21.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 240 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 \times 2 4 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 18.5 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 82.0 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 12090 |
| min bending radius | $[\mathrm{mm}]$ | 1090 |
| open air installation | $[\mathrm{A}]$ | 550 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 500 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 387 |

Table A1.21.b: Cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 240 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix pitch | coefficient | coefficient | coefficient | coefficient |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | $\boldsymbol{a}_{\boldsymbol{2}}$ | $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 0}$ | -0.6166 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 2}$ | -0.4987 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 4}$ | -0.3937 | -4.8682 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 6}$ | -0.2982 | -4.3169 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{1 . 8}$ | -0.2098 | -3.8904 | 0.1 | 1.89 |
| $\mathbf{0 . 0 2 2}$ | $\mathbf{2 . 0}$ | -0.1272 | -3.5513 | 0.1 | 1.93 |

Table A1.22.a: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 300 \mathrm{~mm}^{2}$ : construction and electrical characteristics.

| conductor cross-section | $\left[\mathrm{mm}^{2}\right]$ | $\mathbf{3 \times 3 0 0}$ |
| :---: | :---: | :---: |
| approx conductor diameter | $[\mathrm{mm}]$ | 21.1 |
| insulation thickness | $[\mathrm{mm}]$ | 5.5 |
| max outer diameter | $[\mathrm{mm}]$ | 88.7 |
| approx weight | $[\mathrm{kg} / \mathrm{km}]$ | 14620 |
| min bending radius | $[\mathrm{mm}]$ | 1180 |
| open air installation | $[\mathrm{A}]$ | 630 |
| underground installation $\rho=1^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 561 |
| underground installation $\rho=2^{\circ} \mathrm{Cm} / \mathrm{W}$ | $[\mathrm{A}]$ | 435 |

Table A1.22.b: Cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 300 \mathrm{~mm}^{2}$ : coefficients $a_{0}, a_{1}, a_{2}$, $a_{3}$ for different pitch values.

| helix radius | helix $\boldsymbol{p i t c h}$ | coefficient | coefficient |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\boldsymbol{\alpha}[\mathbf{m}]$ | $\boldsymbol{p}[\mathbf{m}]$ | $\boldsymbol{a}_{\boldsymbol{0}}$ | $\boldsymbol{a}_{\boldsymbol{1}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{2}}$ | coefficient <br> $\boldsymbol{a}_{\boldsymbol{3}}$ |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 0}$ | -0.6152 | -6.6444 | 0.1 | 1.70 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 2}$ | -0.4978 | -5.6066 | 0.1 | 1.76 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 4}$ | -0.3930 | -4.888 | 0.1 | 1.81 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 6}$ | -0.296 | -4.369 | 0.1 | 1.85 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{1 . 8}$ | -0.2094 | -3.8904 | 0.1 | 1.90 |
| $\mathbf{0 . 0 2 4}$ | $\mathbf{2 . 0}$ | -0.1268 | -3.5513 | 0.1 | 1.93 |

## APPENDIX 2

## MAGNETIC FiELD ISOLINES

In Appendix 2 the magnetic field isolines of $1.0 \mu \mathrm{~T}, 3.0 \mu \mathrm{~T}$ ('quality objective'), 5 $\mu \mathrm{T}, 10 \mu \mathrm{~T}$ ('attention value') and $100 \mu \mathrm{~T}$ ('exposure limit') have been calculated and plotted for some frequently used underground cable types, such as ARE4H1RX 12/20 kV [64], ARG7H1RX $12 / 20 \mathrm{kV}$ [63] and RG7H1OR $12 / 20 \mathrm{kV}$ [66], in order to provide a human exposure protection map concerning the magnetic fields generated by these cables carrying current equal at their ampacity. The calculation of the magnetic field isolines was implemented in Matlab ${ }^{\mathrm{TM}}$ environment and based on the exact formulation (eq. 2.23) discussed in Chapter 2, and on formulas (A1.1), (A1.2) and (A1.3) discussed in Appendix 1.


Figure A2.1: Magnetic field isolines for the cable ARE4HIRX 12/20 $\mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}(\mathrm{I}=147 \mathrm{~A})$


Figure A2.2: Magnetic field isolines for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 x 50 \mathrm{~m}^{2}(I=174 \mathrm{~A})$.


Figure A2.3: Magnetic field isolines for the cable ARE4H1RX 12/20 kV $3 x 70 \mathrm{~mm}^{2}(\mathrm{I}=212 \mathrm{~A})$.


Figure A2.4: Magnetic field isolines for the cable ARE4H1RX $12 / 20 \mathrm{kV} 3 \times 95 \mathrm{~mm}^{2}$ (I=253 A).


Figure A2.5: Magnetic field isolines for the cable ARE4H1RX 12/20 kV $3 x 120 \mathrm{~mm}^{2}$ (I=288 A).


Figure A2.6: Magnetic field isolines for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 150 \mathrm{~mm}^{2}(I=322 \mathrm{~A})$.


Figure A2.7: Magnetic field isolines for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}(I=365 \mathrm{~A})$.


Figure A2.8: Magnetic field isolines for the cable ARE4H1RX $12 / 20 \mathrm{kV} 3 x 240 \mathrm{~mm}^{2}$ (I=423 A).


Figure A2.9: Magnetic field isolines for the cable ARE4HIRX $12 / 20 \mathrm{kV} 3 \times 300 \mathrm{~mm}^{2}$ (I=478 A).


Figure A2.10: Magnetic field isolines for the cable ARG7H1RX $12 / 20 \mathrm{kV} 3 x 70 \mathrm{~mm}^{2}$ (I=200 A).


Figure A2.11: Magnetic field isolines for the cable ARG7H1RX 12/20 $\mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}(I=280 \mathrm{~A})$.


Figure A2.12: Magnetic field isolines for the cable ARG7H1RX 12/20 kV $3 x 185 \mathrm{~mm}^{2}(I=360 \mathrm{~A})$.


Figure A2.13: Magnetic field isolines for the cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 25 \mathrm{~mm}^{2}(I=150 \mathrm{~A})$.


Figure A2.14: Magnetic field isolines for the cable RG7HIOR 12/20 $\mathrm{kV} 3 \times 35 \mathrm{~mm}^{2}(\mathrm{I}=175 \mathrm{~A})$.


Figure A2.15: Magnetic field isolines for the cable RG7HIOR 12/20 $\mathrm{kV} 3 \times 50 \mathrm{~mm}^{2}(\mathrm{I}=207 \mathrm{~A})$.


Figure A2.16: Magnetic field isolines for the cable RG7HIOR 12/20 kV $3 x 70 \mathrm{~mm}^{2}(I=253 \mathrm{~A})$.


Figure A2.17: Magnetic field isolines for the cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 95 \mathrm{~mm}^{2}$ ( $I=300 \mathrm{~A}$ ).


Figure A2.18: Magnetic field isolines for the cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 120 \mathrm{~mm}^{2}$ (I=342 A).


Figure A2.19: Magnetic field isolines for the cable RG7H1OR $12 / 20 \mathrm{kV} 3 \times 150 \mathrm{~mm}^{2}$ (I=381 A).


Figure A2.20: Magnetic field isolines for the cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 185 \mathrm{~mm}^{2}(I=431 \mathrm{~A})$.


Figure A2.21: Magnetic field isolines for the cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 240 \mathrm{~mm}^{2}(I=500 \mathrm{~A})$.


Figure A2.22: Magnetic field isolines for the cable RG7HIOR $12 / 20 \mathrm{kV} 3 \times 300 \mathrm{~mm}^{2}(I=561 \mathrm{~A})$.

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